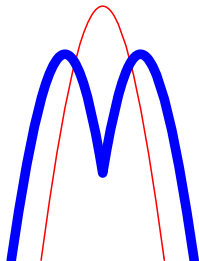


MOLPRO



User's Manual Version 2002.6

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Introduction to MOLPRO

MOLPRO is a complete system of *ab initio* programs for molecular electronic structure calculations, designed and maintained by H.-J. Werner and P. J. Knowles, and containing contributions from a number of other authors. As distinct from other commonly used quantum chemistry packages, the emphasis is on highly accurate computations, with extensive treatment of the electron correlation problem through the multiconfiguration-reference CI, coupled cluster and associated methods. Using recently developed integral-direct local electron correlation methods, which significantly reduce the increase of the computational cost with molecular size, accurate *ab initio* calculations can be performed for much larger molecules than with most other programs.

The heart of the program consists of the multiconfiguration SCF, multireference CI, and coupled-cluster routines, and these are accompanied by a full set of supporting features. The package comprises

- Integral generation for generally contracted symmetry adapted gaussian basis functions (*spdfghi*). There are two programs with identical functionality: the preferred code is SEWARD (R. Lindh) which is the best on most machines; ARGOS (R. M. Pitzer) is available as an alternative, and in some cases is optimum for small memory scalar machines. Also two different gradient integral codes, namely CADPAC (R. Amos) and ALASKA (R. Lindh) are available. Only the latter allows the use of generally contracted symmetry adapted gaussian basis functions.
- Effective Core Potentials (contributions from H. Stoll).
- Many one-electron properties.
- Some two-electron properties, e.g. $L_x^2, L_y^2, L_z^2, L_x L_y$ etc..
- Closed-shell and open-shell (spin restricted and unrestricted) self consistent field.
- Density-functional theory in the Kohn-Sham framework with various gradient corrected exchange and correlation potentials.
- Multiconfiguration self consistent field. This is the quadratically convergent MCSCF procedure described in J. Chem. Phys. 82 (1985) 5053. The program can optimize a weighted energy average of several states, and is capable of treating both completely general configuration expansions and also long CASSCF expansions as described in Chem. Phys. Letters 115 (1985) 259.
- Multireference CI. As well as the usual single reference function approaches (MP2, SDCI, CEPA), this module implements the internally contracted multireference CI method as described in J. Chem. Phys. 89 (1988) 5803 and Chem. Phys. Lett. 145 (1988) 514. Non variational variants (e.g. MR-ACPF), as described in Theor. Chim. Acta 78 (1990) 175, are also available. Electronically excited states can be computed as described in Theor. Chim. Acta, **84** 95 (1992).
- Multireference second-order and third-order perturbation theory (MR-PT2, MR-PT3) as described in Mol. Phys. **89**, 645 (1996) and J. Chem. Phys. **112**, 5546 (2000).
- Møller-Plesset perturbation theory (MPPT), Coupled-Cluster (CCSD), Quadratic configuration interaction (QCISD), and Brueckner Coupled-Cluster (BCCD) for closed shell systems, as described in Chem. Phys. Lett. 190 (1992) 1. Perturbative corrections for triple excitations can also be calculated (Chem. Phys. Letters **227** (1994) 321).

- Open-shell coupled cluster theories as described in *J. Chem. Phys.* **99** (1993) 5219, *Chem. Phys. Letters* **227** (1994) 321.
- Full Configuration Interaction. This is the determinant based benchmarking program described in *Comp. Phys. Commun.* **54** (1989) 75.
- Analytical energy gradients for SCF, DFT, state-averaged MCSCF/CASSCF, MP2 and QCISD methods.
- Analytical non-adiabatic coupling matrix elements for MCSCF.
- Valence-Bond analysis of CASSCF wavefunction, and energy-optimized valence bond wavefunctions as described in *Int. J. Quant. Chem.* **65**, 439 (1997).
- One-electron transition properties for MCSCF and MRCI wavefunctions.
- Spin-orbit coupling, as described in *Mol. Phys.*, **98**, 1823 (2000).
- Some two-electron transition properties for MCSCF wavefunctions (e.g., L_x^2 etc.).
- Population analysis.
- Orbital localization.
- Distributed Multipole Analysis (A. J. Stone).
- Automatic geometry optimization as described in *J. Comp. Chem.* **18**, (1997), 1473.
- Automatic calculation of vibrational frequencies, intensities, and thermodynamic properties.
- Reaction path following, as described in *Theor. Chem. Acc.* **100**, (1998), 21.
- Various utilities allowing other more general optimizations, looping and branching (e.g., for automatic generation of complete potential energy surfaces), general housekeeping operations.
- Geometry output in XYZ, MOLDEN and Gaussian formats; molecular orbital and frequency output in MOLDEN format.
- Integral-direct implementation of all Hartree-Fock, DFT and pair-correlated methods (MP, CCSD, MRCI etc.), as described in *Mol. Phys.*, **96**, (1999), 719. At present, perturbative triple excitation methods are not implemented.
- Local second-order Møller-Plesset perturbation theory (LMP2) as in *Chem. Phys. Lett.* **290**, 143 (1998), *J. Chem. Phys.* **111**, 5691 (1999), and *J. Chem. Phys.* **113**, 9443 (2000), (and references therein).
- Analytical energy gradients for LMP2, as described in *J. Chem. Phys.* **108**, (1998), 5185.
- Parallel execution on distributed memory machines, as described in *J. Comp. Chem.* **19**, (1998), 1215. At present, SCF, DFT, MRCI, MP2, LMP2, CCSD(T) energies and SCF, DFT gradients are parallelized when running with conventional integral evaluation; integral-direct SCF, DFT and LMP2 are also parallel.

The program is written mostly in standard Fortran-90. Those parts which are machine dependent are maintained through the use of a supplied preprocessor, which allows easy interconversion between versions for different machines. Each release of the program is ported and tested on a number of IBM RS/6000, Hewlett-Packard, Silicon Graphics, Compaq, and Linux

systems. A fuller description of the hardware and operating systems of these machines can be found at <http://www.molpro.net/machines.html>. The program additionally runs on Cray, Sun, Convex, Fujitsu and NEC SX4 platforms, as well as older architectures and/or operating systems from the primary list; however, testing is not carried out regularly on these systems, and hand-tuning of code may be necessary on porting. A large library of commonly used orbital basis sets is available, which can be extended as required. There is a comprehensive users' manual, which includes installation instructions. The manual is available in PostScript, PDF and also in HTML for mounting on a Worldwide Web server.

Future enhancements presently under development include

- Local coupled cluster theory (LCCSD) as described in *J. Chem. Phys.* **104**, (1996), 6286 and *J. Chem. Phys.* **114**, 661 (2001), with perturbative treatment of triple excitations, as described in *Chem. Phys. Letters* **318**, 370 (2000) and *J. Chem. Phys.* **113**, 9986 (2000).
- Enhancements to the efficiency of the DFT integration.
- Analytical energy gradients for CCSD, LCCSD, and CAS-PT2.
- Analytical second derivatives for SCF/MCSCF.
- Efficiency improvements for open-shell coupled cluster.
- Further parallelization.
- Open-shell MP2, LMP2 and LCCSD.

Some of these features will be included in the base version, whereas others will be available only as optional modules. The above list is for information only, and no representation is made that any of the above will be available within any particular time.

MOLPRO on the WWW

The latest information on MOLPRO, including program updates, can be found on the worldwide web at location <http://www.molpro.net/>.

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All publications resulting from use of this program must acknowledge the following.

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1 HOW TO READ THIS MANUAL

This manual is organized as follows: The next chapter gives an overview of the general structure of MOLPRO. It is essential for the new user to read this chapter, in order to understand the conventions used to define the symmetry, records and files, orbital spaces and so on. The later chapters, which describe the input of the individual program modules in detail, assume that you are familiar with these concepts. The appendices describe details of running the program, and the installation procedure.

Throughout this manual, words in `Typewriter Font` denote keywords recognized by MOLPRO. In the input, these have to be typed as shown, but may be in upper or lower case. Numbers or options which must be supplied by the user are in *italic*. In some cases, various different forms of an input record are possible. This is indicated as [*options*], and the possible options are described individually in subsequent subsections.

2 GENERAL PROGRAM STRUCTURE

This chapter gives an overview of the most important features of MOLPRO. For the new user, it is essential to understand the strategies and conventions described in this section, in particular the meaning of *files* and *records*, and the use of *symmetry*. This chapter will focus on general aspects; detailed information about each command will be given in later chapters. Information about commands and parameters can also be obtained using the MOLPRO help facility (see section 2.19).

2.1 Running MOLPRO

On Unix systems, MOLPRO is accessed using the *molpro* unix command. The syntax is

```
molpro [options] [datafile]
```

MOLPRO's execution is controlled by user-prepared data; if *datafile* is not given on the command line, the data is read from standard input, and program results go to standard output. Otherwise, data is taken from *datafile*, and the output is written to a file whose name is generated from *datafile* by removing any trailing suffix, and appending `.out`. If the output file already exists, then the old file is appended to the same name with suffix `.out_1`, and then deleted. This provides a mechanism for saving old output files from overwriting. Note that the above behaviour can be modified with the `-o` or `-s` options.

Unless disabled by options, the user data file is prepended by one or more default procedure files, if these files exist. These are, in order of execution, the file `molproi.rc` in the system directory containing the `molpro` command itself, `$HOME/.molproirc` and `./molproi.rc`.

2.1.1 Options

Most options are not required, since sensible system defaults are usually set. Options as detailed below may be given, in order of decreasing priority, on the command line, in the environment variable `MOLPRO`, or in the files `./molpro.rc`, `$HOME/.molprorc`, and `molpro.rc` in the system directory.

`-o | --output outfile` specifies a different output file.

- x | --executable *executable* specifies an alternative MOLPRO executable file.
- d | --directory *directory1:directory2...* specifies a list of directories in which the program will place scratch files. For detailed discussion of optimal specification, see the installation guide.
- s | --nobackup disables the mechanism whereby an existing output file is saved. --backup switches it on again.
- v | --verbose causes the procedure to echo debugging information; --noverbose selects quiet operation (default).
- e | --echo-procedures causes the contents of the default procedure files to be echoed at run time. --noecho-procedures selects quiet operation (default).
- f | --procedures enables the automatic inclusion of default procedure files (the default); --noprocedures disables such inclusion.
- g | --use-logfile causes some long parts of the program output, for example during geometry optimizations and finite-difference frequency calculations, to be diverted to an auxiliary output file whose name is derived from the output file by replacing its suffix (usually .out) by .log. --nouse-logfile disables this facility, causing all output to appear in the normal output file.
- m | --memory *memory* specifies the working memory to be assigned to the program, in 8-byte words. The memory may also be given in units of 1000 words by appending the letter k to the value, or in units of 1000000 with the key m, or 10^9 with g. K, B, G stand for 2^{10} , 2^{20} and 2^{30} .
- I | --main-file-repository *directory* specifies the directory where the permanent copy of any integral file (file 1) resides. This may be a pathname which is absolute or relative to the current directory (e.g., '.' would specify the current directory). Normally, the -I directory should be equal to the -d working directory to avoid copying of large integral files.
- W | --wavefunction-file-repository is similar to --wavefunction-file-repository except that it refers to the directory for the wavefunction files (2,3 and 4).
- L | --library *directory* specifies the directory where the basis set library files (LIBMOL*) are found.
- 1 | --file-1-directory *directory:directory:...* specifies the directory where the runtime file 1 will be placed, overriding --directory for this file only. -2, -3, -4, -5, -6, -7, -8 and -9 may be used similarly. Normally these options should not be given, since the program tries to use what is given in -d to optimally distribute the I/O.

There are a number of other options for tuning and system parameters, but these do not concern the general user.

It is not usually necessary to specify any of these options; the defaults are installation dependent and can be found in the system configuration file `molpro.rc` in the same directory as the `molpro` command itself.

2.1.2 Running MOLPRO on parallel computers

MOLPRO will run on distributed-memory multiprocessor systems, including workstation clusters, under the control of the Global Arrays parallel toolkit. There are also some parts of the code that can take advantage of shared memory parallelism through the OpenMP protocol, although these are somewhat limited, and this facility is not at present recommended. It should be noted that there remain some parts of the code that are not, or only partly, parallelized, and therefore run with replicated work. Additionally, some of those parts which have been parallelized rely on fast inter-node communications, and can be very inefficient across ordinary networks. Therefore some caution and experimentation is needed to avoid waste of resources in a multiuser environment.

Specifying parallel execution The following additional options for the `molpro` command may be used to specify and control parallel execution.

- n | --tasks *tasks/tasks_per_node:smp_threads* *tasks* specifies the number of Global Array processes to be set up, and defaults to 1. *tasks_per_node* sets the number of GA processes to run on each node, where appropriate. The default is installation dependent. In some environments (e.g., IBM running under Loadleveler; PBS batch job), the value given by `-n` is capped to the maximum allowed by the environment; in such circumstances it can be useful to give a very large number as the value for `-n` so that the control of the number of processes is by the batch job specification. *smp_threads* relates to the use of OpenMP shared-memory parallelism, and specifies the maximum number of OpenMP threads that will be opened, and defaults to 1. Any of these three components may be omitted, and appropriate combinations will allow GA-only, OpenMP-only, or mixed parallelism.
- N | --task-specification *user1:node1:tasks1,user2:node2:tasks2... node1, node2* etc. specify the host names of the nodes on which to run. On most parallel systems, *node1* defaults to the local host name, and there is no default for *node2* and higher. On Cray T3E and IBM SP systems, and on systems running under the PBS batch system, if `-N` is not specified, nodes are obtained from the system in the standard way. *tasks1, tasks2* etc. may be used to control the number of tasks on each node as a more flexible alternative to `-n / tasks_per_node`. If omitted, they are each set equal to `-n / tasks_per_node`. *user1, user2* etc. give the username under which processes are to be created. Most of these parameters may be omitted in favour of the usually sensible default values.
- G | --global-memory *memory* Some parts of the program make use of Global Arrays for holding and communicating temporary data structures. This option sets the amount of memory to allocate in total across all processors for such activities.

2.2 Input format

MOLPRO's execution is controlled by an input file. In general, each input record begins with a keyword, which may be followed by data or other keywords. The input is read sequentially by

a controlling program; when the controlling program calls a program module, this module continues to read the input file until it finds an unknown keyword. After the module has performed its function, control is returned to the controller.

The input file can be written in free format. The following conversions take place:

, (comma)	move to next tab stop, i.e. this delimits input fields
;(semicolon)	end of record, i.e. a new record is started
!(exclamation mark)	ignore rest of input line (useful for comments)
--- (three dashes)	end of file (rest of input is ignored)

You may type your input upper or lower case. The input processor will convert all characters to upper case. All integers are appended with “.” (only floating point numbers are read by the program).

Several logical input records can actually be typed on one line and separated by semicolons, i.e., a given input line may contain many actual commands (separated by semicolons), or just one, as you prefer. These basic command units (records) delimited by semicolons are also frequently referred to as *cards* throughout this manual.

Exception to these general rules are:

***	first data line always
INCLUDE	include other input file
FILE	definition of named files
TEXT	prints text
TITLE	defines a title for the run
CON	specifies orbital configurations
---	last line of input

These commands always occupy a whole line. Using INCLUDE it is possible to open secondary input files. If an INCLUDE command is encountered, the new input file is opened and read until its end. Input is then continued after the include card in the first file. INCLUDE's may be nested.

A MOLPRO input record (card) contains a number of input *fields*. Input fields may be up to 256 characters wide and contain either expressions or strings. The fields can be separated by commas or blanks. However, commas are required in certain cases to make the meaning unique. For instance, typing 3 +4 evaluates to one field with value 7, but 3 , +4 is the input for two fields with values 3 and 4. We recommend the general use of commas in order to avoid unexpected results.

Each line may start with a label. A label is separated from the body of the line by a colon (:). The colon is part of the label. The length of the label must not exceed 6 characters (including the colon) and the labels must be unique. Labels may be useful with GOTO commands. Example:

```
GOTO , START :
. . .
START : CCSD ( T )
```

Here START : is a label, and CCSD (T) is a command.

Strings containing blanks can be entered using quotes. For instance, 'This is a string' is interpreted as one string, but This is a string is a sequence of four strings in four subsequent fields. Strings in quotes are not converted to upper case.

Input lines may be concatenated using \ at the end of the line(s) to be continued. Any number of lines may be concatenated up to a total length of 1024 characters (only 500 characters are possible on older IBM systems).

Filenames may be up to 31 characters long, provided that long filenames are supported by the Unix system used. An exception are older CRAY systems, which allow only 8 characters for the names of binary MOLPRO files.

2.3 Input structure

A typical MOLPRO input has the following structure:

```

***,title                !title (optional)
memory,4,m               !memory specification (optional)
file,1,name.int          !permanent named integral file (optional)
file,2,name.wfu         !permanent named wavefunction file (optional)
gprint,options          !global print options (optional)
gthresh,options         !global thresholds (optional)
gdirect[,options]       !global direct (optional)
gexpec,opnames          !global definition of one-electron operators
basis=basisname         !basis specification. If not present, cc-pVDZ is used
var1=value,var2=value,... !setting variables for geometry and/or wavefunction definitions
geometry={...}          !geometry specification
program                 !program or procedure name
---                     !end of input (optional)

```

If the memory card is given, it should be the first card (after the optional title card). If any file cards are given, they should follow immediately. The order of basis, geometry, gprint, gdirect, gthresh, gexpec, and variable definitions is arbitrary. It is possible to call several programs one after each other. It is also possible to redefine basis set and/or geometry between the call to programs; the program will recognize automatically if the integrals have to be recomputed.

2.4 Expressions

In any input field, data can be entered in the form of expressions. Numbers and variables are special cases of expressions. An expression is typed in Fortran style and may contain any number of nested parenthesis. The standard intrinsic functions are also available (see next section). MOLPRO understands both arithmetic and logical expressions. The result of an arithmetic expression is a *real* (double precision) number. Internally, all integers are also converted to *real* numbers. The result of a logical expression is either `.TRUE.` or `.FALSE.` Internally, `.TRUE.` is stored as a one (1.0), and `.FALSE.` as zero (0.0). Expressions may contain any number of variables.

The following standard operations can be performed :

<i>expr</i> + <i>expr</i>	Addition
<i>expr</i> - <i>expr</i>	Subtraction
<i>expr</i> * <i>expr</i>	Multiplication
<i>expr</i> / <i>expr</i>	Division
<i>expr</i> .OR. <i>expr</i>	Logical OR
<i>expr</i> .AND. <i>expr</i>	Logical AND
<i>expr</i> .XOR. <i>expr</i>	Exclusive OR
.NOT. <i>expr</i>	Logical NOT
<i>expr</i> .GT. <i>expr</i>	Greater Than
<i>expr</i> .EQ. <i>expr</i>	Equal
<i>expr</i> .LT. <i>expr</i>	Less Than
<i>expr</i> .GE. <i>expr</i>	Greater Equal

$expr$.LE. $expr$	Less Equal
$expr$.NE. $expr$	Not Equal
$expr$ ** $expr$	Exponentiation
$expr$ ^ $expr$	Exponentiation
($expr$)	Parenthesis (no effect)
- $expr$	Change sign
+ $expr$	Keep sign (no effect)

2.5 Intrinsic functions

Expressions may contain the following intrinsic functions:

ABS ($expr$)	Absolute value
MAX ($expr, expr, \dots$)	Largest value of arbitrary number of numbers or expressions
MIN ($expr, expr, \dots$)	Smallest value of arbitrary number of numbers or expressions
EXP ($expr$)	Exponential
LOG ($expr$)	Natural Logarithm
LOG10 ($expr$)	Common Logarithm
SQRT ($expr$)	Square Root
NINT ($expr$)	Next nearest integer
INT ($expr$)	Truncate to integer
SIN ($expr$)	Sine
COS ($expr$)	Cosine
TAN ($expr$)	Tangent
ASIN ($expr$)	Arcsine
ACOS ($expr$)	Arccosine
ATAN ($expr$)	Arctangent
COSH ($expr$)	Hyperbolic cosine
SINH ($expr$)	Hyperbolic sine
TANH ($expr$)	Hyperbolic tangent
MOD ($expr1, expr2$)	Remainder: $expr1 - \text{INT}(expr1/expr2) * expr2$

Note: all trigonometric functions use or produce angles in degrees.

2.6 Files

MOLPRO uses three sequential text files, namely the *input file*, the *output file*, and the *punch file*. The punch file is a short form of the output which contains the most important data and results, such as geometries, energies, dipole moments, etc. The punch file can be processed by the separate program *READPUN*, which selects specific results by keywords and is able to produce ordered tables in user supplied format. Furthermore, there are up to 9 binary MOLPRO files available, each one known to the program simply by its number (1 to 9). By default, they

are temporary files, usually allocated dynamically by the program, but they can be connected to permanent files with the `FILE` command. Each file is direct access, and word addressable (word=64 bit usually), but is organised in *records* of any length. The name, address and length of each record is held in a directory at the start of the file.

File 1 is the *main file*, holding basis set, geometry, and the one and two electron integrals. By default, file 2 is the *dump file* and used to store the wavefunction information, i.e. orbitals, CI coefficients, and density matrices. File 3 is an auxiliary file which can be used in addition to file 2 for restart purposes. Often files 1 and 2 (and 3) are declared as permanent files (see `FILE`) to enable restarts. Storing the wavefunction information on file 2 is useful, since the integral file is overwritten at each new geometry, while the orbitals and CI coefficients of one calculation can be used as a starting guess for the next calculation at a neighbouring geometry. Files 4 to 8 are used as scratch space, e.g., for sorting the integrals, storage of transformed integrals and of the CI vectors. These files should normally not be made permanent.

2.7 Records

Record names are positive integers, and are usually referred to in the format *record.file*, e.g., 2100.2 means the record called 2100 on file 2. Note that these names are quite arbitrary, and their numerical values have nothing to do with the order of the records in the file. Record names ≤ 2000 are reserved for standard quantities (e.g. integrals, properties etc.) and you should never use these in an input, unless you know exactly what you are doing. Some important default records to remember are

2100	RHF dump record (closed and open-shell)
2200	UHF dump record
2140	MCSCF dump record
5000	MCSCF gradient information
5100	CP-MCSCF gradient information
5200	MP2 gradient information
5300	Frequencies restart information

If an input contains several wavefunction calculations of the same type, e.g., several MCSCF calculations with different active spaces, the record number will be increased by 1 for each calculation of the same type. Thus, the results of the first SCF calculation in an input are stored in dump record 2100.2, the second SCF in record 2101.2, the first MCSCF in 2140.2, the second MCSCF in 2141.2 and so on. Note that these numbers refer to the occurrence in the input and not on the order in which the calculations are performed in the actual run. If an input or part of it is repeated using `DO` loops, this ensures that each calculation will start with the orbitals from the corresponding orbitals from the previous cycle, as long as the order of the commands in the input remains unchanged. If for instance the first SCF would be skipped in the second cycle using some `IF / ENDIF` structure, the second SCF would still use record 2101.2. Thus, under most circumstances the program defaults are appropriate, and the user does not have to specify the records.

After a restart this logic will still work correctly if the number and sequence of `SCF` and `MCSCF` commands is kept unchanged. Thus, if you want to skip certain parts of the input after a restart, it is recommended to use `IF / ENDIF` structures or the `GOTO` command rather than to delete or comment certain commands. If for some reason this is not possible, the `START` and `ORBITAL` directives can be used to specify explicitly the records to be used.

In general we recommend the use of program defaults whenever possible, since this minimizes the probability of input errors and frustration!

After completion of each program step, MOLPRO prints a summary of the records on each file.

2.8 Restart

Information from the permanent files is automatically recovered in subsequent calculations. This can be controlled using the RESTART directive.

2.9 Data set manipulation

It is possible to truncate files and rename or copy records using the DATA command. Several standard matrix operations can be performed with MATROP, e.g., printing records, linearly combining or multiplying matrices, or forming the trace of a product of two matrices.

2.10 Memory allocation

MOLPRO allocates memory dynamically as required by the user on the MEMORY card. Thus it is not necessary to maintain different versions of the program with different memory sizes. If the MEMORY command is omitted, the program will use a default memory size, which depends on the hardware used and how the program was installed. Note that, on Unix machines, the default memory can be set on the molpro command line using the flag -m.

2.11 Variables

The program maintains a set of internal *variables*. These may be used in place of floating point numbers anywhere in the input. Before their use, variables must be defined as described in detail in Section 6. Variables are useful for running the same input with different actual parameters (e.g. geometries or basis function exponents), and to store and manipulate the results. They can also be used to change record names automatically when several geometries are calculated in one run. It is thus possible to save the information for each geometry separately in different records and variables. *Arrays* are variables with an index in parenthesis, e.g., *var(1)*. The number of elements in an array *var* is *#var*. The array length can be reset to zero by the CLEAR directive or simply by modifying *#var*. Variables and variable arrays may be displayed at any place in the output by the SHOW command, and whole tables of variables can be generated using the TABLE command. For more details about variables see section 6.

Variables can also be used for passing input parameters to the program. This is useful for *procedures*, which are described in Section 4.8.

2.12 Multiple passes through the input

It is possible to perform loops over parts of the input using DO loops, very much as in FORTRAN programs. DO loops may be nested to any reasonable depth. This can be conveniently used, for instance, to compute automatically whole potential energy surfaces.

Table 1: The symmetry generators for the point groups

Generators	Point group
(null card)	C_1 (i.e. no point group symmetry)
X (or Y or Z)	C_s
XY	C_2
XYZ	C_i
X, Y	C_{2v}
XY, Z	C_{2h}
XZ, YZ	D_2
X, Y, Z	D_{2h}

2.13 Symmetry

MOLPRO can use Abelian point group symmetry only. For molecules with degenerate symmetry, an Abelian subgroup must be used — e.g., C_{2v} or D_{2h} for linear molecules. The symmetry group which is used is defined in the integral input by combinations of the symmetry elements x , y , and z , which specify which coordinate axes change sign under the corresponding generating symmetry operation. It is usually wise to choose z to be the unique axis where appropriate (essential for C_2 and C_{2h}). The possibilities in this case are shown in Table 1.

Normally, MOLPRO determines the symmetry automatically, and rotates and translates the molecule accordingly. However, explicit symmetry specification is sometimes useful to fix the orientation of the molecule or to use lower symmetries.

The irreducible representations of each group are numbered 1 to 8. Their ordering is important and given in Tables 2 – 4. Also shown in the tables are the transformation properties of products of x , y , and z . s stands for an isotropic function, e.g., s orbital, and for these groups, this gives also the transformation properties of x^2 , y^2 , and z^2 . Orbitals or basis functions are generally referred to in the format *number.irrep*, i.e. 3.2 means the third orbital in the second irreducible representation of the point group used.

2.14 Defining the wavefunction

In all program modules where such information is required, the total symmetry of the N -electron wavefunction is defined on WF (wavefunction) cards in the following way:

WF,*nelec,irrep,spin*

or, alternatively

WF,[NELEC=*nelec*],[SYM[METRY]=*irrep*],[spin=*spin*],[CHARGE=*charge*]

where *nelec* is the total number of electrons, *irrep* is the number of the irreducible representation, and *spin* equals $2 \times S$, where S is the total spin quantum number. Instead of *nelec* also *charge* can be given, which specifies the total charge of the molecule. For instance, for a calculation in C_{2v} symmetry with 10 electrons, WF , 10 , 3 , 0 denotes a 1B_2 state, and WF , 10 , 1 , 2 a 3A_1 state. The charge can also be defined by setting the variable CHARGE:

Table 2: Numbering of the irreducible representations in D_{2h}

D_{2h}		
No.	Name	Function
1	A_g	s
2	B_{3u}	x
3	B_{2u}	y
4	B_{1g}	xy
5	B_{1u}	z
6	B_{2g}	xz
7	B_{3g}	yz
8	A_u	xyz

Table 3: Numbering of the irreducible representations in the four-dimensional groups

No.	C_{2v}		C_{2h}		D_2	
	Name	Function	Name	Function	Name	Function
1	A_1	s, z	A_g	s, xy	A	s
2	B_1	x, xz	A_u	z	B_3	x, yz
3	B_2	y, yz	B_u	x, y	B_2	y, xz
4	A_2	xy	B_g	xz, yz	B_1	xy

Table 4: Numbering of the irreducible representations in the two-dimensional groups

No.	C_s		C_2		C_i	
	Name	Function	Name	Function	Name	Function
1	A'	s, x, y, xy	A	s, z, xy	A_g	s, xy, xz, yz
2	A''	z, xz, yz	B	x, y, xz, yz	A_u	x, y, z

SET, CHARGE=*charge*

This charge will be used in all energy calculations following this input. Note that SET is required, since CHARGE is a system variable (cf. section 6.3).

Although in principle each program unit requires a WF command, in practice it is seldom necessary to give it. The program remembers the information on the WF card, and so one might typically specify the information in an SCF calculation, but then not in subsequent MCSCF or CI calculations; this also applies across restarts. Furthermore, *nelec* defaults to the sum of the nuclear charges, *irrep* to 1 and *spin* to 0 or 1; thus in many cases, it is not necessary to specify a WF card at all.

2.15 Defining orbital subspaces

In the SCF, MCSCF, and CI programs it may be necessary to specify how many orbitals in each symmetry are *occupied* (or *internal* in CI), and which of these are *core* or *closed shell* (doubly occupied in all CSFs). This information is provided on the OCC, CORE, and CLOSED cards in the following way:

OCC, m_1, m_2, \dots, m_8 ; CORE, co_1, co_2, \dots, co_8 ; CLOSED, cl_1, cl_2, \dots, cl_8 ;

where m_i is the number of occupied orbitals (including core and closed), co_i the number of core orbitals, and cl_i is the number of closed-shell orbitals (including the core orbitals) in the irreducible representation i . In general, $m_i \geq cl_i$, and $cl_i \geq co_i$. It is assumed that these numbers refer to the first orbitals in each irrep.

Note that the OCC, CORE, and CLOSED cards have slightly different meanings in the SCF, MCSCF and CI or CCSD programs. In SCF and MCSCF, *occupied* orbitals are those which occur in any of the CSFs. In electron correlation methods (CI, MPn, CCSD etc), however, OCC denotes the orbitals which are occupied in any of the *reference* CSFs. In the MCSCF, *core* orbitals are doubly occupied in all CSFs and *frozen* (not optimized), while *closed* denotes all doubly occupied orbitals (frozen plus optimized). In the CI and CCSD programs, *core* orbitals are those which are not correlated and *closed* orbitals are those which are doubly occupied in all reference CSFs.

OCC, CORE and CLOSED commands are generally required in each program module where they are relevant; however, the program remembers the most recently used values, and so the commands may be omitted if the orbital spaces are not to be changed from their previous values. Note that this information is also preserved across restarts. Note also, as with the WF information, sensible defaults are assumed for these orbital spaces. For full details, see the appropriate program description.

2.16 Selecting orbitals and density matrices (ORBITAL, DENSITY)

As outlined in section 2.7, the information for each SCF or MCSCF calculation is stored in a *dump record*. Dump records contain orbitals, density matrices, orbital energies, occupation numbers, fock matrices and other information as wavefunction symmetries etc. Subsequent calculation can access the orbitals and density matrices from a particular record using the ORBITAL and DENSITY directives, respectively. These input cards have the same structure in all programs. The general format of the ORBITAL and DENSITY directives is as follows.

ORBITAL[, [RECORD=*record*] [, [TYPE=*orbtype*] [, STATE=*state*] [, SYM[METRY]=*symmetry*] [, SPIN=*spin*] [, MS2=*ms2*] [, [N]ELEC=*nelec*] [, SET=*iset*]

DENSITY[, [RECORD=*record*] [, [TYPE=*dentype*] [, STATE=*state*] [, SYM[METRY]=*symmetry*] [, SPIN=*spin*] [, MS2=*ms2*] [, [N]ELEC=*nelec*] [, SET=*iset*]

where the (optional) specifications can be used to select specific orbitals, if several different orbital sets are stored in the same record. The meaning of the individual specifications are as follows:

<i>orbtype</i>	Orbital type. This can be one of CANONICAL: canonical or pseudocanonical orbitals; NATURAL: natural orbitals; LOCAL: localized orbitals; LOCAL (PM) : localized Pipek-Mezey orbitals; LOCAL (BOYS) : localized Boys orbitals; PROJECTED: projected virtual orbitals used in local calculations. Without further specification, the most recently computed orbitals of the specified type are used. If the orbital type is not specified, the program will try to find the most suitable orbitals automatically. For instance, in MRCI calculations NATURAL orbitals will be used preferentially if available; MRPT (CASPT2) calculations will first search for CANONICAL orbitals, and local calculations will first look for LOCAL orbitals. Therefore, in most cases the orbital type needs not to be specified.
<i>state</i>	Specifies a particular state in the form <i>istate.isym</i> . For instance, 2.1 refers to the second state in symmetry 1. This can be used if density matrices or natural orbitals have been computed for different states in a state-averaged CASSCF calculation. If not given, the state-averaged orbitals are used. The specification of <i>isym</i> is optional; it can also be defined using the SYMMETRY key.
<i>dentype</i>	Density type. This can be one of CHARGE: charge density; SPIN: UHF spin density; TRANSITION: transition density matrix; The default is CHARGE.
<i>symmetry</i>	Specifies a particular state symmetry. Alternatively, the state symmetry can be specified using STATE (see above).
<i>spin</i>	Spin quantum number, i.e. 0 for singlet, 1/2 for doublet, 1 for triplet, etc. Alternatively MS2 can be used.
<i>ms2</i>	$2M_S$, i.e. 0 for singlet, 1 for doublet, 2 for triplet etc. Alternatively, SPIN can be used.
<i>nelec</i>	Number of electrons.
<i>iset</i>	Set number of orbitals. The orbital sets are numbered in the order they are stored.

If any of the above options are given, they must be obeyed strictly, i.e., the program aborts if the request cannot be fulfilled.

Examples:

```

ORBITAL,2100.2           !Use SCF orbitals
ORBITAL,2140.2           !Use (state-averaged) MCSCF orbitals
ORBITAL,2140.2,CANONICAL !use canonical MCSCF orbitals
ORBITAL,2140.2,NATURAL,STATE=2.1 !use natural MCSCF orbitals for second state in sym. 1

```


2.17 Summary of keywords known to the controlling program

This is a summary of all keywords presently implemented in the controlling program. Each module knows further keywords, which are described in the chapters about the individual programs. For detailed information about the use of the commands listed below, consult the following chapters.

Program control:

***	indicates start of a new calculation
MEMORY	allocates dynamic memory
PUNCH	opens a punch file
FILE	connects units to permanent files
RESTART	recovers file information
INCLUDE	includes other input files
BASIS	can be used to define default basis sets
GEOMETRY	can be used to specify the geometry
ZMAT	can be used to define the Z-matrix
PARALLEL	can be used to control parallelization
STATUS	checks status of program steps
PRINT,GPRINT	controls global print levels
THRESH,GTHRESH	controls global thresholds
DIRECT,GDIRECT	flags direct computation of integrals and for setting direct options
EXPEC,GEXPEC	controls computation of expectation values
TEXT	prints text
EXIT	stops execution
DO	controls do loops
ENDDO	end of do loops
IF	controls conditional actions
ELSEIF	controls conditional actions
ENDIF	end of IF block
GOTO	used to skip part of input and for loops over input
LABEL	no action
DATA	data set management
DELETE, ERASE	data set deletion
MATROP	performs matrix operations
GRID	Define grid
CUBE	Dump data to grid
CARTESIAN	Use cartesian basis functions
SPHERICAL	Use spherical harmonic basis functions
USER	calls user-supplied subroutine
---	last line of input

Variables:

SET	sets variables (obsolete)
SETI	sets variables or numbers to their inverse (obsolete)
SETA	sets variable arrays (obsolete)
CLEAR	clears variables

CLEARALL	clears all variables
GETVAR	recovers variables from file
SHOW	displays the values of variables
TABLE	prints tables
<i>Wave function optimization:</i>	
INT	calls the machine default integral program. This is optional and needs not to be given.
INTS	calls SEWARD integral program
INTE or INTEGRAL	calls Pitzer's integral program
INTD	flags direct computation of integrals (obsolete, please use GDIRECT instead).
LSINT	calls the spin-orbit integral program
SORT	calls two-electron sorting program
CPP	compute core polarization potential integrals
HF, RHF, HF-SCF, or RHF-SCF	calls spin-restricted Hartree-Fock program (open or closed shell)
UHF or UHF-SCF	calls spin-unrestricted Hartree-Fock program
DFT	calls the density functional program
KS, RKS	call the Kohn-Sham spin restricted density functional program
UKS	call the Kohn-Sham spin-unrestricted density functional program
MULTI, MCSCF, or CASSCF	calls MCSCF/CASSCF program
CASVB	calls the CASVB valence bond program
CI, MRCI, or CI-PRO	calls internally contracted MRCI program
ACPF, AQCC	calls internally contracted MR-ACPF program
CEPA	calls single-reference CEPA program (closed- or open-shell)
RS2, RS3	calls internally contracted multireference perturbation theory
MP2	calls closed-shell MP2 program
MP3	calls closed-shell MP3 program
MP4	calls closed-shell MP4 program
CISD	calls closed-shell CISD program
CCSD	calls closed-shell coupled cluster program
BCCD	calls closed-shell Brueckner CCD program
QCI, QCSID	calls closed-shell quadratic configuration interaction program
UCCSD	calls spin-unrestricted open-shell coupled cluster program
RCCSD	calls spin-restricted open-shell coupled cluster program
FCI or FULLCI	calls determinant based full CI program
<i>Orbital manipulation:</i>	
LOCALI	calls orbital localization program
MERGE	calls orbital manipulation program
<i>Properties and wavefunction analysis:</i>	

POP	calls population analysis program
DMA	calls distributed multipole analysis program
PROPERTY	calls properties program
DIP	adds dipole field to h
QUAD	adds quadrupole field to h
PLOT	calls orbital and density plot program
IGLO, PIGLO, NMR	call magnetic property programs
<i>Gradients and geometry optimization:</i>	
FORCES	calls gradient program
OPT	calls geometry optimization program
OPTG	performs automatic geometry optimization
MIN	performs energy minimization with respect to some parameters
PUT	print or write geometry to a file
HESSIAN	calculate hessian
FREQUENCY	calculate vibrational frequencies
MASS	define atomic masses
DDR	evaluates approximate non-adiabatic coupling matrix elements

2.18 Default procedures

For convenience of use, MOLPRO provides a number of default procedures for standard applications. Each procedure performs a full calculation and prints a summary of the computed energies at the end of the output. It is possible to call several procedures one after the other in the same job.

The procedures are provided in the file `molproi.rc`, which is automatically included at the beginning of each input when the unix `molpro` command is used. Inclusion of the procedures can be avoided using the `molpro -f` option. The user can also define his own procedures in his `molproi.rc` file. For details, see section 4.8.

2.18.1 Procedures for energy calculations

The wavefunction symmetry and spin can be modified using the variables `SYMMETRY` and `SPIN`, respectively. The number of electrons can be modified using the variable `NELEC`. SCF calculations are only done if no orbitals are available or if the symmetry or spin changed since the last SCF optimization. In CASPT and MRCI calculations the CASSCF is only done if the last optimized orbitals are not of MCSCF type.

<code>runscf</code>	Performs SCF calculation.
<code>rundft</code>	Performs DFT calculation. The functional can be specified using either the <code>FUNCTIONAL</code> or <code>DFTNAME</code> variable (default=B3LYP).
<code>runmp2</code>	Performs SCF and MP2 calculations.
<code>runmp3</code>	Performs SCF, MP2 and MP3 calculations.
<code>runmp4</code>	Performs SCF and MP2-MP4 calculations.

<code>runmp4sdq</code>	Same as <code>runmp4</code> , but without triples.
<code>runccsd</code>	Performs SCF and CCSD calculations. Uses RCCSD in open-shell cases.
<code>runccsdt</code>	Performs SCF and CCSD(T) calculations. Uses RCCSD(T) in open-shell cases.
<code>runuccsd</code>	Performs SCF and CCSD calculations. Uses UCCSD in open-shell cases.
<code>runuccsdt</code>	Performs SCF and CCSD(T) calculations. Uses UCCSD(T) in open-shell cases.
<code>runbccd</code>	Performs SCF and BCCD calculations (closed-shell only).
<code>runbccdt</code>	Performs SCF and BCCD(T) calculations (closed-shell only).
<code>runqcisd</code>	Performs SCF and QCISD calculations (closed-shell only).
<code>runqcisdt</code>	Performs SCF and QCISD(T) calculations (closed-shell only).
<code>runcas</code>	Performs SCF and CASSCF calculation. The wavefunctions for state averaged CASSCF calculations can be defined using the variables SYMMETRY, SPIN, STATE, WEIGHT, and NELEC.
<code>caspt2</code>	Performs SCF, CASSCF, and CASPT2 calculations. In case of state-averaged CASSCF reference functions (see <code>runcas</code>) the CASPT2 is performed for each state separately.
<code>caspt3</code>	As CASPT2, but also computes third-order energy.
<code>runcaspt2</code>	Same as <code>caspt2</code> .
<code>runcaspt3</code>	Same as <code>caspt3</code> .
<code>runmrci</code>	Performs SCF, CASSCF, and MRCI calculations. In case of state-averaged CASSCF reference functions (see <code>runcas</code>) the MRCI is performed for each state symmetry separately. Several states in the same symmetry are treated simultaneously. Transition moments are automatically computed between all states. However, these are not printed in the summary at the end of the output and must be extracted from the output or punch file.
<code>runacpf</code>	Performs SCF, CASSCF, and MR-ACPF calculations. At present, multiple-state MR-ACPF calculations are not possible.

2.18.2 Procedures for geometry optimizations

If the geometry is given as Z-matrix and depends on variables, only these are optimized and other numerical parameters are kept fixed. This behaviour can be modified by setting the variable `OPTFULL= .TRUE. .`, which causes the geometry optimization be done in cartesian coordinates for all degrees of freedom. In this case the variables are not modified. If the geometry is given as XYZ-form or the Z-matrix does not depend on variables, all degrees of freedom are optimized.

<code>optscf</code>	Performs SCF geometry optimization.
<code>optdft</code>	Performs DFT geometry optimization. The functional can be specified using either the <code>FUNCTIONAL</code> or <code>DFTNAME</code> variable (default=B3LYP).
<code>optmp2</code>	Performs MP2 geometry optimization.
<code>optcas</code>	Performs CASSCF geometry optimization. The procedure does not support state-averaged calculations.

2.18.3 Procedures for frequency calculations

In all cases the geometry is optimized for all degrees of freedom before computing the frequencies. Variables on which the Z-matrix depends are not modified. The hessian is computed by finite differences from analytical energy gradients. Frequencies, intensities, and thermodynamic quantities are evaluated and printed in the output. The summary at the end shows only the energies.

<code>freqscf</code>	Performs SCF frequency calculation.
<code>freqdft</code>	Performs DFT frequency calculation.
<code>freqmp2</code>	Performs MP2 frequency calculation.
<code>freqcas</code>	Performs CASSCF frequency calculation. The procedure does not support state-averaged calculations.

Note: in most cases the SCF calculation is not performed if starting orbitals are found. Several procedures may be specified after each other.

2.19 MOLPRO help

The help command can be used to obtain a short description of commands, input parameters, and variables. The syntax is:

```
HELP,set,name,[keys]
```

where *set* is either COMMAND, VARIABLE, or the name of the input set (e.g., THRESH, PRINT, LOCAL, EOM, CFIT), and *name* is the name of the parameter. If *name* is blank, all parameters of the set are shown. Optionally, *keys* can be specified to request specific information (e.g., short_description, long_description, default_value, type, program). If *keys* are not given, short_description is assumed.

Currently, help is only available for a limited number of parameters and commands. However, the database will be extended in the near future.

3 INTRODUCTORY EXAMPLES

This section explains some very simple calculations in order to help the new user to understand how easy things can be.

3.1 Using the molpro command

1. Perform a simple SCF calculation for molecular hydrogen. The input is typed in directly and the output is sent to the terminal:

```
molpro <<!
basis=vdz;
geometry={angstrom;h1;h2,h1,.74}
hf
!
```

2. The same calculation, with the data taken from the file `h2.com`. The output is sent to `h2.out`. On completion, the file `h2.pun` is returned to the current directory and the file `h2.wf` to the directory `$HOME/wfu` (this is the default):

```
molpro h2.com
```

`h2.com` contains:

```
***,H2
file,2,h2.wf;
punch,h2.pun;
basis=vdz;
geometry={angstrom;h1;h2,h1,.74}
hf
```

examples/
h2.com

3. As before, but the file `h2.wf` is sent to the directory `/tmp/wfu`:

```
molpro -W /tmp/wfu h2.com
```

3.2 Simple SCF calculations

The first example does an SCF calculation for H_2O , using all possible defaults.

```
***,h2o !A title
r=1.85,theta=104 !set geometry parameters
geometry={0; !z-matrix geometry input
H1,O,r;
H2,O,r,H1,theta}
hf !closed-shell scf
```

examples/
h2o_scf.com

In the above example, the default basis set (VDZ) is used. We can modify the default basis using a `BASIS` directive.

```
***,h2o cc-pVTZ basis !A title
r=1.85,theta=104 !set geometry parameters
geometry={0; !z-matrix geometry input
H1,O,r;
H2,O,r,H1,theta}
basis=VTZ !use VTZ basis
hf !closed-shell scf
```

examples/
h2o_scf_vtz.com

3.3 Geometry optimizations

Now we can also do a geometry optimization, simply by adding the card `OPTG`.

```
***,h2o !A title
r=1.85,theta=104 !set geometry parameters
geometry={0; !z-matrix geometry input
H1,O,r;
H2,O,r,H1,theta}
basis=6-31g** !use Pople basis set
hf !closed-shell scf
optg !do scf geometry optimization
```

examples/
h2o_scfopt_631g.com

3.4 CCSD(T)

The following job does a CCSD(T) calculation using a larger (VTZ) basis (this includes an f function on oxygen and a d function on the hydrogens).

```

***,h2o                !A title
r=1.85,theta=104       !set geometry parameters
geometry={O;           !z-matrix geometry input
          H1,O,r;
          H2,O,r,H1,theta}
basis=VTZ              !use VTZ basis
hf                    !closed-shell scf
ccsd(t)               !do ccsd(t) calculation

```

examples/
h2o_ccsdt_vtz.com

3.5 CASSCF and MRCI

Perhaps you want to do a CASSCF and subsequent MRCI for comparison. The following uses the full valence active space in the CASSCF and MRCI reference function.

```

***,h2o                !A title
r=1.85,theta=104       !set geometry parameters
geometry={O;           !z-matrix geometry input
          h1,O,r;
          h2,O,r,H1,theta}
basis=vtz              !use VTZ basis
hf                    !closed-shell scf
ccsd(t)               !do ccsd(t) calculation
casscf                !do casscf calculation
mrci                  !do mrci calculation

```

examples/
h2o_mrci_vtz.com

3.6 Tables

You may now want to print a summary of all results in a table. To do so, you must store the computed energies in variables:

```

***,h2o                !A title
r=1.85,theta=104       !set geometry parameters
geometry={O;           !z-matrix geometry input
          h1,O,r;
          h2,O,r,H1,theta}
basis=vtz              !use VTZ basis
hf                    !closed-shell scf
e(1)=energy            !save scf energy in variable e(1)
method(1)=program      !save the string 'HF' in variable method(1)

ccsd(t)               !do ccsd(t) calculation
e(2)=energy            !save ccsd(t) energy in variable e(2)
method(2)=program      !save the string 'CCSD(T)' in variable method(2)

casscf                !do casscf calculation
e(3)=energy            !save scf energy in variable e(3)
method(3)=program      !save the string 'CASSCF' in variable method(3)
mrci                  !do mrci calculation
e(4)=energy            !save scf energy in variable e(4)
method(4)=program      !save the string 'MRCI' in variable method(4)

table,method,e         !print a table with results
title,Results for H2O, basis=$basis !title for the table

```

examples/
h2o_table.com

This job produces the following table:

Results for H2O, basis=VTZ

METHOD	E
HF	-76.05480122
CCSD(T)	-76.33149220
CASSCF	-76.11006259
MRCI	-76.31960943

3.7 Procedures

You could simplify this job by defining a procedure `SAVE_E` as follows:

```

proc save_e                !define procedure save_e
if(#i.eq.0) i=0            !initialize variable i if it does not exist
i=i+1                     !increment i
e(i)=energy                !save scf energy in variable e(i)
method(i)=program         !save the present method in variable method(i)
endproc                   !end of procedure

***,h2o                   !A title

r=1.85,theta=104          !set geometry parameters
geometry={o;              !z-matrix geometry input
          h1,O,r;
          h2,O,r,H1,theta}
basis=vtz                 !use VTZ basis
hf                         !closed-shell scf
save_e                     !call procedure, save results
ccsd(t)                   !do ccscd(t) calculation
save_e                     !call procedure, save results
casscf                    !do casscf calculation
save_e                     !call procedure, save results
mrci                      !do mrci calculation
save_e                     !call procedure, save results

table,method,e            !print a table with results
title,Results for H2O, basis=$basis !title for the table

```

examples/
h2o_proce.com

The job produces the same table as before. If you put the procedure `SAVE_E` in a file `molproi.rc` or `$HOME/.molproirc`, it would be automatically included in all your jobs (`./molproi.rc` is searched first; if this file does not exist, molpro looks for `$HOME/.molproirc`. If this also does not exist, molpro uses the default file in the system directory).

3.8 Using default Procedures

Using the default procedures described in section 2.18, the above inputs can be simplified even more.

3.8.1 RCCSD(T) for different states

The following calculation performs RCCSD(T) calculations for the ground and first excited state of the OH radical.

```

***,OH
geometry={O;H,O,1.83}
runccsd
set,symmetry=2,spin=1 ! X 2 Pi state
runccsd
epi=energy
set,symmetry=1,spin=1 ! A 2 Sigma+ state
runccsd
esig=energy
de=(esig-epi)*toev !Excitation energy in eV

```

examples/
oh_runccsd.com

This produces the following results:

RESULTS

METHOD	STATE	S	ENERGY
HF-SCF	1.2	0.5	-75.39004124
RCCSD	1.2	0.5	-75.55736436
RCCSD(T)	1.2	0.5	-75.55912676
RCCSD[T]	1.2	0.5	-75.55916157
HF-SCF	1.1	0.5	-75.22787407
RCCSD	1.1	0.5	-75.39738157
RCCSD(T)	1.1	0.5	-75.39914839
RCCSD[T]	1.1	0.5	-75.39915981

SETTING DE = 4.35323484 EV

3.8.2 SA-CASSCF and MRCI

The following calculation performs state-averaged CASSCF and subsequent MRCI calculations for the ground and first excited state of the OH radical. A full valence active space is used.

```

***,OH
geometry={O;H,O,1.83}
set,symmetry=[2,3,1],spin=1 ! 2 Pix, 2Piy, and 2 Sigma+ states
runmrci ! SA-CASSCF and MRCI

```

examples/
oh_runmrci.com

The following table is printed at the end of the output:

RESULTS

METHOD	STATE	S	ENERGY	DIPX	DIPY	DIPZ
CASSCF	1.2	0.5	-75.41331789	0.0	0.0	0.67158730
CASSCF	1.3	0.5	-75.41331789	0.0	0.0	0.67158730
CASSCF	1.1	0.5	-75.24125256	0.0	0.0	0.69975340
MRCI	1.2	0.5	-75.55518444	0.0	0.0	0.66457191
MRCI+D	1.2	0.5	-75.56014871	0.0	0.0	0.66457191
MRCI+P	1.2	0.5	-75.55853208	0.0	0.0	0.66457191
MRCI	1.3	0.5	-75.55518444	0.0	0.0	0.66457191
MRCI+D	1.3	0.5	-75.56014871	0.0	0.0	0.66457191
MRCI+P	1.3	0.5	-75.55853208	0.0	0.0	0.66457191
MRCI	1.1	0.5	-75.39442202	0.0	0.0	0.70484623
MRCI+D	1.1	0.5	-75.40040680	0.0	0.0	0.70484623
MRCI+P	1.1	0.5	-75.39846312	0.0	0.0	0.70484623

This calculation performs MRCI calculations for both $^2\Pi_x$ and $^2\Pi_y$. The procedure is not clever enough to recognize that they are degenerate. However, one can easily modify the input to eliminate this drawback.

```

***,OH
geometry={O;H,O,1.83}
set,symmetry=[2,3,1],spin=1 ! 2Pix, 2Piy, and 2Sigma+ states (specifying spin=1 is optional)
runcas ! SA-CASSCF for all three states examples/
set,symmetry=[2,1] ! remove Piy oh_runmrci2.com
runmrci ! MRCI for 2Pix and 2Sigma+

```

This produces

RESULTS

METHOD	STATE	S	ENERGY	DIPX	DIPY	DIPZ
CASSCF	1.2	0.5	-75.41331789	0.0	0.0	0.67158730
CASSCF	1.3	0.5	-75.41331789	0.0	0.0	0.67158730
CASSCF	1.1	0.5	-75.24125256	0.0	0.0	0.69975340
MRCI	1.2	0.5	-75.55518444	0.0	0.0	0.66457191
MRCI+D	1.2	0.5	-75.56014871	0.0	0.0	0.66457191
MRCI+P	1.2	0.5	-75.55853208	0.0	0.0	0.66457191
MRCI	1.1	0.5	-75.39442202	0.0	0.0	0.70484623
MRCI+D	1.1	0.5	-75.40040680	0.0	0.0	0.70484623
MRCI+P	1.1	0.5	-75.39846312	0.0	0.0	0.70484623

You may want to extend the active space to include the 2π orbitals. This can be achieved by setting the variable OCC.

```

***,OH
geometry={O;H,O,1.83} ! Geometry definition
set,symmetry=[2,3,1],spin=1 ! 2Pix, 2Piy, and 2Sigma+ states
occ=[4,2,2] ! 4 sigma and 2 pi occupied
runcas ! SA-CASSCF for all three states examples/
set,symmetry=[2,1] ! remove Piy oh_runmrci3.com
runmrci ! MRCI for 2Pix and 2Sigma+

```

For accurate calculations of the electronic transition moment, also the 1δ orbitals contribute significantly. These are in symmetry 1 ($\delta_{x^2-y^2}$) and 4 (δ_{xy}). In order to force the $5a_1$ orbital to become the δ , we must use the SYM directive in the SCF calculation. Since it is not possible to insert this into the procedure, we must write the SCF input explicitly.

```

***,OH
basis=avqz ! Use aug-cc-pVQZ basis
geometry={O;H,O,1.83} ! Geometry definition
set,symmetry=[2,3,1],spin=1 ! 2Pix, 2Piy, and 2Sigma+ states
occ=[5,2,2,1] ! 4 sigma, 2 pi, 1 delta occupied
hf ! do scf calculation
sym,1, 1,1,1,1,2 ! 5th orbital in symmetry 1 has extra symmetry 2 (=oh_runmrci4.com)
runcas ! SA-CASSCF for all three states
set,symmetry=[2,1] ! remove Piy
runmrci ! MRCI for 2Pix and 2Sigma+

```

Note that this calculation is quite expensive!

3.8.3 MP2 geometry optimization

The following input performs an MP2 geometry optimization for water.

```

***,H2O                                !title
basis=vtz                               !use cc-pVTZ basis
geometry={O;                             !Z-matrix for water
          H1,O,R;
          H2,O,R,H1,THETA}
R=0.96 Ang                               !start bond distance
Theta=104                                !start bond angle
optmp2                                   !do MP2 geometry optimization
show,R,Theta                             !show optimized geometry parameters

```

examples/
h2o_optmp2.com

At the end of the output the following summary of results is printed.

RESULTS FOR BASIS=VTZ

METHOD	STATE	S	ENERGY	DIPX	DIPY	DIPZ
MP2(D)	1.1	0.0	-76.31865774	0.0	0.0	0.76152434
R / ANG		=		0.95906396		
THETA / DEGREE		=		103.51638614		

The next calculation optimizes the geometry at the MP2 level and subsequently performs MP4 and CCSD(T) calculations at the optimized geometry.

```

geometry={O;                             !Z-matrix for water
          H1,O,R;
          H2,O,R,H1,THETA}
basis=vtz                               !use VTZ basis
R=0.96 Ang                               !start bond distance
Theta=104                                !start bond angle
optmp2                                   !optimize energy at mp2 level
runmp4                                   !do single-point MP4 at optimized mp2 geometry
runccsd(t)                               !do single-point ccSD(T) calculation

```

examples/
h2o_optmp2_runccsd(t).com

At the end of the output the following summary of results is printed.

RESULTS FOR BASIS=VTZ

METHOD	STATE	S	ENERGY
MP2(D)	1.1	0.0	-76.31865774
MP2	1.1	0.0	-76.31865774
MP3	1.1	0.0	-76.32273584
MP4(SDQ)	1.1	0.0	-76.32484084
MP4(SDTQ)	1.1	0.0	-76.33305159
MP2	1.1	0.0	-76.31865774
CCSD	1.1	0.0	-76.32454712
CCSD(T)	1.1	0.0	-76.33221602
CCSD[T]	1.1	0.0	-76.33240959

The MP2 energy appears repeatedly, since it is computed in the MP2, MP4, and CCSD(T) calculations.

3.8.4 DFT frequency calculation

The following input performs a DFT/B3LYP geometry optimization and frequency calculation for water:

```

geometry={O;                                !Z-matrix for water
          H1,O,R;
          H2,O,R,H1,THETA}
R=0.96 Ang                                  !start bond distance
Theta=104                                   !start bond angle
basis=6-31g**                               !Pople basis set
functional=b3lyp                             !define functional (optional, b3lyp is default)
freqdft                                      !run frequency calculation

```

examples/
h2o_freqdft.com

The results are

RESULTS FOR BASIS=6-31G**

METHOD	STATE	S	ENERGY	DIPX	DIPY	DIPZ
FREQ[KS/B3LYP]	1.1	0.0	-76.38101813	0.0	0.0	0.80039692
ZPE	1.1	0.0	0.02697203	0.0	0.0	0.80039692
E+ZPE	1.1	0.0	-76.35404610	0.0	0.0	0.80039692
HTOTAL	1.1	0.0	-76.35023712	0.0	0.0	0.80039692
GTOTAL	1.1	0.0	-76.37236272	0.0	0.0	0.80039692

The frequencies, intensities and further details can be found in the output file.

3.9 Do loops

Now you have the idea that one geometry is not enough. Why not compute the whole surface? DO loops make it easy. Here is an example, which computes a whole potential energy surface for H₂O.

```

***,H2O potential
geometry={x;                                !use cs symmetry
          o;                                !z-matrix
          h1,o,r1(i);
          h2,o,r2(i),h1,theta(i) }
basis=vdz                                    !define basis set
angles=[100,104,110]                        !list of angles
distances=[1.6,1.7,1.8,1.9,2.0]            !list of distances
i=0                                          !initialize a counter
do ith=1,#angles                            !loop over all angles H1-O-H2
do ir1=1,#distances                         !loop over distances for O-H1
do ir2=1,ir1                                !loop over O-H2 distances(r1.ge.r2)
i=i+1                                       !increment counter
r1(i)=distances(ir1)                       !save r1 for this geometry
r2(i)=distances(ir2)                       !save r2 for this geometry
theta(i)=angles(ith)                       !save theta for this geometry
hf;                                         !do SCF calculation
escf(i)=energy                              !save scf energy for this geometry
ccsd(t);                                    !do CCSD(T) calculation
eccsd(i)=energc                             !save CCSD energy
eccsd(t)=energy                            !save CCSD(T) energy
enddo                                       !end of do loop ith
enddo                                       !end of do loop ir1
enddo                                       !end of do loop ir2
table,r1,r2,theta,escf,eccsd,eccsd(t)       !produce a table with results
head, r1,r2,theta,scf,ccsd,ccsd(t)         !modify column headers for table
save,h2o.tab                                !save the table in file h2o.tab
title,Results for H2O, basis $basis        !title for table
sort,3,1,2                                  !sort table

```

examples/
h2o_pes_ccsd.com

This produces the following table.

Results for H2O, basis VDZ

R1	R2	THETA	SCF	CCSD	CCSD(T)
1.6	1.6	100.0	-75.99757338	-76.20140563	-76.20403920
1.7	1.6	100.0	-76.00908379	-76.21474489	-76.21747582
1.7	1.7	100.0	-76.02060127	-76.22812261	-76.23095473
...					
2.0	1.9	110.0	-76.01128923	-76.22745359	-76.23081968
2.0	2.0	110.0	-76.00369171	-76.22185092	-76.22537212

You can use also use DO loops to repeat your input for different methods.

```

***,h2o benchmark
method=[hf, fci, ci, cepa(0), cepa(1), cepa(2), cepa(3), mp2, mp3, mp4, \
      qci, ccscf, bccsd, qci(t), ccscf(t), bccsd(t), casscf, mrci, acpf]
basis=dz                                !Double zeta basis set
geometry={o;h1,o,r;h2,o,r,h1,theta}     !Z-matrix for geometry
r=1 ang, theta=104                       !Geometry parameters
do i=1,#method                            !Loop over all requested methods
$method(i);                               !call program
e(i)=energy                               !save energy for this method
enddo
escf=e(1)                                 !scf energy
efci=e(2)                                 !fci energy
table,method,e,e-escf,e-efci             !print a table with results
!Title for table:
title,Results for H2O, basis $basis, R=$r Ang, Theta=$theta degree

```

examples/
h2o_manymethods.com

This calculation produces the following table.

Results for H2O, basis DZ, R=1 Ang, Theta=104 degree

METHOD	E	E-ESCF	E-EFCI
HF	-75.99897339	.00000000	.13712077
FCI	-76.13609416	-.13712077	.00000000
CI	-76.12844693	-.12947355	.00764722
CEPA(0)	-76.13490643	-.13593304	.00118773
CEPA(1)	-76.13304720	-.13407381	.00304696
CEPA(2)	-76.13431548	-.13534209	.00177868
CEPA(3)	-76.13179688	-.13282349	.00429728
MP2	-76.12767140	-.12869801	.00842276
MP3	-76.12839400	-.12942062	.00770015
MP4	-76.13487266	-.13589927	.00122149
QCI	-76.13461684	-.13564345	.00147732
CCSD	-76.13431854	-.13534515	.00177561
BCCD	-76.13410586	-.13513247	.00198830
QCI(T)	-76.13555640	-.13658301	.00053776
CCSD(T)	-76.13546225	-.13648886	.00063191
BCCD(T)	-76.13546100	-.13648762	.00063315
CASSCF	-76.05876129	-.05978790	.07733286
MRCI	-76.13311835	-.13414496	.00297580
ACPF	-76.13463018	-.13565679	.00146398

One can do even more fancy things, like, for instance, using macros, stored as string variables. See example `oh_macros.com` for a demonstration.

4 PROGRAM CONTROL

4.1 Starting a job (***)

The first card of each input should be:

```
***,text
```

where *text* is arbitrary. If file 1 is restarted, *text* must always be the same. The effect of this card is to reset all program counters, etc. If the *** card is omitted, *text* assumes its default value, which is all blank.

4.2 Ending a job (---)

The end of the input is signalled by either an end of file, or a

```
---
```

card. All input following the --- card is ignored.

Alternatively, a job can be stopped at some place by inserting an EXIT card. This could also be in the middle of a DO loop or an IF block. If in such a case the --- card would be used, an error would result, since the ENDDO or ENDDIF cards would not be found.

4.3 Restarting a job (RESTART)

In contrast to MOLPRO92 and older versions, the current version of MOLPRO attempts to recover all information from all permanent files by default. If a restart is unwanted, the NEW option can be used on the FILE directive. The RESTART directive as described below can still be used as in MOLPRO92, but is usually not needed.

```
RESTART,r1,r2,r3,r4,...;
```

The r_i specify which files are restarted. These files must have been allocated before using FILE cards. There are two possible formats for the r_i :

- a) $0 < r_i < 10$: Restart file r_i and restore all information.
- b) $r_i = \text{name.nr}$: Restart file nr but truncate before record *name*.

If all $r_i = 0$, then all permanent files are restarted. However, if at least one r_i is not equal to zero, only the specified files are restarted.

Examples:

```
RESTART;           will restart all permanent files allocated with FILE cards (default)
RESTART, 1;       will restart file 1 only
RESTART, 2;       will restart file 2 only
RESTART, 1, 2, 3; will restart files 1-3
RESTART, 2000.1;  will restart file 1 and truncate before record 2000.
```


4.4 Including secondary input files (INCLUDE)

```
INCLUDE,file,echo;
```

Insert the contents of the specified *file* in the input stream. In most implementations the file name given is used directly in a Fortran open statement. If the parameter *echo* is nonzero, the included file is echoed to the output in the normal way, but by default its contents are not printed. The included file may itself contain INCLUDE commands up to a maximum nesting depth of 10.

4.5 Allocating dynamic memory (MEMORY)

```
MEMORY,n,scale;
```

Sets the limit on dynamic memory to *n* floating point words. If *scale* is given as K, *n* is multiplied by 1000; if *scale* is M, *n* is multiplied by 1 000 000.

Note: The MEMORY card must precede all FILE cards!

Examples:

```
MEMORY,90000      allocates 90 000 words of memory
MEMORY,500,K      allocates 500 000 words of memory
MEMORY,2,M        allocates 2 000 000 words of memory
```

4.6 DO loops (DO/ENDDO)

DO loops can be constructed using the DO and ENDDO commands. The general format of the DO command is similar to Fortran:

```
DO variable=start, end [[,]increment] [[,]unit]
```

where *start*, *end*, *increment* may be expressions or variables. The default for *increment* is 1. In contrast to Fortran, these variables can be modified within the loop (to be used with care!). For instance:

```
DR=0.2
DO R=1.0,6.0,DR,ANG
IF (R.EQ.2) DR=0.5
IF (R.EQ.3) DR=1.0
....
ENDDO
```

performs the loop for the following values of R: 1.0, 1.2, 1.4, 1.6, 1.8, 2.0, 2.5, 3.0, 4.0, 5.0, 6.0 Ångström. The same could be achieved as follows:

```
RVEC=[1.0,1.2,1.4,1.6,1.8,2.0,2.5,3.0,4.0,5.0,6.0] ANG
DO I=1,#RVEC
R=RVEC(I)
....
ENDDO
```

Up to 20 DO loops may be nested. Each DO must end with its own ENDDO.

Jumps into DO loops are possible if the DO variables are known. This can be useful in restarts, since it allows to continue an interrupted calculation without changing the input (all variables are recovered in a restart).

4.6.1 Examples for do loops

The first example shows how to compute a potential energy surface for water.

```

***,H2O potential
geometry={x;                                !use cs symmetry
         o;                                !z-matrix
         h1,o,r1(i);
         h2,o,r2(i),h1,theta(i) }
basis=vdz                                  !define basis set
angles=[100,104,110]                       !list of angles
distances=[1.6,1.7,1.8,1.9,2.0]           !list of distances
i=0                                         !initialize a counter
do ith=1,#angles                           !loop over all angles H1-O-H2
do ir1=1,#distances                        !loop over distances for O-H1
do ir2=1,ir1                                !loop over O-H2 distances(r1.ge.r2)
i=i+1                                       !increment counter
r1(i)=distances(ir1)                       !save r1 for this geometry
r2(i)=distances(ir2)                       !save r2 for this geometry
theta(i)=angles(ith)                       !save theta for this geometry
hf;                                         !do SCF calculation
escf(i)=energy                             !save scf energy for this geometry
ccsd(t);                                   !do CCSD(T) calculation
eccsd(i)=energc                             !save CCSD energy
eccsd(t)=energy                             !save CCSD(T) energy
enddo                                       !end of do loop ith
enddo                                       !end of do loop ir1
enddo                                       !end of do loop ir2
table,r1,r2,theta,escf,eccsd,eccsd(t)      !produce a table with results
head, r1,r2,theta,scf,ccsd,ccsd(t)         !modify column headers for table
save,h2o.tab                               !save the table in file h2o.tab
title,Results for H2O, basis $basis        !title for table
sort,3,1,2                                 !sort table

```

examples/
h2o_pes_ccsd.com

The next example shows how to loop over many methods.

```

***,h2o benchmark
method=[hf, fci, ci, cepa(0), cepa(1), cepa(2), cepa(3), mp2, mp3, mp4, \
       qci, ccsd, bccd, qci(t), ccsd(t), bccd(t), casscf, mrci, acpf]
basis=dz                                   !Double zeta basis set
geometry={o;h1,o,r;h2,o,r,h1,theta}       !Z-matrix for geometry
r=1 ang, theta=104                         !Geometry parameters
do i=1,#method                             !Loop over all requested methods
$method(i);                               !call program
e(i)=energy                                 !save energy for this method
enddo
escf=e(1)                                  !scf energy
efci=e(2)                                  !fci energy
table,method,e,e-escf,e-efci               !print a table with results
!Title for table:
title,Results for H2O, basis $basis, R=$r Ang, Theta=$theta degree

```

examples/
h2o_manymethods.com

4.7 Branching (IF/ELSEIF/ENDIF)

IF blocks and IF/ELSEIF blocks can be constructed as in FORTRAN.

4.7.1 IF statements

IF blocks have the same form as in Fortran:

```
IF (logical expression) THEN
statements
ENDIF
```

If only one statement is needed, the one-line form

```
IF (logical expression) statement
```

can be used, except if *statement* is a procedure name.

ELSE and ELSE IF can be used exactly as in Fortran. IF statements may be arbitrarily nested. Jumps into IF or ELSE IF blocks are allowed. In this case no testing is performed; when an ELSE is reached, control continues after ENDIF.

The logical expression may involve logical comparisons of algebraic expressions or of strings. Examples:

```
IF(STATUS.LT.0) THEN
TEXT,An error occurred, calculation stopped
STOP
ENDIF
```

```
IF($method.eq.'HF') then
...
ENDIF
```

In the previous example the dollar and the quotes are optional:

```
IF(METHOD.EQ.HF) then
...
ENDIF
```

4.7.2 GOTO commands

GOTO commands can be used to skip over parts of the input. The general form is

```
GOTO , command , [ n ] , [ nrep ]
```

Program control skips to the $|n|$ 'th occurrence of *command* (Default: $n = 1$). *command* must be a keyword in the first field of an input line. If *n* is positive, the search is forward starting from the current position. If *n* is negative, search starts from the top of the input. The GOTO command is executed at most *nrep* times. The default for *nrep* is 1 if $n < 0$ and infinity otherwise. We recommend that GOTO commands are never used to construct loops.

Alternatively, one can jump to labels using

```
GOTO , label
```

Since labels must be unique, the search starts always from the top of the input. It is required that the *label* ends with a colon.

4.7.3 Labels (LABEL)

```
LABEL
```

This is a dummy command, sometimes useful in conjunction with GOTO.

4.8 Procedures (PROC/ENDPROC)

Procedures can be defined at the top of the input, in the default file `molproi.rc`, or in INCLUDE files as follows:

```
PROC name
statements
ENDPROC
```

Alternatively, one can use the form

```
PROC name [=] {statements}
```

In the latter case, it is required that the left curly bracket (`{`) appears on the same line as `PROC`, but *statements* can consist of several lines. If in the subsequent input *name* is found as a command in the first field of a line, it is substituted by the *statements*. Example:

```
PROC SCF
IF (#SPIN.EQ.0.OR.MOD(SPIN,2).NE.MOD(NELEC,2)) SET,SPIN=MOD(NELEC,2)
IF (SPIN.EQ.0) THEN
  HF
ELSE
  RHF
ENDIF
ENDPROC
```

Alternatively, this could be written as

```
PROC SCF={
IF (#SPIN.EQ.0.OR.MOD(SPIN,2).NE.MOD(NELEC,2)) SET,SPIN=MOD(NELEC,2)
IF (SPIN.EQ.0) THEN; HF; ELSE; RHF; ENDIF}
```

Procedures may be nested up to a depth of 10. In the following example `SCF` is a procedure:

```
PROC CC
SCF
IF (SPIN.EQ.0) THEN
  CCSD
ELSE
  RCCSD
ENDPROC
```

Note: Procedure names are substituted only if found in the first field of an input line. Therefore, they must not be used on one-line `IF` statements; please use `IF / ENDIF` structures instead.

If as first statement of a procedure `ECHO` is specified, the substituted commands of the present and lower level procedures will be printed. If `ECHO` is specified in the main input file, all subsequent procedures are printed.

Certain important input data can be passed to the program using variables. For instance, occupancy patterns, symmetries, number of electrons, and multiplicity can be defined in this way (see section 6.7 for more details). This allows the quite general use of procedures. For example, assume the following procedure has been defined in `molproi.rc`:

```
PROC MRCI
IF (INTDONE.EQ.0) INT
IF (SCFDONE.EQ.0) THEN
SCF
```

```
ENDIF
MULTI
CI
ENDPROC
```

This procedure can be used for a calculation of a vertical ionization potential of H₂O as follows:

```
R=1 ANG          !Set bond distance
THETA=104 DEGREE !Set bond angle

BASIS=VTZ        !Define basis set

GEOMETRY         !Geometry input block
O                !Z-matrix
H1,O,R
H2,O,R,H1,THETA
ENDG             !End of geometry input

MRCI             !Compute mrci energy of water using defaults
EH2O=ENERGY      !save mrci energy in variable EH2O

SET,NELEC=9      !Set number of electrons to 9
SET,SYMMETRY=2   !Set wavefunction symmetry to 2
MRCI             !Compute mrci energy of H2O+ (2B2 state)

IPCI=(ENERGY-EH2O)*TOEV  compute MRCI ionization potential in eV
```

Note: At present, all variables are *global*, i.e., variables are commonly known to all procedures and all variables defined in procedures will be subsequently known outside the procedures as well. The reason is that procedures are included into the internal input deck at the beginning of the job and not at execution time; for the same reason, variable substitution of procedure names is not possible, e.g. one cannot use constructs like

```
method=scf
$method      !this does not work!
```

4.9 Text cards (TEXT)

```
TEXT,xxxxxx
```

will just print *xxxxxx* in the output. If the text contains variables which are preceded by a dollar (\$), these are replaced by their actual values, e.g.

```
r=2.1
text,Results for R=\$r
```

will print

```
Results for R=2.1
```

4.10 Checking the program status (STATUS)

```
STATUS,nr,prog1,...,prog6
```

This command reads the status information from file *nr* and checks the status of the specified program steps. *prog1* to *prog6* may be HF, RHF, UHF, MCSCF, CI, MULTI,

FORCES. If none of these is specified, the status of the last step is checked. If one of *prog*₁ to *prog*₆ is CRASH or STOP the program will either crash or stop if status was not o.k. (STOP is default). If CLEAR is specified, a bad status is cleared, so there will be no crash at subsequent status checks.

Examples:

STATUS, 1, HF, CRASH; will check the status of the last HF-SCF step and crash if it was not o.k. (i.e. no convergence). CRASH is useful to avoid that the next program in a chain is executed.

STATUS, 2, MULTI, CI, STOP; will check the status of the most previous MULTI and CI steps which had allocated file 2 and stop if something was wrong.

STATUS, 1, RHF, CLEAR; will clear status flag for last RHF. No action even if RHF did not converge.

Note that the status variables are recovered in a restart.

4.11 Global Thresholds (GTHRESH)

A number of global thresholds can be set using the GTHRESH command outside the individual programs (the first letter G is optional, but should be used to avoid confusion with program specific THRESH cards). The syntax is

GTHRESH, *key1=value1, key2=value2, . . .*

key can be one of the following.

ZERO	Numerical zero (default 1.d-12)
ONEINT	Threshold for one-electron integrals (default 1.d-12, but not used at present)
TWOINT	Threshold for the neglect of two-electron integrals (default 1.d-12)
PREFAC	Threshold for test of prefactor in TWOINT (default 1.d-14)
LOCALI	Threshold for orbital localization (default 1.d-8)
EORDER	Threshold for reordering of orbital after localization (default 1.d-4)
ENERGY	Convergence threshold for energy (default 1.d-6)
GRADIENT	Convergence threshold for orbital gradient in MCSCF (default 1.d-2)
STEP	Convergence threshold for step length in MCSCF orbital optimization (default 1.d-3)
ORBITAL	Convergence threshold for orbital optimization in the SCF program (default 1.d-5).
CIVEC	Convergence threshold for CI coefficients in MCSCF and reference vector in CI (default 1.-d.5)
COEFF	Convergence threshold for coefficients in CI and CCSD (default 1.d-4)
PRINTCI	Threshold for printing CI coefficients (default 0.05)
PUNCHCI	Threshold for punching CI coefficients (default 99 - no punch)

4.12 Global Print Options (GPRINT/NOGPRINT)

Global print options can be set using the GPRINT command outside the individual programs (the first letter G is optional, but should be used to avoid confusion with program specific PRINT cards). The syntax is

```
GPRINT,key1[=value1],key2[=value2],...
NOGPRINT,key1,key2,...
```

Normally, *value* can be omitted, but values > 0 may be used for debugging purposes, giving more information in some cases. The default is no print for all options, except for DISTANCE, ANGLES (default=0), and VARIABLE. NOGPRINT,*key* is equivalent to PRINT,*key*=-1. *key* can be one of the following:

BASIS	Print basis information
DISTANCE	Print bond distances (default)
ANGLES	Print bond angle information (default). If > 0, dihedral angles are also printed.
ORBITAL	Print orbitals in SCF and MCSCF
CIVECTOR	Print CI vector in MCSCF
PAIRS	Print pair list in CI, CCSD
CS	Print information for singles in CI, CCSD
CP	Print information for pairs in CI, CCSD
REF	Print reference CSFs and their coefficients in CI
PSPACE	Print p-space configurations
MICRO	Print microiterations in MCSCF and CI
CPU	Print detailed CPU information
IO	Print detailed I/O information
VARIABLE	Print variables each time they are set or changed (default).

4.13 One-electron operators and expectation values (GEXPEC)

The operators for which expectation values are requested, are specified by keywords on the global GEXPEC directive. The first letter G is optional, but should be used to avoid confusion with program specific EXPEC cards, which have the same form as GEXPEC. For all operators specified on the GEXPEC card, expectation values are computed in all subsequent programs (if applicable).

For a number of operators it is possible to use *generic* operator names, e.g., DM for dipole moments, which means that all three components DMX, DMY, and DMZ are computed. Alternatively, individual components may be requested.

The general format is as follows:

```
[ G ] EXPEC,opname[,][icen,[x,y,z]],...
```

where

opname operator name (string), either generic or component.

icen z-matrix row number or z-matrix symbol used to determine the origin (x,y,z must not be specified).
If *icen*= 0 or blank, the origin must be specified in x,y,z

Several GEXPEC cards may follow each other, or several operators may be specified on one card.

Examples:

GEXPEC ,QM computes quadrupole moments with origin at (0,0,0),

GEXPEC ,QM1 computes quadrupole moments with origin at centre 1.

GEXPEC ,QM ,O1 computes quadrupole moments with origin at atom O1.

GEXPEC ,QM,,1 , 2 , 3 computes quadrupole moments with origin at (1,2,3).

The following table summarizes all available operators:

Expectation values are only nonzero for symmetric operators (parity=1). Other operators can be used to compute transition quantities (spin-orbit operators need a special treatment). By default, the dipole moments are computed.

4.13.1 Example for computing expectation values

The following job computes dipole and quadrupole moments for H₂O.

```

***,h2o properties
geometry={o;h1,o,r;h2,o,r,h1,theta} !Z-matrix geometry input
r=1 ang !bond length
theta=104 !bond angle
gexpec,dm,sm,qm !compute dipole and quarupole moments
methods=[hf,multi,ci] !do hf, casscf, mrci
do i=1,#methods !loop over methods
$methods(i) !run energy calculation
e(i)=energy
dip(i)=dmz !save dipole moment in variable dip
quadxx(i)=qmxx !save quadrupole moments
quady(i)=qmyy
quadzz(i)=qmzz
smxx(i)=xx !save second moments
smyy(i)=yy
smzz(i)=zz
enddo
table,methods,dip,smxx,smyy,smzz !print table of first and second moments
table,methods,e,quadxx,quady,quadzz !print table of quadrupole moments

```

examples/
h2o_gexpec2.com

This Job produces the following tables

METHODS	DIP	SMXX	SMYY	SMZZ
HF	0.82747571	-5.30079792	-3.01408114	-4.20611391
MULTI	0.76285513	-5.29145148	-3.11711397	-4.25941000
CI	0.76868508	-5.32191822	-3.15540500	-4.28542917

METHODS	E	QUADXX	QUADYY	QUADZZ
HF	-76.02145798	-1.69070039	1.73937477	-0.04867438
MULTI	-76.07843443	-1.60318949	1.65831677	-0.05512728
CI	-76.23369821	-1.60150114	1.64826869	-0.04676756

4.13.2 Example for computing relativistic corrections

```

***,ar2
geometry={ar1;ar2,ar1,r} !geometry definition
r=2.5 ang !bond distance
hf; !non-relativistic scf calculation
expec,rel,darwin,massv !compute relativistic correction using Cowan-Griffin operator
e_nrel=energy !save non-relativistic energy in variable enrel
show,massv,darwin,erel !show individual contribution and their sum

dkroll=1 !use douglas-kroll one-electron integrals examples/
hf; !relativistic scf calculation ar2_rel.com
e_dk=energy !save relativistic scf energy in variable e_dk.
show,massv,darwin,erel !show mass-velocity and darwin contributions and their sum
show,e_dk-e_nrel !show relativistic correction using Douglas-Kroll

```

This jobs shows at the end the following variables:

```

MASSV / AU = -14.84964285
DARWIN / AU = 11.25455679
EREL / AU = -3.59508606

```

Table 5: One-electron operators and their components

Generic name	Parity	Components	Description
OV	1		Overlap
EKIN	1		Kinetic energy
POT	1		potential energy
DELTA	1		delta function
DEL4	1		Δ^4
DARW	1		one-electron Darwin term, i.e., DELT with appropriate factors summed over atoms.
MASSV	1		mass-velocity term, i.e., DEL4 with appropriate factor.
REL	1		total Cowan-Griffin Relativistic correction, i.e., DARW+MASSV.
DM	1	DMX, DMY, DMZ	dipole moments
SM	1	XX, YY, ZZ, XY, XZ, YZ	second moments
TM	1	XXX, XXY, XXZ, XYY, XYZ, XZZ, YYY, YYZ, YZZ, ZZZ	third moments
MLTP n	1	all unique cartesian products of order n	multipole moments
QM	1	QMXX, QMYY, QMZZ, QMXY, QMXZ, QMYZ, QMR=XX + YY + ZZ, QMXX=(3 XX - RR)/2, QMXY=3 XY / 2 etc.	quadrupole moments and R^2
EF	1	EFX, EFY, EFZ	electric field
FG	1	FGXX, FGYY, FGZZ, FGXY, FGXZ, FGYZ	electric field gradients
DMS	1	DMSXX, DMSYX, DMSZX, DMSXY, DMSYY, DMSZY, DMSXZ, DMSYZ, DMSZZ	diamagnetic shielding tensor
LOP	-1	LX, LY, LZ	Angular momentum operators $\hat{L}_x, \hat{L}_y, \hat{L}_z$
LOP2	1	LXLX, LYLY, LZLZ, LXLY, LXLZ, LYLZ The symmetric combinations	one electron parts of products of angular momentum operators. $\frac{1}{2}(\hat{L}_x\hat{L}_y + \hat{L}_y\hat{L}_x)$ etc. are computed
VELO	-1	D/DX, D/DY, D/DZ	velocity
LS	-1	LSX, LSY, LSZ	spin-orbit operators
ECPLS	-1	ECPLSX, ECPLSY, ECPLSZ	ECP spin-orbit operators

5 FILE HANDLING

5.1 FILE

The FILE directive is used to open permanent files, which can be used for later restarts. The syntax in MOLPRO94 and later versions is

FILE,*file,name,[status]*

file is the logical MOLPRO file number (1-9). *name* is the file name (will be converted to lower case). *status* can be one of the following:

UNKNOWN	A permanent file is opened. If it exists, it is automatically restarted. This is the default.
OLD	Same effect as UNKNOWN. No error occurs if the file does not exist.
NEW	A permanent file is opened. If it already exists, it is erased and not restarted.
ERASE	Same effect as NEW.
SCRATCH	A temporary file is opened. If it already exists, it is erased and not restarted. After the job has finished, the file is no longer existent.
DELETE	Same effect as SCRATCH.

Note that RESTART is now the default for all permanent files. All temporary files are usually allocated automatically where needed. I/O buffers are allocated at the top of the dynamic memory, and the available memory decreases by the size of the buffers. The MEMORY card must therefore be presented before the first FILE card!

Examples:

FILE , 1 , H2O . INT allocates permanent file 1 with name H2O . INT. Previous information on the file is recovered.

FILE , 2 , H2O . WFU , NEW allocates permanent file 2 with name H2O . WFU. All previous information on the file is erased.

Note that filenames are converted to lower case on unix machines.

5.2 DELETE

DELETE,*file1, file2, ...*

Deletes the specified files. *file* refers to the logical MOLPRO file numbers as specified on the FILE card.

5.3 ERASE

ERASE,*file1, file2, ...*

Erases the specified files. *file* refers to the logical MOLPRO file numbers as specified on the FILE card.

5.4 DATA

The DATA command can be used to modify the MOLPRO binary files.

UNIT	Alias for NPL (should never be used)
RENAME, <i>rec1,rec2</i>	used to rename <i>rec1</i> to <i>rec2</i> . <i>rec1</i> and <i>rec2</i> must be given in the form <i>name.ifil</i> , where <i>ifil</i> is the number of a MOLPRO binary file (alias for NAME).
TRUNCATE , <i>nen</i>	used to truncate files after <i>nen-1</i> records (alias for NEN).
TRUNCATE , <i>rec</i>	used to truncate before record <i>rec</i> . <i>rec</i> must be given in the form <i>name.ifil</i> , where <i>ifil</i> is the number of a MOLPRO binary file.
COUNT	Alias for NRE (presently not used)
COPY, <i>rec1,rec2</i>	Copies record <i>rec1</i> to <i>rec2</i> . <i>rec1</i> and <i>rec2</i> must be given in the form <i>nam1.ifil1, nam2.ifil2</i> . If <i>nam2=0</i> , <i>nam2=nam1</i> . If <i>nam1=0</i> , all records are copied from file <i>ifil1</i> to file <i>ifil2</i> .

5.5 Assigning punch files (PUNCH)

PUNCH,*filename*, [REWIND]

Opens punch file named *filename*. If this file already exists, it is appended, unless the REWIND or NEW option is specified; in that case, any previous information on the punch file is overwritten. See FILE for machine dependent interpretation of filename. The punch file contains all important results (geometries, energies, dipole, transition moments etc). It can be read by a separate program READPUN, which can produce tables in user supplied format.

Example:

```
PUNCH , H2O . PUN  allocates punch file H2O . PUN
```

Note that the file name is converted to lower case on unix machines.

5.6 MOLPRO system parameters (GPARAM)

The GPARAM card allows to change MOLPRO system parameters. This should only be used by experts!

GPARAM,*option=value*,...

The following options can be given in any order.

NOBUFF	if present, disable system buffering
LSEG	disk sector length
INTREL	number of integer words per real word (should never be modified!)
IBANK	number of memory banks. Default is 2, which should always be o.k.
IVECT	0=scalar, 1=vector machine
MINVEC	minimum vector length for call to mxmb
LTRACK	page size in buffer routines (must be multiple of <i>lseg</i>)

LENBUF	length of integral buffer (file 1)
NTR	length of integral records (must be multiple of 3· <i>ltrack</i>)
LTR	disk sector length assumed in CI (default 1 is reasonable)
NCACHE	machine cache size in bytes
IASYN	if nonzero, use asynchronous I/O on CONVEX
MXMBLK	column/row block size for mxma
MXMBLN	link block size for mxma
NCPUS	maximum number of cpus to be used in multitasking
MINBR1	min number of floating point ops per processor
MXDMP	highest file number to be treated as dump file with full functionality ($1 \leq \text{MXDMP} \leq .3$).

The MXDMP option is for experts only! This prevents basis and geometry information from being written to dump files with higher file number than the given *value*, and can sometimes be useful for counterpoise corrected geometry optimizations. Note that some functionality is lost by giving this option, and errors will result unless all input is correct!

6 VARIABLES

Data may be stored in *variables*. A variable can be of type *string*, *real* or *logical*, depending on the type of the expression in its definition. Any sequence of characters which is not recognized as expression or variable is treated as string. In this section, we will discuss only *real* and *logical* variables. *String* variables will be discussed in more detail in section 6.2. Variables can be used anywhere in the input, but they can be set only outside the input blocks for specific programs. For example, if a variable is used within the input block for HF, it must have been set before the HF{ . . . } input block.

MOLPRO automatically stores various results and data in system variables (see section 6.7.1), which can be used for further processing. A new feature of MOLPRO2002 is that most system variables are write protected and cannot be overwritten by the user. The input is automatically checked before the job starts, and should a system variable be set in the input the job will stop immediately with an error message. Only in some exceptions (see section 6.3), system variables can be modified using the SET command (but not with the simple NAME=*value* syntax). Note that due to the changed usage and syntax of the SET command, compatibility with MOLPRO92 input syntax is no longer maintained.

6.1 Setting variables

A variable can be defined using

```
variable = expression [,] [unit], ...
```

unit is an optional string, which can be used to associate a unit to a variable. A variable definition is recognized by the equals sign in the *first* field of the input card. For instance,

```
R=1 ANG, THETA=100 DEGREE
```

defines the variables R and THETA, but

```
THRESH, ENERGY=1.d-8, GRADIENT=1.d-5
```

does *not* define variables; here ENERGY and GRADIENT are keywords to be recognized by the program.

6.2 String variables

String variables can be set as other variables in the form

variable = *string*

If *string* contains blanks or other special characters (like +, -, /), it must be given in quotes. Instead of *string*, also another string variable can be used, e.g.,

```
METHOD=PROGRAM
```

where PROGRAM is a string variable set by the program (see section *special variables*). The same name must not be used for a *string* variable and a *real* or *logical* variable.

As a general rule, string variables are replaced by their value only if they are preceded by a dollar (\$) (exceptions: in variable definitions, on SHOW cards, and in logical expressions on IF cards, the dollar is optional). This is a precaution to avoid commands which have the same name as a variable being interpreted as variables. Variables may also appear on TEXT or TITLE cards or in strings, but must be preceded by \$ in these cases. Example:

```
METHOD=MCSCF
R=1.5
TEXT,$method results for R=$R Bohr
```

prints

```
MCSCF results for R=1.5 Bohr
```

String variables can be concatenated with strings or other string variables in the following way. Assume that variable PROGRAM has the value MRCI. Setting

```
METHOD= '$PROGRAM+Q'
```

sets METHOD to MRCI+Q. Alternatively, if we would also have a variable VERSION with value Q, we could write

```
METHOD= '$PROGRAM+$VERSION'
```

Again, the value of METHOD would be MRCI+Q. Note that the quotes are necessary in these cases.

6.3 System variables

As mentioned above, most system variables cannot be written by the user. In some exceptions, it is possible to redefine them using the SET command:

```
SET,variable = expression [,] [unit]
```

This holds for the following variables:

CHARGE	Total charge of the molecule
NELEC	Number of electrons
SPIN	Spin quantum number, given as $2 \cdot M_S$ (integer)

SCFSPIN	Same as SPIN, but only for HF
MCSPIN	Same as SPIN, but only for MCSCF
CISPIN	Same as SPIN, but only for MRCI
STATE	State to be optimized
MCSTATE	Same as STATE but only for MCSCF
CISTATE	Same as STATE but only for MRCI
SYMMETRY	State symmetry
SCFSYM[METRY]	Same as SYMMETRY but only for HF
MCSYM[METRY]	Same as SYMMETRY but only for MCSCF
CISYM[METRY]	Same as SYMMETRY but only for MRCI
ZSYMEL	Symmetry elements
LQUANT	Lambda quantum number for linear molecules
OPTCONV	Geometry optimization convergence criterion
PROGRAM	Last program name
CPUSTEP	CPU-time of last program step
SYSSTEP	System-time of last program step
WALLSTEP	Elapsed-time of last program step
FOCKDONE	Indicates if closed-shell fock operator is available.

6.4 Macro definitions using string variables

String variables for which the stored string has the form of an algebraic expression are evaluated to a number if they are preceded by two dollars (\$\$). Example:

```
string='a+b'
a=3
b=4
text,This is string $string which evaluates to $$string
```

prints

```
** This is string a+b which evaluates to 7
```

This can be used to define simple macros, which can be used at various places in the subsequent input. For instance,

```
ECORR='ENERGY-ESCF' !define a macro
HF !do SCF calculation
ESCF=ENERGY !store SCF energy in variable ESCF
MULTI !do CASSCF
DEMC=$$ECORR !store CASSCF correlation energy in variable DEMC
MRCI !do MRCI
DECI=$$ECORR !store MRCI correlation energy in variable DECI
```

Here is an example of advanced use of macros and string variables:

```

***,test for parser
text,This fancy input demonstrates how string variables and macros can be used
text
basis=vdz                !define basis set
geometry={O;H,O,r}      !define geometry (z-matrix)

text,methods
method=[rhf,2[casscf,2[mrci]]]
text,active spaces
spaces=[ '[3,1,1]',3['[4,2,2]'],3['[5,2,2]'] ]
text,symmetries
symset=[ '1',2['[1,2,3]'],'1','2' ]
text,weight factors for state averaged casscf
weights=[ '1','[1,1,1]',2[' '],'[1,0.5,0.5]',2[' ' ] ]
text,scf occupation
scfocc=[3,2[1]]
text,bond distance
r=1.85

hf
do i=1,#method          !loop over methods
occ=$$spaces(i)        !set active space for this run
set,symmetry=$$symset(i) !set symmetries for this run
weight=$$weights(i)    !set weights for this run
$method(i)             !now run method
e(i)='$energy'          !save energies in strings
dipol(i)='$dmz'        !save dipole moments in strings
enddo
table,method,spaces,symset,weights,e,dipol
title,Results for OH, r=$r, basis=$basis
head,method,spaces,symmetries,weights,energies,'dipole moments'
exit

```

examples/
oh_macros.com

6.5 Indexed Variables (Vectors)

Variables may be indexed, but only one-dimensional arrays (vectors) are supported. The index may itself be a variable. For instance

```
METHOD(I)=PROGRAM
E(I)=ENERGY
```

are valid variable definitions, provided I, PROGRAM, and ENERGY are also defined variables. Indices may be nested to any depth.

Different elements of an array can be of different type (either *real* or *logical*). However, only one *unit* can be assigned to an array. String variables have no associated value and cannot be mixed with the other variable types. Therefore, a given variable name can only be used either for a *string* variable or a *real (logical)* variable.

Vectors (arrays) can be conveniently defined using square brackets:

```
R=[1.0,1.2,1.3] ANG
```

This defines an array with three elements, which can be accessed using indices; for instance, R(2) has the value 1.2 ANG. A repeat specifier can be given in front of the left bracket: 5[0] is equivalent to [0,0,0,0,0]. Brackets can even be nested: for instance, 2[1,2,2[2.1,3.1]] is equivalent to [1,2,2.1,3.1,2.1,3.1,1,2,2.1,3.1,2.1,3.1].

Arrays can be appended from a given position just by entering additional elements; for instance,

```
R(4)=[1.4,1.5] ANG
```


or

```
R(4:)= [ 1.4 , 1.5 ] ANG
```

extends the above array to length 5. Previously defined values can be overwritten. For instance

```
R(2)= [ 1.25 , 1.35 , 1.45 ]
```

modifies the above vector to (1.0, 1.25, 1.35, 1.45, 1.5).

If no index is given on the left hand side of the equal sign, an existing variable of the same name is replaced by the new values, and all old values are lost. For instance

```
THETA= [ 100 , 110 , 120 , 130 ] set four values
```

...

```
THETA(1)=104      replace THETA(1) by a new value; THETA(2:4) are unchanged
```

...

```
THETA= [ 140 , 150 ] old variable THETA is replaced; THETA(3:4) are deleted
```

Square brackets can also be used to define an array of strings, e.g.,

```
METHOD= [ INT , HF , CASSCF , MRCI ]
```

These could be used as follows:

```
DO I=1,4
$METHOD(I)
ENDDO
```

The above input would be equivalent to

```
INT
HF
CASSCF
MRCI
```

The current length of an array can be accessed by preceding # to the variable name. For instance, in the above examples #R and #METHOD have the values 5 and 4, respectively. If a variable is not defined, zero is returned but no error occurs. This can be used to test for the existence of a variable, for example:

```
IF( #SPIN.EQ.0.AND.#NELEC.EQ.1) SET,SPIN=MOD(NELEC,2)
```

This defines variable SPIN if it is unknown and if NELEC is a scalar (one dimensional) variable.

6.6 Vector operations

The following simple vector operations are possible:

- Copying or appending a vector to another vector. For instance S=R copies a vector R to a vector S. S(3)=R copies R to S(3), S(4), . . . S(#S+1)=R appends vector R to vector S. It is also possible to access a range of subsequent elements in a vector: S=R(2:4) copies elements 2 to 4 of R to S(1), S(2), S(3). Note that R(2:) denotes elements R(2) to R(#R), but R(2) denotes a single element of R.

- Vector-scalar operations: $R=R*2$ multiplies each element of R by 2. Instead of the number 2, also scalar (one dimensional) variables or expressions can be used, e.g., $R=R*ANG$ converts all elements of R from Ångström to bohr, or $Z=R*COS(THETA)$ creates a vector Z with elements $Z(i) = R(i)*COS(THETA)$. All other algebraic operators can be used instead of “*”.
- Vector-vector operations: If A and B are vectors of the same length, then $A \times B$ is also a vector of this length. Here \times stands for any algebraic operator, and the operation is done for each pair of corresponding elements. For instance, $A + B$ adds the vectors A and B , and $A * B$ multiplies their elements. Note that the latter case is not a scalar product. If an attempt is made to connect two vectors of different lengths by an algebraic operator, an error occurs.
- Intrinsic functions: Assume $THETA=[100,110,120,-130]$ to be a vector of angles (in degrees). In this case $X=2*COS(THETA)$ is also a vector containing the cosines of each element of $THETA$ multiplied by two, i.e., $X(i) = 2*COS(THETA(i))$. $MAX(THETA)$ or $MIN(THETA)$ return the maximum and minimum values, respectively, in array $THETA$. Vector operations can also be nested, e.g., $MAX(ABS(THETA))$ returns the maximum value in array $ABS(THETA)$.

At present, vector operations are not supported with `string` variables.

6.7 Special variables

6.7.1 Variables set by the program

A number of variables are predefined by the program. The following variables can be used to convert between atomic units and other units:

```

EV=1.d0/27.2113961d0 HARTREE
KELVIN=1.d0/3.157733d5 HARTREE
KJOULE=1.d0/2625.500d0 HARTREE
KCAL=1.d0/627.5096d0 HARTREE
CM=1.d0/219474.63067d0 HARTREE
CM-1=1.d0/219474.63067d0 HARTREE
HZ=1.d0/6.5796838999d15 HARTREE
HERTZ=1.d0/6.5796838999d15 HARTREE
ANG=1.d0/0.529177249d0 BOHR
ANGSTROM=1.d0/0.529177249d0 BOHR

TOEV=27.2113961d0 EV
TOK=3.157733d5 K
TOKELVIN=3.157733d5 K
TOCM=219474.63067d0 CM-1
TOHERTZ=6.5796838999d15 HZ
TOHZ=6.5796838999d15 HZ
TOKJ=2625.500d0 KJ/MOL
TOKJOULE=2625.500d0 KJ/MOL
TOKCAL=627.5096d0 KCAL/MOL
TOA=0.529177249d0 ANGSTROM
TOANG=0.529177249d0 ANGSTROM
TODEBYE=2.54158d0 DEBYE

```

Further variables which are set during execution of the program:

INTYP	defines integral program to be used. Either INTS (Seward) or INTP (Argos).
INTDONE	has the value <code>.true.</code> if the integrals are done for the current geometry.
CARTESIAN	Set to one if cartesian basis functions are used.
SCFDONE	has the value <code>.true.</code> if an SCF calculation has been done for the current geometry.
NUMVAR	number of variables presently defined
STATUS	status of last step (1=no error, -1=error or no convergence)
CHARGE	Total charge of the molecule
NELEC	number of electrons in last wavefunction
SPIN	spin multiplicity minus one of last wavefunction
ORBITAL	record of last optimized orbitals (set but never used in the program)
LASTORB	Type of last optimized orbitals (RHF, UHF, UHFNAT, or MCSCF).
LASTSYM	Symmetry of wavefunction for last optimized orbitals.
LASTSPIN	$2 * M_S$ for wavefunctions for last optimized orbitals.
LASTNELEC	Number of electrons in wavefunction for last optimized orbitals.
ENERGR(<i>istate</i>)	Reference energy for state <i>istate</i> in MRCI and CCSD.
ENERGY(<i>istate</i>)	last computed total energy for state <i>istate</i> for the method specified in the input (e.g., HF, MULTI, CCSD(T), or CCSD[T]).
ENERGD(<i>istate</i>)	Total energy for state <i>istate</i> including Davidson correction (set only in CI).
ENERGP(<i>istate</i>)	Total energy for state <i>istate</i> including Pople correction (set only in CI).
ENERGT(1)	Total energy including perturbative triples (T) correction (set only in CCSD(T), QCI(T)).
ENERGT(2)	Total energy including perturbative triples [T] correction (set only in CCSD(T), QCI(T)).
ENERGT(3)	Total energy including perturbative triples -t correction (set only in CCSD(T), QCI(T)).
EMP2	holds MP2 energy in MPn, CCSD, BCCD, or QCISD calculations, and RS2 energy in MRPT2 (CASPT2) calculations.
EMP3	holds MP3 energy in MP3 and MP4 calculations, and RS3 energy in MRPR3 (CASPT3) calculations.
EMP4	holds MP4(SDQ) energy in MP4 calculations. The MP4(SDTQ) energy is stored in variable ENERGY.
METHODC	String variable holding name of the methods used for ENERGC, e.g., CCSD, BCCD, QCI.
METHODT(1)	String variable holding name of the methods used for ENERGT(1), e.g., CCSD(T), BCCD(T), QCI(T).
METHODT(2)	String variable holding name of the methods used for ENERGT(2), e.g., CCSD[T], BCCD[T], QCI[T].
METHODT(3)	String variable holding name of the methods used for ENERGT(3), e.g., CCSD-T, BCCD-T, QCI-T.

ENERGC	Total energy excluding perturbative triples correction (set only in QCI or CCSD with triples correction enabled).
DFTFUN	total value of density functional in DFT or KS.
DFTFUNS(<i>ifun</i>)	value of <i>ifun</i> 'th component of density functional in DFT or KS.
DFTNAME(<i>ifun</i>)	name of <i>ifun</i> 'th component of density functional in DFT or KS.
DFTFAC(<i>ifun</i>)	factor multiplying <i>ifun</i> 'th component of density functional in DFT or KS.
DFTEXFAC	factor multiplying exact exchange in KS.
PROP(<i>istate</i>)	computed property for state <i>istate</i> . See below for the names PROP of various properties.
PROGRAM	last program called, as specified in the input (e.g., HF, CCSD(T), etc.)
ITERATIONS	Number of iterations used. Set negative if no convergence or max number of iterations reached.
CPUSTEP	User-CPU time in seconds for last program called.
SYSSTEP	System-CPU time in seconds for last program called.
WALLSTEP	Elapsed time in seconds for last program called.

The variable names for properties are the same as used on the EXPEC input cards.

OV	Overlap
EKIN	Kinetic energy
POT	Potential
DELT	Delta function
DEL4	∇^4
DARWIN	Darwin term of relativistic correction
MASSV	Mass-velocity term of relativistic correction
EREL	Total relativistic correction
DMX, DMY, DMZ	Dipole moments
XX, YY, ZZ, XY, XZ, YX	Second moments
XXX, XXY, XXZ, XYY, XYZ, XZZ, YYY, YYZ, YZZ, ZZZ	Third moments
QMXX, QMYX, QMZZ, QMXY, QMXZ, QMZY	Quadrupole moments
EFX, EFX, EFZ	Electric field
FGXX, FGYY, FGZZ, FGXY, FGXZ, FGZY	Electric field gradients
D/DX, D/DY, D/DZ	Velocity
LSX, LSY, LSZ	One-electron spin-orbit
LL	Total angular momentum squared L^2
LX, LY, LZ	Electronic angular momentum
LXLX, LYLY, LZLZ, LXLY, LXLZ, LYLZ	Two-electron angular momentum

By default, only the dipole moments are computed and defined. The values of other properties are only stored in variables if they are requested by EXPEC cards. If more than one state is computed (e.g., in state-averaged MCSCF, corresponding arrays `PROP(istate)` are returned. If properties are computed for more than one center, the center number is appended to the name, e.g. EFX1, EFX2 etc.

If transition properties are computed, their values are stored in corresponding variables with prefix TR, e.g., TRDMX, TRDMY, TRDMZ for transition dipole moments. If more than two states are computed, the index is $(i-1)*(i-2)/2+j$, where $i > j \geq 1$ are state numbers. In a state-averaged calculation, states are counted sequentially for all state symmetries.

For instance, in the following state-averaged MCSCF

```
MULTI ; WF , 14 , 1 , 0 ; STATE , 3 ; WF , 14 , 2 , 0 ; STATE , 2 ; WF , 3 , 0
```

the states are counted as

<i>i</i>	1	2	3	4	5	6
Symmetry	1	1	1	2	2	3
Root in Sym.	1	2	3	1	2	1

6.7.2 Variables recognized by the program

All variables described below are checked by the program, but not set (except NELEC and SPIN). If these are not defined by the user, the program uses its internal defaults. The variables have no effect if the corresponding input cards are present.

Variables recognized by the SCF program:

CHARGE	Total charge of the molecule (can be given instead of nelec)
NELEC	number of electrons
SPIN	spin multiplicity minus one
SCFSYM[METRY]	wavefunction symmetry
SYMMETRY	as SCFSYMM; only used if SCFSYMM is not present.
SCFOC[C]	number of occupied orbitals in each symmetry for SCF
SCFCL[OSED]	number of closed-shell orbitals in each symmetry for SCF
SCFORB	record of saved orbitals in SCF
SCFSTART	record of starting orbitals used in SCF

Variables recognized by the MCSCF program:

CHARGE	Total charge of the molecule (can be given instead of nelec)
NELEC	number of electrons
MCSYM[METRY]	wavefunction symmetry. This can be an array for state-averaged calculations.
SYMMETRY	as MCSYMM; only used if MCSYMM is not present.
MCSPIN	spin multiplicity minus one. This can be an array for state-averaged calculations, but different spin multiplicities can only be used in determinant CASSCF. If only one value is specified, this is used for all states

SPIN	as MCSPIN; only used if MCSPIN is not present.
MCSTATE	number of states for each symmetry in MCSCF
STATE	as MCSTATE; only used if MCSTATE is not present.
WEIGHT	weight factors for all states defined by SYMMETRY and STATE
LQUANT	Eigenvalues of L_z^2 for linear molecules for each state defined by SYMMETRY and STATE.
MCSELECT	records from which configurations can be selected and selection threshold
SELECT	as MCSELECT; only used if MCSELECT is not present.
MCRESTRIC	can be used to define occupancy restrictions
RESTRICT	as MCRESTRIC; only used if MCRESTRIC is not present:
CONFIG	if set to .true. or to one triggers use of CSFs
MCOCC [C]	number of occupied orbitals in each symmetry
OCC	as MCOCC; only used if MCOCC is not present.
MCCLOS [SED]	number of optimized closed-shell orbitals in each symmetry
CLOSED	as MCCLOSED; only used if MCCLOSED is not present.
MCCO [RE]	number of frozen core orbitals in each symmetry
CORE	as MCCORE; only used if MCCORE is not present.
MCSTART	record of starting orbitals
COREORB	record of frozen core orbitals
MCORB	record for saving optimized orbitals
MCSAVE	records for saving CI wavefunction (like SAVE card in MCSCF)

Variables recognized by the CI/CCSD program:

CHARGE	Total charge of the molecule (can be given instead of nelec)
NELEC	number of electrons
SPIN	spin multiplicity minus one
CISYM [METRY]	wavefunction symmetry. If this is an array, only SYMMETRY(1) is used.
SYMMETRY	as CISYMM; only used if CISYMM is not present.
CISTATE	number of states in CI
STATE	as CISTATE, only used if CISTATE is not present.
CISELECT	records from which configurations can be selected
SELECT	as CISELECT; only used if CISELECT is not present.
CIRESTRIC	defines occupancy restrictions
RESTRICT	as RESTRICT; only used if CIRESTRIC is not present.
CIOCC [C]	number of occupied orbitals in each symmetry
OCC	as CIOCC; only used if CIOCC is not present.
CICLOS [SED]	number of closed-shell orbitals in each symmetry
CLOSED	as CICLOSED; only used if CICLOSED is not present.

CICO[RE]	number of core orbitals in each symmetry
CORE	as CICO; only used if CICO is not present.
CIORB	record of orbitals used in CI
CISAVE	records for saving CI wavefunction (like SAVE card in CI)
CISTART	records for restarting with previous CI wavefunction (like START card in CI)

Variables recognized by the DFT/KS program:

DF(ifun) or DFTNAME(ifun)	name of ifun'th component of density functional.
DFTFAC(ifun)	factor multiplying ifun'th component of density functional.
DFTEXFAC	factor multiplying exact exchange in KS.

Example for the use of these variables for a state-averaged MCSCF (note that system variables can only be modified using the SET command, see section 6.3):

SET, NELEC=9	defines number of electrons
SET, SPIN=1	defines wavefunction to be a doublet
SET, SYMMETRY=[1, 2, 3]	defines wavefunction symmetries for state averaged calculation
SET, STATE=[2, 1, 1]	defines number of states to be averaged in each symmetry
WEIGHT=[2, 2, 1, 1]	defines weights for the above four states
OCC=[5, 2, 2]	number of occupied orbitals in each symmetry
CLOSED=2	number of closed-shell orbitals in symmetry 1
MCORB=3100.2	record for optimized orbitals
MULTI	do mcscf with above parameters

6.8 Displaying variables

Variables or the results of expressions can be displayed in the output using SHOW and TABLE.

6.8.1 The SHOW command

The general form of the SHOW command is as follows:

SHOW[*ncol*, *format*], *expression*

where *expression* can be an expression or variable, *ncol* is the number of values printed per line (default 6), and *format* is a format (default 6F15.8). This can be used to print vectors in matrix form. The specification of *ncol* and *format* is optional. Assume that E is a vector:

SHOW, E	prints E using defaults.
SHOW[<i>n</i>], E	prints E with <i>n</i> elements per line; (if <i>n</i> >6, more than one line is needed, but in any case a new line is started after <i>n</i> elements).
SHOW[<i>n</i> , 10f10.4], E	prints E in the format given, with newline forced after <i>n</i> elements.

Note that the total length of the format should not exceed 100 characters (a left margin of 30 characters is always needed).

A *wild card* format can be used to show several variables more easily:

```
SHOW , qm* , dm*
```

shows all variables whose names begin with QM and DM. Note that no letters must appear after the *, i.e., the wild card format is less general than in UNIX commands.

See the TABLE command for another possibility to tabulate results.

6.9 Clearing variables

Variables can be deleted using

```
CLEAR,name1, name2, ...
```

Wild cards can be used as in SHOW, e.g.,

```
CLEAR , ENERG*
```

clears all variables whose names begin with ENERG. All variables can be cleared using

```
CLEARALL
```

The length of vectors can be truncated simply by redefining the length specifier: #R=2 truncates the array R to length 2. Higher elements are no longer available (but could be redefined). Setting #R=0 is equivalent to the command CLEAR , R.

7 TABLES AND PLOTTING

7.1 Tables

Variables can be printed in Table form using the command

```
TABLE,var1,var2,...
```

The values of each variable are printed in one column, so all variables used must be defined for the same range, and corresponding elements should belong together. For example, if in a calculation one has stored $R(i)$, $THETA(i)$, $ECI(i)$ for each geometry i , one can print these data simply using

```
TABLE , R , THETA , ECI
```

By default, the number of rows equals the number of elements of the first variable. This can be changed, however, using the RANGE subcommand.

The first ten columns of a table may contain string variables. For instance,

```
hf;etot(1)=energy;method(1)=program;cpu(1)=cpustep
ccsd;etot(2)=energy;method(2)=program;cpu(2)=cpustep
qci;etot(3)=energy;method(3)=program;cpu(3)=cpustep
table,method,etot,cpu
```

prints a table with the SCF, CCSD, and QCI results in the first, second, and third row, respectively. For other use of string variables and tables see, e.g. the examples `h2o_tab.com` and `oh_macros.com`

The appearance of the table may be modified using the following commands, which may be given (in any order) directly after the the TABLE card:

HEADING , <i>head1</i> , <i>head2</i> ,...	Specify a heading for each column. By default, the names of the variables are used as headings.
FORMAT , <i>format</i>	Specify a format for each row in fortran style. <i>format</i> must be enclosed by quotes. Normally, the program determines automatically an appropriate format, which depends on the type and size of the printed data.
FTYP, <i>typ1</i> , <i>typ2</i> , <i>typ3</i> , ...	Simplified form to modify the format. This gives the type (A, F, or D) for each column (sensible defaults are normally used).
DIGITS, <i>dig1</i> , <i>dig2</i> , <i>dig3</i> , ...	Give the number of digits after the decimal points to be printed for each column (sensible defaults are normally used).
SAVE, <i>file</i> , <i>status</i>	Specify a file on which the table will be written. If status is NEW, the file is rewound, otherwise it is appended.
TITLE , <i>title</i>	Specify one line of a title (several TITLE cards may follow each other). Note that titles are only displayed in the SAVE file, if the SAVE command is given before the TITLE card.
SORT , <i>col1</i> , <i>col2</i> ,...	Sort rows according to increasing values of the given columns. The columns are sorted in the order they are specified.
PRINT , <i>key1</i> , <i>key2</i> ,...	Specify print options (TABLE, HEADING, TITLE, WARNING, FORMAT, SORT). The default is print for the first three, and noprint for the last three.
NOPRINT , <i>key1</i> , <i>key2</i> ,...	Disable print for given keys.
NOPUNCH	Don't write data to the punch file (data are written by default).
RANGE , <i>start</i> , <i>end</i>	Specify start and end indices of the variables to be printed.
STATISTICS	Print also linear regression and quadratic fits of the data columns.

7.2 Plotting

[PLOT , [[CMD=]*unix_plot_command*],[FILE=*plotfile*],[NOPLOT]

Execute a plotting program using the table as data. PLOT is a subcommand of TABLE and must follow TABLE or any of its valid subcommands given in the previous section. *unix_plot_command* consists of the unix command needed to start the plotting program, followed by any required options. The whole thing should normally be enclosed in quotation marks to preserve lower-case letters. The default is '*xmgrace*'. At present, only the *xmgrace*, *grace*, *gracebat* and *xmgr* programs with all numerical data are supported, although use of *xmgr* is deprecated, and may not be possible in future versions.

By default the input file for the plotting program is saved in *molpro_plot.dat*. The name of the plotfile can be modified using the FILE (or PLOTFILE) directive. FILE implies that the plot is not shown on the screen but all plot data are saved in the given file. The plot on the screen can also be suppressed with the NOPLOT option.

8 INTEGRAL-DIRECT CALCULATIONS (GDIRECT)

References:

Direct methods, general: M. Schütz, R. Lindh, and H.-J. Werner, *Mol. Phys.* **96**, 719 (1999).
 Linear scaling LMP2: M. Schütz, G. Hetzer, and H.-J. Werner *J. Chem. Phys.* **111**, 5691 (1999).

All methods implemented in MOLPRO apart from full CI (FCI) and perturbative triple excitations (T) can be performed integral-direct, i.e., the methods are integral driven with the two-electron integrals in the AO basis being recomputed whenever needed, avoiding the bottleneck of storing these quantities on disk. For small molecules, this requires significantly more CPU time, but reduces the disk space requirements when using large basis sets. However, due to efficient prescreening techniques, the scaling of the computational cost with molecular size is lower in integral-direct mode than in conventional mode, and therefore integral-direct calculations for extended molecules may even be less expensive than conventional ones. The break-even point depends strongly on the size of the molecule, the hardware, and the basis set. Depending on the available disk space, calculations with more than 150–200 basis functions in one symmetry should normally be done in integral-direct mode.

Integral-direct calculations are requested by the DIRECT or GDIRECT directives. If one of these cards is given outside the input of specific programs it acts globally, i.e. all subsequent calculations are performed in integral-direct mode. On the other hand, if the DIRECT card is part of the input of specific programs (e.g. HF, CCSD), it affects only this program. The GDIRECT directive is not recognized by individual programs and always acts globally. Normally, all calculations in one job will be done integral-direct, and then a DIRECT or GDIRECT card is required before the first energy calculation. However, further DIRECT or GDIRECT directives can be given in order to modify specific options or thresholds for particular programs.

The integral-direct implementation in MOLPRO involves three different procedures: (i) Fock matrix evaluation (DFOCK), (ii) integral transformation (DTRAF), and (iii) external exchange operators (DKEXT). Specific options and thresholds exist for all three programs, but it is also possible to specify the most important thresholds by general parameters, which are used as defaults for all programs.

Normally, appropriate default values are automatically used by the program, and in most cases no parameters need to be specified on the DIRECT directive. However, in order to guarantee sufficient accuracy, the default thresholds are quite strict, and in calculations for extended systems larger values might be useful to reduce the CPU time.

The format of the DIRECT directive is

```
DIRECT, key1=value1, key2=value2...
```

The following table summarizes the possible keys and their meaning. The default values are given in the subsequent table. In various cases there is a hierarchy of default values. For instance, if THREST_D2EXT is not given, one of the following is used: [THR_D2EXT, THREST_DTRAF, THR_DTRAF, THREST, *default*]. The list in brackets is checked from left to right, and the first one found in the input is used. *default* is a default value which depends on the energy threshold and the basis set (the threshold is reduced if the overlap matrix contains very small eigenvalues).

General Options (apply to all programs):

THREST	Integral prescreening threshold. The calculation of an integral shell block is skipped if the product of the largest estimated integral value (based on the Cauchy-Schwarz inequality) and the largest density matrix element contributing to the shell block is
--------	--

smaller than this value. In DTRAF and DKEXT effective density matrices are constructed from the MO coefficients and amplitudes, respectively.

THRINT	Integral prescreening threshold. This applies to the product of the exact (i.e. computed) integral value and a density matrix. This threshold is only used in DTRAF and DKEXT. A shell block of integrals is skipped if the product of the largest integral and the largest element of the effective density matrix contributing to the shell block is smaller than this threshold. If it set negative, no computed integrals will be neglected.
THRPROD	Prescreening threshold for products of integrals and MO-coefficients (DTRAF) or amplitudes (DKEXT). Shell blocks of MO coefficients or amplitudes are neglected if the product of the largest integral in the shell block and the largest coefficient is smaller than this value. If this is set negative, no product screening is performed.
THRMAX	Initial value of the prescreening threshold THREST for DFOCK and DKEXT in iterative methods (SCF, CI, CCSD). If nonzero, it will also be used for DKEXT in MP3 and MP4 (SDQ) calculations. The threshold will be reduced to THREST once a certain accuracy has been reached (see VARRED), or latest after MAXRED iterations. In CI and CCSD calculations, also the initial thresholds THRINT_DKEXT and THRPROD_DKEXT are influenced by this value. For a description, see THRMAX_DKEXT. If THRMAX=0, the final thresholds will be used from the beginning in all methods.
SCREEN	Enables or disables prescreening. SCREEN ≥ 0: full screening enabled. SCREEN < 0: THRPROD is unused. No density screening in direct SCF. SCREEN < -1: THRINT is unused. SCREEN < -2: THREST is unused.
MAXRED	Maximum number of iterations after which thresholds are reduced to their final values in CI and CCSD calculations. If MAXRED=0, the final thresholds will be used in CI and CCSD from the beginning (same as THRMAX=0, but MAXRED has no effect on DSCF. In the latter case a fixed value of 10 is used.
VARRED	Thresholds are reduced to their final values if the sum of squared amplitude changes is smaller than this value.
SWAP	Enables or disables label swapping in SEWARD. Test purpose only.

Specific options for direct SCF (DFOCK):

THREST_DSCF	Final prescreening threshold in direct SCF. If given, it replaces the value of THREST.
THRMAX_DSCF	Initial prescreening threshold in direct SCF. This is used for the first 7-10 iterations. Once a certain accuracy is reached, the threshold is reduced to THREST_DSCF
DISKSIZE	Max disk size in MB allowed in semi-direct SCF calculations.

THRDISK	Only integrals larger than this threshold are stored on disk in semi-direct SCF calculations.
MEM_DFOCK	Integral buffer size for semi-direct SCF. If this is larger than DISKSIZE, the calculation is done semi-direct incore (no integrals are written to disk). If MEM_DFOCK is negative (default), a default buffer size is used. If MEM_DFOCK is zero, all available memory is used.
SWAP_DFOCK	Enables or disables label swapping in fock matrix calculation (test purpose only).

General options for direct integral transformation (DTRAF):

PAGE_DTRAF	Selects the transformation method. PAGE_DTRAF=0: use minimum memory algorithm, requiring four integral evaluations. PAGE_DTRAF=1: use paging algorithm, leading to the minimum CPU time (one integral evaluation for DMP2/LMP2 and two otherwise).
SCREEN_DTRAF	If given, replaces value of SCREEN for DTRAF.
MAXSHLQ1_DTRAF	Maximum size of merged shells in the first quarter transformation step (0: not used).
MINSHLQ1_DTRAF	Shells are only merged if their size is smaller than this value (0: not used).
MAXSHLQ2_DTRAF	Maximum size of merged shells in the second quarter transformation step (0: not used).
MINSHLQ2_DTRAF	Shells are only merged if their size is smaller than this value (0: not used).
MAXCEN_DTRAF	Maximum number of centres in merged shells (0: no limit).
PRINT_DTRAF	Print parameter for DTRAF.

General thresholds for all direct integral transformations:

THR_DTRAF	General threshold for DTRAF. If given, this is taken as default value for all thresholds described below.
THREST_DTRAF	AO prescreening threshold for DTRAF. Defaults: [THR_DTRAF, THREST, <i>default</i>].
THRINT_DTRAF	Integral threshold for DTRAF. Defaults: [THR_DTRAF, THRINT, <i>default</i>].
THRPROD_DTRAF	Product threshold for DTRAF. Defaults: [THR_DTRAF, THRPROD, <i>default</i>].

Thresholds specific to direct integral transformations:

THR_D2EXT	General threshold for generation of 2-external integrals. If given, this is used as a default for all D2EXT thresholds described below.
THREST_D2EXT	Prescreening threshold for generation of 2-external integrals. Defaults: [THR_D2EXT, THREST_DTRAF, THR_DTRAF, THREST, <i>default</i>].

THRINT_D2EXT	Integral threshold for generation of 2-external integrals. Defaults: [THR_D2EXT, THRINT_DTRAF, THR_DTRAF, THRINT, <i>default</i>].
THRPROD_D2EXT	Product threshold for generation of 2-external integrals. Defaults: [THR_D2EXT, THRPROD_DTRAF, THR_DTRAF, THRPROD, <i>default</i>].
THR_D3EXT	General threshold for generation of 3-external integrals. If given, this is used as a default for all D3EXT thresholds described below.
THREST_D3EXT	Prescreening threshold for generation of 3-external integrals. Defaults: [THR_D3EXT, THREST_DTRAF, THR_DTRAF, THREST, <i>default</i>].
THRINT_D3EXT	Integral threshold for generation of 3-external integrals. Defaults: [THR_D3EXT, THRINT_DTRAF, THR_DTRAF, THRINT, <i>default</i>].
THRPROD_D3EXT	Product threshold for generation of 3-external integrals. Defaults: [THR_D3EXT, THRPROD_DTRAF, THR_DTRAF, THRPROD, <i>default</i>].
THR_D4EXT	General threshold for generation of 4-external integrals. If given, this is used as a default for all D4EXT thresholds described below.
THREST_D4EXT	Prescreening threshold for generation of 4-external integrals. Defaults: [THR_D4EXT, THREST_DTRAF, THR_DTRAF, THREST, <i>default</i>].
THRINT_D4EXT	Integral threshold for generation of 4-external integrals. Defaults: [THR_D4EXT, THRINT_DTRAF, THR_DTRAF, THRINT, <i>default</i>].
THRPROD_D4EXT	Product threshold for generation of 4-external integrals. Defaults: [THR_D4EXT, THRPROD_DTRAF, THR_DTRAF, THRPROD, <i>default</i>].
THR_DCCSD	General threshold for generalized transformation needed in each CCSD iteration. If given, this is used as a default for THREST_DCCSD, THRINT_DCCSD, and THRPROD_DCCSD described below.
THREST_DCCSD	Prescreening threshold for DCCSD transformation. Defaults: [THR_DCCSD, THREST_DTRAF, THR_DTRAF, THREST, <i>default</i>].
THRINT_DCCSD	Integral threshold for DCCSD transformation. Defaults: [THR_DCCSD, THRINT_DTRAF, THR_DTRAF, THRINT, <i>default</i>].
THRPROD_DCCSD	Product threshold for DCCSD transformation. Defaults: [THR_DCCSD, THRPROD_DTRAF, THR_DTRAF, THRPROD, <i>default</i>].
THRMAX_DCCSD	Initial value for THREST_DCCSD in CCSD calculations. The threshold will be reduced to THREST_DCCSD once a certain accuracy has been reached (see VARRED), or latest after MAXRED

iterations. The initial thresholds THRINT_DCCSD and THRPROD_DCCSD are obtained by multiplying their input (or default) values by THRMAX_DCCSD/THREST_DCCSD, with the restriction that the initial values cannot be smaller than the final ones.

Specific options for direct MP2 (DMP2):

DMP2	<p>Selects the transformation method for direct MP2: DMP2=-1: automatic selection, depending on the available memory. DMP2=0: use fully direct method for DMP2 (min. two integral evaluations, possibly multipassing, no disk space). DMP2=1: use semi-direct method for DMP2 (one to four integral evaluations, depending on PAGE_DTRAF). DMP2=2: use DKEXT to compute exchange operators in DMP2 (one integral evaluation). This is only useful in local DMP2 calculations with many distant pairs.</p>
THR_DMP2	<p>General threshold for generation of 2-external integrals in DMP2. If given, this is used as a default for all DMP2 thresholds described below.</p>
THREST_DMP2	<p>Prescreening threshold for generation of 2-external integrals. Defaults: [THR_DMP2, THREST_DTRAF, THR_DTRAF, THREST, <i>default</i>].</p>
THRINT_DMP2	<p>Integral threshold for generation of 2-external integrals. Defaults: [THR_DMP2, THRINT_DTRAF, THR_DTRAF, THRINT, <i>default</i>].</p>
THRPROD_DMP2	<p>Product threshold for generation of 2-external integrals. Defaults: [THR_DMP2, THRPROD_DTRAF, THR_DTRAF, THRPROD, <i>default</i>].</p>

Specific options for direct local MP2 (LMP2):

DTRAF	<p>Selects the transformation method for direct LMP2: DTRAF ≥ 0: generates the 2-external integrals (exchange operators) first in AO basis and transforms these thereafter in a second step to the projected, local basis. The disk storage requirements hence scale cubically with molecular size. DTRAF = -1: generates the 2-external integrals (exchange operators) directly in projected basis. The disk storage requirements hence scale linearly with molecular size. This (together with PAGE_DTRAF = 0) is the recommended algorithm for very large molecules (cf. linear scaling LMP2, chapter 23). DTRAF = -2: alternative algorithm to generate the exchange operators directly in projected basis. Usually, this algorithm turns out to be computationally more expensive than the one selected with DTRAF = -1. Note, that neither DTRAF = -1 nor DTRAF = -2 work in the context of LMP2 gradients.</p>
THR_LMP2	<p>General threshold for generation of 2-external integrals in linear scaling LMP2. If given, this is used as a default for all LMP2 thresholds described below.</p>

THREST_LMP2	Prescreening threshold for generation of 2-external integrals. Defaults: [THR_LMP2, THREST_DTRAF, THR_DTRAF, THREST, <i>default</i>].
THRQ1_LMP2	Threshold used in the first quarter transformation. Defaults: [THR_LMP2, THRPROD_DTRAF, THR_DTRAF, THRPROD, <i>default</i>].
THRQ2_LMP2	Threshold used in the second and subsequent quarter transformations. Defaults: [THR_LMP2, THRINT_DTRAF, THR_DTRAF, THRINT, <i>default</i>].
THRAO_ATTEN	Special threshold for prescreening of attenuated integrals ($\mu\mu vv$) Default: THREST_LMP2

Options for integral-direct computation of external exchange operators (DKEXT):

DKEXT	Selects driver for DKEXT. DKEXT=-1: use paging algorithm (minimum memory). This is automatically used if in-core algorithm would need more than one integral pass. DKEXT=0: use in-core algorithm, no integral triples. DKEXT=1: use in-core algorithm and integral triples. DKEXT=2: use in-core algorithm and integral triples if at least two integrals of a triple differ. DKEXT=3: use in-core algorithm and integral triples if all integrals of a triple differ.
SCREEN_DKEXT	if given, replaces value of SCREEN for DKEXT.
MAXSIZE_DKEXT	Largest size of merged shells in DKEXT (0: not used).
MINSIZE_DKEXT	Shells are only merged if their size is smaller than this value. (0: not used).
MAXCEN_DKEXT	Maximum number of centres in merged shells (0: no limit).
SCREEN_DKEXT	Enables or disables screening in DKEXT.
PRINT_DKEXT	Print parameter for DKEXT.
SWAP_DKEXT	Enables or disables label swapping in DKEXT (test purpose only)
MXMBLK_DKEXT	Largest matrix block size in DKEXT (only used with DKEXT \geq 1).

Thresholds for integral-direct computation of external exchange operators (DKEXT):

THR_DKEXT	General threshold for DKEXT. If given, this is used as a default for all DKEXT thresholds described below.
THREST_DKEXT	Prescreening threshold for DKEXT. Defaults: [THR_DKEXT, THREST, <i>default</i>].
THRINT_DKEXT	Integral threshold for DKEXT. Defaults: [THR_DKEXT, THRINT, <i>default</i>].
THRPROD_DKEXT	Product threshold for DKEXT. Defaults: [THR_DKEXT, THRPROD, <i>default</i>].

THRMAX_DKEXT Initial value for THREST_DKEXT in CI, and CCSD calculations. If nonzero, it will also be used for DKEXT in MP3 and MP4 (SDQ) calculations. The threshold will be reduced to THREST_DKEXT once a certain accuracy has been reached (see VARRED), or latest after MAXRED iterations. The initial thresholds THRINT_DKEXT and THRPROD_DKEXT are obtained by multiplying their input (or default) values by THRMAX_DKEXT/THREST_DKEXT with the restriction that the initial values cannot be smaller than the final ones.

For historical reasons, many options have alias names. The following tables summarize the default values for all options and thresholds and also gives possible alias names.

Table 6: Default values and alias names for direct options.

Parameter	Alias	Default value
SCREEN		1
MAXRED		7
VARRED		1.d-7
SWAP		1
DISKSIZE		0
THRDISK		1.d-3
MEM_DFOCK	BUF_DFOCK	-1
SWAP_DFOCK		SWAP
DMP2	DTRAF	-1
PAGE_DTRAF	PAGE	1
SCREEN_DTRAF		SCREEN
MAXSHLQ1_DTRAF	NSHLQ1	32
MINSHLQ1_DTRAF		0
MAXSHLQ2_DTRAF	NSHLQ2	16
MINSHLQ2_DTRAF		0
MAXCEN_DTRAF		0
PRINT_DTRAF		-1
SWAP_DTRAF		SWAP
DKEXT	DRVKEXT	3
SCREEN_DKEXT		SCREEN
MAXSIZE_DKEXT		0
MINSIZE_DKEXT		5
MAXCEN_DKEXT		1
PRINT_DKEXT		-1
SWAP_DKEXT		SWAP
MXMBLK_DKEXT		depends on hardware (-B parameter on molpro command)

Table 7: Default thresholds and alias names for direct calculations

Parameter	Alias	Default value
THREST	THRAO	$\min(\Delta E \cdot 1.d - 2, 1.d - 9)^{a,b}$
THRINT	THRSO	$\min(\Delta E \cdot 1.d - 2, 1.d - 9)^{a,b}$
THRPROD	THRP	$\min(\Delta E \cdot 1.d - 3, 1.d - 10)^{a,b}$
THRMAX		$1.d \cdot 8^b$
THREST_DSCF	THRDSCF	$\leq 1.d \cdot 10$ (depending on accuracy and basis set)
THRMAX_DSCF	THRDSCF_MAX	THRMAX
THR_DTRAF	THRDTRAF	
THREST_DTRAF	THRAO_DTRAF	[THR_DTRAF, THREST]
THRINT_DTRAF	THRAO_DTRAF	[THR_DTRAF, THRINT]
THRPROD_DTRAF	THRP_DTRAF	[THR_DTRAF, THRPROD]
THR_D2EXT	THR2EXT	THR_DTRAF
THREST_D2EXT	THRAO_D2EXT	[THR_D2EXT, THREST_DTRAF]
THRINT_D2EXT	THRSO_D2EXT	[THR_D2EXT, THRINT_DTRAF]
THRPROD_D2EXT	THRP_D2EXT	[THR_D2EXT, THRPROD_DTRAF]
THR_D3EXT	THR3EXT	THR_DTRAF
THREST_D3EXT	THRAO_D3EXT	[THR_D3EXT, THREST_DTRAF]
THRINT_D3EXT	THRSO_D3EXT	[THR_D3EXT, THRINT_DTRAF]
THRPROD_D3EXT	THRP_D3EXT	[THR_D3EXT, THRPROD_DTRAF]
THR_D4EXT	THR4EXT	THR_DTRAF
THREST_D4EXT	THRAO_D4EXT	[THR_D4EXT, THREST_DTRAF]
THRINT_D4EXT	THRSO_D4EXT	[THR_D4EXT, THRINT_DTRAF]
THRPROD_D4EXT	THRP_D4EXT	[THR_D4EXT, THRPROD_DTRAF]
THR_DCCSD	THRCCSD	THR_DTRAF
THREST_DCCSD	THRAO_DCCSD	[THR_DCCSD, THREST_DTRAF]
THRINT_DCCSD	THRSO_DCCSD	[THR_DCCSD, THRINT_DTRAF]
THRPROD_DCCSD	THRP_DCCSD	[THR_DCCSD, THRPROD_DTRAF]
THRMAX_DCCSD	THRMAX_DTRAF	THRMAX
THR_DMP2	THRDMP2	THR_DTRAF
THREST_DMP2	THRAO_DMP2	[THR_DMP2, THREST_DTRAF, <i>default^c</i>]
THRINT_DMP2	THRSO_DMP2	[THR_DMP2, THRINT_DTRAF, <i>default^c</i>]
THRPROD_DMP2	THRP_DMP2	[THR_DMP2, THRPROD_DTRAF, <i>default^c</i>]
THR_LMP2	THRLMP2	THR_DTRAF
THREST_LMP2	THRAO_LMP2	[THR_LMP2, THREST_DTRAF, <i>default^c</i>]
THRQ1_LMP2	THRQ1	[THR_LMP2, THRPROD_DTRAF, <i>default^c</i>]
THRQ2_LMP2	THRQ2	[THR_LMP2, THRINT_DTRAF, <i>default^c</i>]
THRAO_ATTEN]	THRATTEN	THREST_LMP2
THR_DKEXT	THRKEXT	
THREST_DKEXT	THRAO_DKEXT	[THR_DKEXT, THREST]
THRINT_DKEXT	THRSO_DKEXT	[THR_DKEXT, THRINT]
THRPROD_DKEXT	THRP_DKEXT	[THR_DKEXT, THRPROD]
THRMAX_DKEXT		THRMAX

a) ΔE is the requested accuracy in the energy (default 1.d-6).

b) The thresholds are reduced if the overlap matrix has small eigenvalues.

c) The default thresholds for DMP2 and LMP2 are $0.1 \cdot \Delta E$.

8.1 Example for integral-direct calculations

```

memory,2,m
method=[hf,mp2,ccsd,qci,bccd,multi,mrci,acpf,rs3]
basis=vdz
geometry={o;h1,o,r;h2,o,r,h1,theta}
gdirect
r=1 ang,theta=104
do i=1,#method
$method(i)
e(i)=energy
dip(i)=dmz
enddo
table,method,e,dip

```

```

!some methods
!basis
!geometry
!direct option
!bond length and angle
!loop over methods/
!run method(i) h2o_direct.com
!save results in variables

!print table of results

```

This jobs produces the following table:

METHOD	E	DIP
HF	-76.02145798	0.82747348
MP2	-76.22620591	0.00000000
CCSD	-76.23580191	0.00000000
QCI	-76.23596211	0.00000000
BCCD	-76.23565813	0.00000000
MULTI	-76.07843443	0.76283026
MRCI	-76.23369819	0.76875001
ACPF	-76.23820180	0.76872802
RS3	-76.23549448	0.75869972

9 GEOMETRY SPECIFICATION AND INTEGRATION

Before starting any energy calculations, MOLPRO checks if the one- and two-electron integrals are available for the current basis set and geometry and automatically computes them if necessary. It is therefore not necessary any more to call the integral program explicitly, as was done in older MOLPRO versions using the `INT` command. The program also recognizes automatically if only the nuclear charges have been changed, as is the case in counterpoise calculations. In this case, the two-electron integrals are not recomputed.

Before any energy calculation, the geometry and basis set must be defined in `GEOMETRY` and `BASIS` blocks, respectively.

9.1 Sorted integrals

By default, two electron integrals are evaluated once and stored on disk. This behaviour may be overridden by using the input command `gdirect` (see section 8) to force evaluation of integrals on the fly. If the integrals are stored on disk, immediately after evaluation they are sorted into complete symmetry-packed matrices, so that later program modules that use them can do so as efficiently as possible. The options for the integral sort can be specified using the `AOINT` parameter set, using the input form

```
AOINT, key1=value1, key2=value2, ...
```

The following summarizes the possible keys, together with their meaning, and default values.

<code>c_final</code>	Integer specifying the compression algorithm to be used for the final sorted integrals. Possible values are 0 (no compression), 1 (compression using 1, 2, 4 or 8-byte values), 2 (2, 4 or 8 bytes), 4 (4, 8 bytes) and 8. Default: 0
<code>c_sort1</code>	Integer specifying the compression algorithm for the intermediate file during the sort. Default: 0
<code>c_seward</code>	Integer specifying the format of label tagging and compression written by the integral program and read by the sort program. Default: 0
<code>compress</code>	Overall compression; <code>c_final</code> , <code>c_seward</code> and <code>c_sort1</code> are forced internally to be not less than this parameter. Default: 1
<code>thresh</code>	Real giving the truncation threshold for compression. Default: 0.0, which means use the integral evaluation threshold (<code>GTHRESH</code> , <code>TWOINT</code>)
<code>io</code>	String specifying how the sorted integrals are written. Possible values are <code>molpro</code> (standard MOLPRO record on file 1) and <code>eaf</code> (Exclusive-access file). <code>eaf</code> is permissible only if the program has been configured for MPP usage, and at present <code>molpro</code> is implemented only for serial execution. <code>molpro</code> is required if the integrals are to be used in a restart job. For maximum efficiency on a parallel machine, <code>eaf</code> should be used, since in that case the integrals are distributed on separate processor-local files.

For backward-compatibility purposes, two convenience commands are also defined: `COMPRESS` is equivalent to `AOINT, COMPRESS=1`, and `UNCOMPRESS` is equivalent to `AOINT, COMPRESS=0`.

9.2 Symmetry specification

If standard Z-matrix input is used, MOLPRO determines the symmetry automatically by default. However, sometimes it is necessary to use a lower symmetry or a different orientation than obtained by the default, and this can be achieved by explicit specification of the symmetry elements to be used, as described below.

On the first card of the integral input (directly after the INT card or as first card in a geometry block), generating symmetry elements can be given, which uniquely specify the point group. The dimension of the point group is 2^{**} (number of fields given). Each field consists of one or more of X, Y, or Z (with no intervening spaces) which specify which coordinate axes change sign under the corresponding generating symmetry operation. It is usually wise to choose z to be the unique axis where appropriate (essential for C_2 and C_{2h}). In that case, the possibilities are:

(null card)	C_1 (i.e., no point group symmetry)
Z	C_s
XY	C_2
XYZ	C_i
X, Y	C_{2v}
XY, Z	C_{2h}
XZ, YZ	D_2
X, Y, Z	D_{2h}

Note that Abelian point group symmetry only is available, so for molecules with degenerate symmetry, an Abelian subgroup must be used — e.g. C_{2v} or D_{2h} for linear molecules.

See section 2.13 for more details of symmetry groups and ordering of the irreducible representations. Also see section 9.3.1 for more information about automatic generation of symmetry planes.

9.3 Geometry specifications

The geometry may be given in standard Z-matrix form, XYZ form, or cartesian and polar coordinate MOLPRO92 format. The geometry specifications are given in the form

```
geometry={ ,
options
atom specifications
}
```

The following are permitted as options:

	Any valid combination of symmetry generators, as described in the previous section.
NOSYM	Disable use of symmetry.
ANGSTROM	Bond lengths specified by numbers, or variables without associated units, are assumed to be in Å.
CHARGE	Orient molecule such that origin is centre of charge, and axes are eigenvectors of quadrupole moment.

MASS	Orient molecule such that origin is centre of mass, and axes are eigenvectors of inertia tensor (default).
NOORIENT	Disable re-orientation of molecule.
PLANEXZ	For the C_{2v} and D_{2h} point groups, force the primary plane to be xz instead of the default yz . The geometry builder attempts by swapping coordinate axes to place as many atoms as possible in the primary plane, so for the particular case of a planar molecule, this means that all the atoms will lie in the primary plane. The default implements recommendation 5a and the first part of recommendation 5b specified in J. Chem. Phys. 55, 1997 (1955). PLANEYZ and PLANEXY may also be specified, but note that the latter presently generates an error for C_{2v} .

9.3.1 Z-matrix input

The general form of an atom specification line is

[*group* [,]] *atom*, p_1 , r , p_2 , α , p_3 , β , J

or, alternatively,

[*group* [,]] *atom*, p_1 , x , y , z

where

<i>group</i>	atomic group number (optional). Can be used if different basis sets are used for different atoms of the same kind. The basis set is then referred to by this group number and not by the atomic symbol.
<i>atom</i>	chemical symbol of the new atom placed at position p_0 . This may optionally be appended (without blank) by an integer, which can act as sequence number, e.g., C1, H2, etc. Dummy centres with no charge and basis functions are denoted either Q or X, optionally appended by a number, e.g. Q1; note that the first atom in the z-matrix must not be called X, since this may be confused with a symmetry specification (use Q instead).
p_1	atom to which the present atom is connected. This may be either a number n , where n refers to the n 'th line of the Z-matrix, or an alphanumeric string as specified in the <i>atom</i> field of a previous card, e.g., C1, H2 etc. The latter form works only if the atoms are numbered in a unique way.
r	Distance of new atom from p_1 . This value is given in bohr, unless ANG has been specified directly before or after the symmetry specification.
p_2	A second atom needed to define the angle $\alpha(p_0, p_1, p_2)$. The same rules hold for the specification as for p_1 .
α	Internuclear angle $\alpha(p_0, p_1, p_2)$. This angle is given in degrees and must be in the range $0 < \alpha < 180^\circ$.
p_3	A third atom needed to define the dihedral angle $\beta(p_0, p_1, p_2, p_3)$. Only applies if $J = 0$, see below.

β	Dihedral angle $\beta(p_0, p_1, p_2, p_3)$ in degree. This angle is defined as the angle between the planes defined by (p_0, p_1, p_2) and (p_1, p_2, p_3) ($-180^\circ \leq \beta \leq 180^\circ$). Only applies if $J = 0$, see below.
J	If this is specified and nonzero, the new position is specified by two bond angles rather than a bond angle and a dihedral angle. If $J = \pm 1$, β is the angle $\beta(p_0, p_1, p_3)$. If $J = 1$, the triple vector product $(\mathbf{p}_1 - \mathbf{p}_0) \cdot [(\mathbf{p}_1 - \mathbf{p}_2) \times (\mathbf{p}_1 - \mathbf{p}_3)]$ is positive, while this quantity is negative if $J = -1$.
x, y, z	Cartesian coordinates of the new atom. This form is assumed if $p_1 \leq 0$; if $p_1 < 0$, the coordinates are frozen in geometry optimizations.

All atoms, including those related by symmetry transformations, should be specified in the Z-matrix. Note that for the first atom, no coordinates need be given, for the second atom only p_1, r are needed, whilst for the third atom p_3, β, J may be omitted. The 6 missing coordinates are obtained automatically by the program, which translates and re-orientates the molecule such that the origin is at the centre of mass, and the axes correspond to the eigenvectors of the inertia tensor (see also CHARGE option above).

Once the reorientation has been done, the program then looks for symmetry (D_{2h} and subgroups), unless the NOSYM option has been given. It is possible to request that reduced symmetry be used by using appropriate combinations of the options X, Y, Z, XY, XZ, YZ, XYZ. These specify symmetry operations, the symbol defining which coordinate axes change sign under the operation. The point group is constructed by taking all combinations of specified elements. If symmetry is explicitly specified in this way, the program checks to see that the group requested can be used, swapping the coordinate axes if necessary. This provides a mechanism for ensuring that the same point group is used, for example, at all points in the complete generation of a potential energy surface, allowing the safe re-utilization of neighbouring geometry molecular orbitals as starting guesses, etc..

9.3.2 XYZ input

Simple cartesian coordinates in Ångstrom units can be read as an alternative to a Z matrix. This facility is triggered by setting the MOLPRO variable GEOMTYP to the value XYZ before the geometry specification is given. The geometry block should then contain the cartesian coordinates in Minnesota Computer Centre, Inc. XYZ format. Variable names may be used as well as fixed numerical values.

The XYZ file format consists of two header lines, the first of which contains the number of atoms, and the second of which is a title. The remaining lines each specify the coordinates of one atom, with the chemical symbol in the first field, and the x, y, z coordinates following. A sequence number may be appended to the chemical symbol; it is then interpreted as the atomic group number, which can be used when different basis sets are wanted for different atoms of the same kind. The basis set is then specified for this group number rather than the atomic symbol.

```

geomtyp=xyz
geometry={
3          ! number of atoms
This is an example of geometry input for water with an XYZ file
O ,0.0000000000,0.0000000000,-0.1302052882
H ,1.4891244004,0.0000000000, 1.0332262019
H,-1.4891244004,0.0000000000, 1.0332262019
}
hf

```

examples/
h2o_xyzinput.com

The XYZ format is specified within the documentation distributed with MSCI's XMol package. Note that MOLPRO has the facility to write XYZ files with the PUT command (see section 9.4).

9.3.3 MOLPRO92 input

A subset of the MOLPRO92 atom specification commands are retained for compatibility. These may be interspersed with Z-matrix lines, and are of the form

$A[group],atom,x,y,z$

$A[group],atom,POL,r,\theta,\phi$

giving, respectively, cartesian or polar coordinates of the atom to be added. Note that the internal coordinate specifications NPCC, CCPA, TCT, LC, RCP, RCF are no longer available, and Z-matrix input should be used instead.

If any MOLPRO92-style atom specifications appear in the input, the NOORIENT option is enforced, and the handling of symmetry is slightly different. No automatic search for symmetry takes place, and all symmetry required should be specified. Furthermore, only symmetry-unique atoms need be given, the others being generated automatically.

9.4 Writing Gaussian, XMol or MOLDEN input (PUT)

The PUT command may be used at any point in the input to print, or write to a file, the current geometry. The syntax is

$PUT,style,file,status,info$

If *style* is GAUSSIAN, a complete Gaussian input file will be written; in that case, *info* will be used for the first (route) data line, and defaults to '# SP'.

If *style* is XYZ, an XYZ file will be written (see also section 9.3.2). If *style* is CRD, the coordinates will be written in CHARMm CRD format.

If *style* is MOLDEN, an interface file for the MOLDEN visualization program is created; further details and examples are given below.

If *style* is omitted, the Z-matrix, current geometry, and, where applicable, gradient are written.

file specifies a file name to which the data is written; if blank, the the data is written to the output stream. If *status* is omitted or set to NEW, any old contents of the file are destroyed; otherwise the file is appended.

9.4.1 Visualization of results using Molden

Geometry, molecular orbital, and normal mode information, when available, is dumped by PUT, MOLDEN in the format that is usable by MOLDEN.

The interface to the gOpenMol program offers an alternative visualization possibility, and is described in section 25.7.

The example below generates all the information required to plot the molecular orbitals of water, and to visualize the normal modes of vibration:

```
***,H2O
geometry={angstrom;o;h,o,roh;h,o,roh,h,theta};
roh=1.0
theta=104.0
rhf;
optg;
frequencies;
print,low,img;
put,molden,h2o.molden;
```

examples/
h2o_put_molden.com

The example below does a difference density by presenting its natural orbitals to MOLDEN. Note that it although MOLDEN has internal features for difference density plots, the approach show here is more general in that it bypasses the restriction to STO-3G, 3-21G, 4-31G and 6-31G basis sets.

```
gprint,orbitals
geometry={y;planexz;O;H1,O,r;h2,O,r,h1,alpha}
r=1.8
alpha=104
int;
hf;wf,10,1;orbital,2100.2
multi;wf,10,1;orbital,2140.2;

matrop
load,dscf,density,2100.2      !load scf density
load,dmcsf,density,2140.2    !load mcscf density
add,ddiff,dmcsf,-1,dscf     !compute dmcsf-dscf
natorb,neworb1,dscf
natorb,neworb2,dmcsf
natorb,neworbs,ddiff
save,neworbs,2110.2
save,ddiff,2110.2

put,molden,h2o_ddens.molden;orb,2110.2
```

examples/
h2o_diffden_molden.com

9.5 Geometry Files

Using the format

```
GEOMETRY=file
```

the geometry definitions are read from *file*, instead of inline. This file must contain all information of the symmetry block, i.e. symmetry specifications (optional), z-matrix, or xyz-input.

9.6 Lattice of point charges

```
LAT,x,y,z,q;
```

Immediately following the Z-matrix specification may appear one or more “LAT” cards, defining a lattice of point charges. *x,y,z* are the cartesian coordinates and *q* the charge of the point. Note that internal coordinates may not be used, and that only the symmetry unique lattice points should be defined.

9.7 Redefining and printing atomic masses

The current masses of all atoms can be printed using

```
MASS, PRINT
```

The atomic masses can be redefined using

```
MASS, [type,] [symbol=mass, ...]
```

The optional keyword *type* can take either the value AVER[AGE] for using average isotope masses, or ISO[TOPE] for using the masses of the most abundant isotopes. This affects only the rotational constants and vibrational frequencies. As in most quantum chemistry packages, the default for *type* is AVERAGE.

Individual masses can be changed by the following entries, where *symbol* is the chemical symbol of the atom and *mass* is the associated mass. Several entries can be given on one MASS card, and/or several MASS cards can immediately follow each other. If there is any other input between two sets of mass cards, all mass definitions but not `type`) from the first one are overwritten by the last one.

Note that specifying different isotope masses for symmetry related atoms lowers the symmetry of the system if the molecular centre of mass is taken as the origin. This effect can be avoided by using the charge centre as origin, i.e., specifying CHARGE as first entry in the GEOMETRY input:

```
GEOMETRY={CHARGE; ...}
```

9.8 Dummy centres

```
DUMMY,atom1,atom2,...
```

Sets nuclear charges on atoms 1,2 etc. to zero, for doing counterpoise calculations, for example. *atom1, atom2, ...* can be atom numbers or tag names. Note that the current setting of dummies is remembered by the program across restarts via the MOLPRO variable DUMMYATOMS. Dummies can be reset to their original charges using a DUMMY card with no entries.

The program does not recognize automatically if the symmetry is reduced by defining dummy atoms. Therefore, for a given dummy atom, either all symmetry equivalent atoms must also be dummies, or the symmetry must be reduced manually as required. An error will result if the symmetry is not consistent with the dummy center definitions.

9.8.1 Counterpoise calculations

Counterpoise corrections are easily performed using dummy cards. One first computes the energy of the total system, and then for the subsystems using dummy cards.

9.8.2 Example: interaction energy of OH-Ar

```

***,OH(2Sig+)-Ar linear
memory,2,m
geometry={q1;                                !dummy center in center of mass
o,q1,ro;h,q1,rh,o,180;                       !geometry of OH
ar,q1,rar,o,theta,h,0}                       !geometry of Ar
roh=1.8                                       !OH bond-length
rar=7.5                                       !distance of Ar from center of mass
theta=0                                       !angle OH-Ar
ro=roh*16/17                                 !distance of O from center of mass
rh=roh*1/17                                  !distance of H from center of mass
basis=avdz                                   !basis set

text,calculation for complex
rhf;occ,8,3,3;wf,27,1,1                     !RHF for total system
rccsd(t)                                     !CCSD(T) for total system
e_ohar=energy                                !save energy in variable e_ohar

text,cp calculation for OH
dummy,ar                                     !make Ar a dummy center
rhf;occ,3,1,1;wf,9,1,1                     !RHF for OH
rccsd(t)                                     !CCSD(T) for OH
e_oh=energy                                  !save energy in variable e_oh

text,cp calculation for Ar
dummy,o,h                                    !make OH dummy
hf                                           !scf for Ar
ccsd(t)                                     !CCSD(T) for Ar
e_ar=energy                                  !save energy in variable e_ar

text,separate calculation for OH
geometry={O;H,O,roh}                         !geometry for OH alone
rhf;occ,3,1,1;wf,9,1,1                     !RHF for OH
rccsd(t)                                     !CCSD(T) for OH
e_oh_inf=energy                              !save energy in variable e_oh_inf

text,separate calculation for Ar
geometry={AR}                                !geometry for OH alone
hf                                           !scf for Ar
ccsd(t)                                     !CCSD(T) for Ar
e_ar_inf=energy                              !save energy in variable e_ar_inf

de=(e_ohar-e_oh_inf-e_ar_inf)*to cm         !compute uncorrected interaction energy
de_cp=(e_ohar-e_oh-e_ar)*to cm              !compute counter-poise corrected interaction energy
bsse_oh=(e_oh-e_oh_inf)*to cm               !BSSE for OH
bsse_ar=(e_ar-e_ar_inf)*to cm               !BSSE for Ar
bsse_tot=bsse_oh+bsse_ar                    !total BSSE

```

examples/
ohar_bsse.com

For performing counterpoise corrected geometry optimizations see section 32.2.19.

10 BASIS INPUT

The basis set may either be taken from the program library, or may be specified explicitly, or any combination. Optionally, the basis function type can be chosen using the CARTESIAN or SPHERICAL commands.

10.1 Cartesian and spherical harmonic basis functions

MOLPRO uses spherical harmonics (5d, 7f, etc) by default, even for Pople basis sets like 6-31G**. This behaviour may be different to that of other programs; However, cartesian

functions can be requested using the CARTESIAN command.

CARTESIAN

If this command is encountered, the logical MOLPRO variable CARTESIAN is set to true (1.0), and all subsequent calculations use cartesian basis functions. This is remembered across restarts. One can switch back to spherical harmonics using the command

SPHERICAL

10.2 The basis set library

The basis set library consists of a set of plain text files, together with an associated index, that constitute a database of commonly-used basis sets (primitive gaussians and associated contractions) and effective core potentials. These files can be found in the source tree as `lib/*.libmol` and `lib/libmol.index`, but it is usually more convenient to query the database using one of the provided tools.

Many of the basis sets are taken directly from the Pacific Northwest National Laboratory basis set database, but there are others, notably the Stuttgart effective core potentials and bases.

A simple command-line interface to the database is provided through the `libmol` program. It requires the environment variable `LIBMOL` to point to the `lib/` directory, but this will default to the location of the source tree at compile time, so it is often not necessary to specify it. The command-line syntax is

```
libmol [-p print] [-e element] [-k key] [-t type] [-f format]
```

where the parameters are

<i>print</i> :	Output level; 0 means list matching keys, 1 means print also the entry.
<i>element</i> :	Specify chemical element. If omitted, all elements are searched.
<i>key</i> :	Specify record key. If omitted, all keys are searched.
<i>type</i> :	Specify entry type, i.e. <i>s</i> , <i>p</i> , ... If omitted, all types are searched.
<i>format</i> :	One of <code>text</code> (default), <code>molpro</code> (MOLPRO input format), <code>table</code> (tabular) or <code>html</code> (html table) to govern the output format.

A more convenient way of browsing the basis library is through a web-based interface. The CGI script `molpro_basis` presents a graphical and forms based interface for performing searches. It may be installed locally, but is also normally available at

http://www.molpro.net/current/molpro_basis.

10.3 Default basis sets

If a basis is not specified at all for any unique atom group, then the program assumes a global default. Presently, this default is VDZ, but may be overridden using

`BASIS,basis`

or

`BASIS=basis`

basis is looked up in the file `lib/defbas`, which generates an appropriate request for a complete contracted set, together in some cases with an ECP, from the library. This mapping includes the following commonly-used basis sets.

- All of the Dunning correlation-consistent sets, through the use of either the standard name of the basis set (e.g., `aug-cc-pVDZ`) or an abbreviation (e.g., `AVDZ`).
- The older segmented Dunning/Hay double-zeta sets for the first row (`DZ` and `DZP`).
- The Roos ANO basis sets (`ROOS`).
- The Stuttgart ECPs and associated basis sets (e.g., `ECP10MWB`).
- The Hay ECPs and corresponding basis sets (`ECP1` and `ECP2`).
- Some of the Karlsruhe basis sets (`SV`, `TZV`, and, for some elements, `SVP`, `TZVP`, `TZVPP`, `TZVPPP`).
- The Binning/Curtiss sets for Ga–Kr (`BINNING-SV`, `BINNING-SVP`, `BINNING-VTZ` and `BINNING-VTZP`).
- Most of the Pople basis sets, using their standard names (e.g., `6-31G*`, `6-311++G(D,P)`, etc.). Note that specially in this case, the mechanism described below using parenthesized modifiers to restrict the basis set is disabled to allow the full range of standard basis sets to be specified.

Example:

```
BASIS=VTZ
```

generates valence triple zeta basis set for all atoms. Thus, the input

```
***,h2o cc-pVTZ basis      !A title
r=1.85,theta=104          !set geometry parameters
geometry={O;              !z-matrix geometry input
      H1,O,r;
      H2,O,r,H1,theta}
basis=VTZ                 !use VTZ basis
hf                        !closed-shell scf
```

examples/
h2o_scf_vtz.com

is entirely equivalent to

```
***,h2o cc-pVTZ basis      !A title
r=1.85,theta=104          !set geometry parameters
geometry={O;              !z-matrix geometry input
      H1,O,r;
      H2,O,r,H1,theta}
basis={
  spdf,o,vtz;c;
  spd,h,vtz;c}
hf;
```

examples/
h2o_scf_vtz_explicit.com

Default basis sets can be defined anywhere in the input before the energy calculation to which it should apply using a single `BASIS` cards. The default basis set applies to all types of atoms but can be superseded by different basis sets for specific atoms as explained later. Some restrictions concerning the maximum angular momentum functions to be used, or the number of contracted functions are possible as follows:

The maximum angular momentum in the basis set can be reduced using syntax such as

```
BASIS , VQZ ( D )
```

which would omit the f and g functions that would normally be present in the VQZ basis set.

```
BASIS , VQZ ( D / P )
```

would specify additionally a maximum angular momentum of 1 on hydrogen, i.e. would omit d orbitals on hydrogen.

For generally contracted basis sets, an extended syntax can be used to explicitly give the number of contracted functions of each angular momentum. For example,

```
BASIS , ROOS ( 3s2p1d / 2s )
```

generates a 6-31G*-sized basis set from the Roos ANO compilation.

10.4 Default basis sets for individual atoms

More specific basis set definitions for individual atoms can be given BASIS input blocks, which have the following general form:

```
BASIS
SET=type                ! type can be ORBITAL, DENSITY or any other name,
                        ! as used in basis specifications for density
                        ! fitting; optional; default=ORBITAL
DEFAULT=name            ! sets the default basis to name ;
atom1=name1            ! Use basis name1 for atom1
atom2=name2            ! Use basis name1 for atom2
primitive basis set specifications !additional basis functions
SET=type                ! specify basis of another type in following lines
...
END
```

Any number of basis sets can be given in a basis block.

The default and atom specifications can also be merged to one line, separated by commas:

```
DEFAULT=name,atom1=name1,atom2=name2
```

Here the basis sets $name1$, $name2$ overwrite the default basis set $name$ for specific atoms $atom1$, $atom2$, respectively. For instance,

```
DEFAULT=VTZ , O=AVTZ , H=VDZ
```

uses VTZ as the default basis sets, but sets the basis for oxygen to AVTZ and for hydrogen to VDZ

This name conventions for the atom specific basis sets work exactly as described above for default basis sets. The keyword DEFAULT can be abbreviated by DEF. Any DEFAULT basis set defined in a basis set block supercedes a previous one given outside the basis block.

The specifications SET, DEFAULT, atom=name are all optional. If DEFAULT is not given, the previous default, as specified on the last previous BASIS card, is used.

If no further primitive basis set specifications follow, one can also use the one-line form

```
BASIS , DEFAULT=VTZ , O=AVTZ , H=VDZ
```

or

```
BASIS=VTZ , O=AVTZ , H=VDZ
```

Both of these are equivalent to

```
BASIS
DEFAULT=VTZ
O=AVTZ
H=VDZ
END
```

Note that any new BASIS card supercedes all previous basis input, except for the default basis (unless this is given).

The optional additional primitive basis set specifications (see next section) are appended to the given atom-specific basis sets, i.e., the union of atom-specific and primitive basis set definitions is used for the atom.

Examples:

```
BASIS
DEFAULT=VTZ           ! use cc-pVTZ basis as default
H=VDZ                 ! use cc-pVDZ for H-atoms
END
```

This could also be written as

```
BASIS={DEF=VTZ , H=VDZ }
```

```
BASIS
DEFAULT=VTZ           ! use cc-pVTZ basis as default
H=VDZ                 ! use cc-pVDZ for H-atoms
D,H,VTZ              ! add the VTZ d-function to the VDZ basis for H
END
```

```
BASIS
SPD,O,VTZ            !use uncontracted s,p,d functions of basis VTZ for oxygen
S,H,H07              !use Huzinaga 7s for Hydrogen
C,1.4                !contract first four s-functions
P,H,1.0,0.3          !add two p-functions for hydrogen
END
```

Several BASIS cards and/or blocks can immediately follow each other. Always the last specification for a given atom and type is used. Defaults given using BASIS commands can be overwritten by specifications in the integral input. If an individual basis function type is specified for an atom, it is required that all other types are also defined. For example, in the above example, no *f*-functions are included for O, even if the global default would include *f*-functions. Also, defining the *s* functions for hydrogen switches off the default basis set for hydrogen, and so the *p* functions must be defined. Instead of the atomic symbol, the atom group number can also be used.

The same input forms are also possible as direct input to the integral program. In contrast to MOLPRO92, now the atomic symbol can be used in field 2 of a basis specification instead of the atom group number:

```
SPD , O , VTZ           !use VTZ basis for all oxygen atoms
```

```
SPD, 1, V TZ           !use V TZ basis for atom group 1
```

Instead of the BASIS ... END block one can also use the structure BASIS [=] { . . . }

If a basis is not specified at all for any unique atom group, then the program assumes a default. For further details, including respecifying the default to be used, see the specification of the BASIS subcommand below.

10.5 Primitive set definition

A group of basis functions is defined by a data card specifying a set of primitive gaussians, optionally followed by one or more cards specifying particular contractions of primitives to be included in the final basis. When all contraction definitions have been read (delimited by the next data card other than a contraction definition), the remaining primitives in the set which have not been included in any contraction set are added uncontracted to the basis set.

There are four different input forms, as explained below under a) to d):

a) *type,atom,key,scale,nprim*;

Load basis named *key* from the library with angular symmetry *type* (S, P, D, F, G, H, or I). This basis is added from all atoms with number $i=atom$ on the A cards (see above). If *scale* is present, all exponents are scaled by $scale^{**2}$. If *nprim* is specified, the first *nprim* exponents only are taken from the library. If *nprim* is negative, the last $|nprim|$ basis functions from the library set are deleted. Associated with the library basis may be a set of default contraction coefficients which may be accessed in subsequent contraction cards. *type* can include several types, e.g., SPD or DF. This usually makes sense only with default contractions, i.e., such cards should be followed only by "C" without any other specifications for contractions.

b) *type,atom,exp1,exp2,...expn;expn+1,...*;

General specification of exponents; continuation onto subsequent cards (separated by semi-colon) is permitted as shown (the first card can hold up to 19 exponents (cray 13), each following card 20 exponents (cray 15)).

The exponents (and other numerical parameters described below such as numbers of functions, and contraction coefficients) can be given as general input expressions, possibly involving variables. It is important to note, however, that these expressions are evaluated typically just once, at the same time as the complete basis set is parsed. This generally happens the first time that the basis set is required, perhaps before the first SCF calculation can be done. If the variables on which the basis depends are altered, this will not be noticed by the program, and the new basis set will not be used for subsequent stages of the computation. If, however, a new basis block is presented in the input, then the program marks as outdated any quantities such as integrals that have been calculated with the old basis set; subsequent job steps will then use the new basis. Such behaviour can be forced at any stage by issuing the command

```
INT
```

and this is the recommended way of ensuring that a new basis set is adopted; note, however, that INT will cause atomic orbital integrals to be evaluated, and this overhead should of course not be incurred unnecessarily in non-direct calculations.

c) *type,atom,EVEN,nprim,ratio,centre,dratio*

Generates a generalized even tempered set of functions. The number of functions n is specified by *nprim*, their geometric mean c by *centre*, the mean ratio of successive exponents r by *ratio*,

and the variation of this ratio, d , by d ratio. If $centre$ is not given, the previous basis of the same type is extended by diffuse functions. If in this case $ratio$ is not given, r is determined from the exponents of the last two function of the previous basis. If this is not possible, the default $r = 2.5$ is adopted. $d = 1$ (the default) specifies a true even-tempered set, but otherwise the ratio between successive exponents changes linearly; the exponents are given explicitly by

$$\log e_i = \log c + ((n+1)/2 - i) \log r + \frac{1}{2}((n+1)/2 - i)^2 \log d \quad i = 1, 2, \dots, n$$

Example 1 SP, 1, VTZ; C; SP, 1, EVEN, 1;
generates the generally contracted s and p triple-zeta basis sets for atom 1 and extends these by one diffuse function.

Example 2 SPD, 1, VTZ,, -1; C; SP, 1, EVEN, 2, 2.5;
generates the generally contracted s , p triple-zeta basis sets for atom 1. Two energy optimized d -functions of Dunning are included. The last s and p functions are deleted and replaced by two even tempered functions with ratio 2.5.

d) *type,atom,EVENR,nprim,aa,ap,bb,bp*

Generates an even tempered set of $nprim$ functions according to the “regular” prescription described in M W Schmidt and K Ruedenberg, J. Chem. Phys. 71 (1970) 3951. If any of the parameters aa , ap , bb , bp is zero or omitted, the values are taken from table III of the above.

10.6 Contracted set definitions

a) *C,first.last,c1,c2,...,cn;cn+1,...;*

General specification of a contracted function. *first.last* defines the range of primitives to be contracted. $c1$, $c2$... are the contraction coefficients. Continuation onto a subsequent card is permitted as shown.

b) *C;*

Use default contractions from the library. This applies to both the number of contracted primitives and also to the number of different contraction sets.

c) *nC,first.last;*

n contracted functions taken from library. *first.last* defines the range of primitives to be contracted. If n is omitted and *first.last* is specified, $n = 1$. If *first.last* is omitted, the library default values are used. If both n and *first.last* are omitted, default values for both are used.

d) *nC,first.last,record.file,orb.sym;*

n contracted functions taken from orbitals orb , $orb + 1$, ..., $orb + n - 1$ of symmetry sym on molpro file *record.file*. The first nonzero coefficient in the specified orbital corresponds to the first associated basis function. *first.last* specifies the range of primitives to be contracted. If *first.last* is omitted, all coefficients from the specified orbitals are used.

Example 2C, 1.12, 2100.2, 1.1
generates two contractions, using the first 12 coefficients from orbitals 1.1 and 2.1. The orbitals are read from record 2100.2.

10.7 Examples

This shows the use of default basis sets for H₂O:

```
***,H2O
basis=VQZ(f/p)
R=0.95 ANG,THETA=104 DEGREE
geometry={O;H1,O,R;H2,O,R,H1,THETA}
hf                !do closed-shell SCF
```

examples/
h2o_vqz_fp.com

This is equivalent to the explicit input form

```
***,H2O
R=0.95 ANG,THETA=104 DEGREE
geometry={O;H1,O,R;H2,O,R,H1,THETA}
basis={spdf,o,vqz;c;sp,h,vqz,c;}
hf                !do closed-shell SCF
```

examples/
h2o_vqz_fp_explicit.com

11 EFFECTIVE CORE POTENTIALS

Pseudopotentials (effective core potentials, ECPs) may be defined at the beginning of BASIS blocks.

The general form of the input cards is

```
ECP,atom,[ECP specification]
```

which defines a pseudopotential for an atom specified either by a chemical symbol or a group number. The *ECP specification* may consist either of a single keyword, which references a pseudopotential stored in the library, or else of an explicit definition (extending over several input cards), cf. below.

11.1 Input from ECP library

The basis set library presently contains the pseudopotentials and associated valence basis sets by a) the Los Alamos group (P. J. Hay and W. R. Wadt, J. Chem. Phys. **82**, 270 (1985) and following two papers), and b) the Stuttgart/Köln group (e.g., A. Nicklass, M. Dolg, H. Stoll and H. Preuß, J. Chem. Phys. **102**, 8942 (1995); for more details and proper references, see the web page <http://www.theochem.uni-stuttgart.de/pseudopotentials/>). Pseudopotentials a) are adjusted to orbital energies and densities of a suitable atomic reference state, while pseudopotentials b) are generated using total valence energies of a multitude of atomic states.

Library keywords in case a) are ECP1 and ECP2; ECP2 is used when more than one pseudopotential is available for a given atom and then denotes the ECP with the smaller core definition. (For Cu, e.g., ECP1 refers to an Ar-like $18e^-$ -core, while ECP2 simulates a Ne-like $10e^-$ one with the $3s$ and $3p$ electrons promoted to the valence shell). For accurate calculations including electron correlation, promotion of all core orbitals with main quantum number equal to any of the valence orbitals is recommended.

Library keywords in case b) are of the form ECP nXY ; n is the number of core electrons which are replaced by the pseudopotential, X denotes the reference system used for generating the pseudopotential ($X = S$: single-valence-electron ion; $X = M$: neutral atom), and Y stands for

the theoretical level of the reference data ($Y = HF$: Hartree-Fock, $Y = WB$: quasi-relativistic; $Y = DF$: relativistic). For one- or two-valence electron atoms $X = S$, $Y = DF$ is a good choice, while otherwise $X = M$, $Y = WB$ (or $Y = DF$) is recommended. (For light atoms, or for the discussion of relativistic effects, the corresponding $Y = HF$ pseudopotentials may be useful.) Additionally, spin-orbit (SO) potentials and core-polarization potentials (CPP) are available, to be used in connection with case b) ECPs, but these are not currently contained in the library, so explicit input is necessary here (cf. below).

In both cases, a) and b), the same keywords refer to the pseudopotential and the corresponding basis set, with a prefix MBS-... in case a).

11.2 Explicit input for ECPs

For each of the pseudopotentials the following information has to be provided:

- a card of the form
ECP,*atom*,*n_{core}*,*l_{max}*,*l'_{max}*;
where *n_{core}* is the number of core electrons replaced by the pseudopotential V_{ps} , *l_{max}* is the number of semi-local terms in the scalar-relativistic part of V_{ps} , while *l'_{max}* is the corresponding number of terms in the SO part:

$$V_{ps} = -\frac{Z - n_{core}}{r} + V_{l_{max}} + \sum_{l=0}^{l_{max}-1} (V_l - V_{l_{max}}) \mathcal{P}_l + \sum_{l=1}^{l'_{max}} \Delta V_l \mathcal{P}_l \vec{l} \cdot \vec{s} \mathcal{P}_l;$$

the semi-local terms (with angular-momentum projectors \mathcal{P}_l) are supplemented by a local term for $l = l_{max}$.

- a number of cards specifying $V_{l_{max}}$, the first giving the expansion length $n_{l_{max}}$ in

$$V_{l_{max}} = \sum_{j=1}^{n_{l_{max}}} c_j r^{m_j-2} e^{-\gamma_j r^2}$$

and the following $n_{l_{max}}$ ones giving the parameters in the form

$$m_1, \gamma_1, c_1; m_2, \gamma_2, c_2; \dots$$

- a number of cards specifying the scalar-relativistic semi-local terms in the order $l = 0, 1, \dots, l_{max} - 1$. For each of these terms a card with the expansion length n_l in

$$V_l - V_{l_{max}} = \sum_{j=1}^{n_l} c_j^l r^{m_j^l-2} e^{-\gamma_j^l r^2}$$

has to be given, and immediately following n_l cards with the corresponding parameters in the form $m_1^l, \gamma_1^l, c_1^l; m_2^l, \gamma_2^l, c_2^l; \dots$

- analogously, a number of cards specifying the coefficients of the radial potentials ΔV_l of the SO part of V_{ps} .

11.3 Example for explicit ECP input

```

***,CU
! SCF d10s1 -> d9s2 excitation energy of the Cu atom
! using the relativistic Ne-core pseudopotential
! and basis of the Stuttgart/Koeln group.
gprint,basis,orbitals
geometry={cu}
basis
ECP,1,10,3; ! ECP input
 1; ! NO LOCAL POTENTIAL
 2,1.,0.;
 2; ! S POTENTIAL
 2,30.22,355.770158;2,13.19,70.865357;
 2; ! P POTENTIAL
 2,33.13,233.891976;2,13.22,53.947299;
 2; ! D POTENTIAL
 2,38.42,-31.272165;2,13.26,-2.741104;
! (8s7p6d)/[6s5p3d] BASIS SET
s,1,27.69632,13.50535,8.815355,2.380805,.952616,.112662,.040486,.01;
c,1,3,.231132,-.656811,-.545875;
p,1,93.504327,16.285464,5.994236,2.536875,.897934,.131729,.030878;
c,1,2,.022829,-1.009513;C,3,4,.24645,.792024;
d,1,41.225006,12.34325,4.20192,1.379825,.383453,.1;
c,1,4,.044694,.212106,.453423,.533465;
end
rhf;
e1=energy
rhf;occ,4,1,1,1,1,1;closed,4,1,1,1,1,1;wf,19,7,1;
e2=energy
de=(e2-e1)*toev ! Delta E = -0.075 eV

```

examples/
cu_ecp_explicit.com

11.4 Example for ECP input from library

```

***,AuH
! CCSD(T) binding energy of the AuH molecule at r(exp)
! using the scalar-relativistic 19-valence-electron
! pseudopotential of the Stuttgart/Koeln group
gprint,basis,orbitals;
geometry={au}
basis={
ecp,au,ECP60MWB; ! ECP input
spd,au,ECP60MWB;c,1,2; ! basis set
f,au,1.41,0.47,0.15;
g,au,1.2,0.4;
spd,h,avtz;c;
}
rhf;
rccsd(t);core,1,1,1,,1;
e1=energy
geometry={h}
rhf
e2=energy;
rAuH=1.524 ang ! molecular calculation
geometry={au;h,au,rAuH}
hf;
ccsd(t);core,2,1,1;
e3=energy
de=(e3-e2-e1)*toev ! binding energy = 3.11 eV

```

examples/
auh_ecp_lib.com

12 CORE POLARIZATION POTENTIALS

12.1 Input options

The calculation of core-polarization matrix elements is invoked by the CPP card, which can be called at an arbitrary position in the MOLPRO input, provided the integrals have been calculated before. The CPP card can have the following three formats:

- CPP,INIT,*ncenters*;
- CPP,ADD[,*factor*];
- CPP,SET[,*fcpp*];

CPP,INIT,< *ncenters* >;

abs(< *ncenters* >) further cards will be read in the following format:

< *atomtype* >, < *ntype* >, < α_d >, < α_q >, < β_d >, < *cutoff* >;

< *atomtype* > corresponds to the recognition of the atomic centres in the integral part of the program,

< *ntype* > fixes the form of the cutoff-function (choose 1 for Stoll/Fuentealba and 2 for Mueller/Meyer);

< α_d > is the static dipole polarizability,

< α_q > is the static quadrupole polarizability,

< β_d > is the first non-adiabatic correction to the dipole-polarizability and

< *cutoff* > is the exponential parameter of the cutoff-function.

When < *ncenters* > is lower than zero, only the integrals are calculated and saved in the record 1490.1. Otherwise, the h_0 matrix (records 1200.1 and 1210.1) and the two-electron-integrals (record 1300.1) will be modified.

CPP,ADD,< *factor* >;

With this variant, previously calculated matrix elements of the polarization matrix can be added with the variable factor < *factor* > (default: < *factor* > = 1) to the h_0 -matrix as well as to the two-electron-integrals. In particular, CPP,ADD,-1.; can be used to retrieve the integrals without the polarization contribution.

CPP,SET,< *fcpp* >;

normally not necessary but may be used to tell MOLPRO after a restart, with what factor the polarization integrals are effective at the moment.

12.2 Example for ECP/ CPP

```

***,Na2
! Potential curve of the Na2 molecule
! using 1-ve ECP + CPP
gprint,basis,orbitals;
rvec=[2.9,3.0,3.1,3.2,3.3] ang
do i=1,#rvec
rNa2=rvec(i)
geometry={na;na,na,rNa2}
basis={
ecp,na,ecp10sdf;          ! ecp input
s,na,even,8,3,.5;        ! basis input
p,na,even,6,3,.2;
d,na,.12,.03;
}
cpp,init,1;              ! CPP input
na,1,.9947,,,.62;
hf;
ehf(i)=energy
cisd;core;
eci(i)=energy
enddo
table,rvec,ehf,eci
---
```

examples/
na2_ecp_cpp.com

13 RELATIVISTIC CORRECTIONS

There are three ways in MOLPRO to take into account scalar relativistic effects:

1. Use the Douglas-Kroll relativistic one-electron integrals.
2. Compute a perturbational correction using the Cowan-Griffin operator (see section 4.13).
3. Use relativistic effective core potentials (see section 11).

For all-electron calculations, the preferred way is to use the Douglas-Kroll Hamiltonian. It is simply activated by setting

```
DKROLL=1
```

somewhere in the input before the first energy calculation.

13.0.1 Example for computing relativistic corrections

```

***,ar2
geometry={ar1;ar2,ar1,r} !geometry definition
r=2.5 ang                !bond distance
hf;                      !non-relativistic scf calculation
expec,rel,darwin,massv   !compute relativistic correction using Cowan-Griffin operator
e_nrel=energy            !save non-relativistic energy in variable enrel
show,massv,darwin,erel  !show individual contribution and their sum

dkroll=1                 !use douglas-kroll one-electron integrals
hf;                      !relativistic scf calculation
e_dk=energy              !save relativistic scf energy in variable e_dk.
show,massv,darwin,erel  !show mass-velocity and darwin contributions and their sum
show,e_dk-e_nrel        !show relativistic correction using Douglas-Kroll
```

examples/
ar2_rel.com

14 THE SCF PROGRAM

The Hartree-Fock self-consistent field program is invoked by one of the following commands:

HF or RHF	calls the spin-restricted Hartree-Fock program
UHF or UHF-SCF	calls the spin-unrestricted Hartree-Fock program

In contrast to older versions of MOLPRO, the HF and RHF directives have identical functionality and can both be used for closed-shell or open-shell calculations. Other aliases are HF-SCF or RHF-SCF.

Often, no further input is necessary. By default, the number of electrons is equal to the nuclear charge, the wavefunction is assumed to be totally symmetric (symmetry 1), and the spin multiplicity is 1 (singlet) for an even number of electrons and 2 (doublet) otherwise. The Aufbau principle is used to determine the occupation numbers in each symmetry. Normally, this works well in closed-shell cases, but sometimes wrong occupations are obtained or the wavefunction alternates between different orbital spaces. In such cases, the OCC directive must be used to force convergence to the desired state.

In open-shell cases, we recommend to use the WF, OCC, CLOSED, or OPEN cards to define the wavefunction uniquely. Other commands frequently used are START and ORBITAL (or SAVE) to modify the default records for starting and optimized orbitals, respectively. The SHIFT directive allows to modify the level shift in the RHF program, and EXPECT to calculate expectation values of one-electron operators (see section 4.13).

14.1 Defining the wavefunction

The number of electrons and the total symmetry of the wavefunction are specified on the WF card:

WF,*elec,sym,spin*

where

<i>elec</i>	is the number of electrons
<i>sym</i>	is the number of the irreducible representation
<i>spin</i>	defines the spin symmetry, $spin = 2 * S$ (singlet=0, doublet=1, triplet=2 etc.)

Note that these values take sensible defaults if any or all are not specified (see section 2.13).

14.1.1 Defining the number of occupied orbitals in each symmetry

OCC,*n₁,n₂,...,n₈*

To avoid convergence problems in cases with high symmetry, this card should be included whenever the occupation pattern is known in advance. n_i is the number of occupied orbitals in the irreducible representation i . The total number of orbitals must be equal to $(elec+spin)/2$ (see WF card).

14.1.2 Specifying closed-shell orbitals

CLOSED, n_1, n_2, \dots, n_8

This optional card can be used in open-shell calculations to specify the number of closed-shell orbitals in each symmetry. This makes possible to force specific states in the absence of an OPEN card.

14.1.3 Specifying open-shell orbitals

OPEN, $orb_1.sym_1, orb_2.sym_2, \dots, orb_n.sym_n$

This optional card can be used to specify the singly occupied orbitals. The number of singly occupied orbitals must be equal to $spin$, and their symmetry product must be equal to sym (see WF card). If the OPEN card is not present, the open shell orbitals are selected automatically. The algorithm tries to find the ground state, but it might happen that a wrong state is obtained if there are several possibilities for distributing the open shell electrons among the available orbitals. This can also be avoided using the CLOSED card. If $orb_i.sym$ is negative, this orbital will be occupied with negative spin (only allowed in UHF).

14.2 Saving the final orbitals

ORBITAL,*record.file*

SAVE,*record.file*

The optimized orbitals, and the corresponding density matrix, fock matrix, and orbital energies are saved on *record.file*. SAVE is an alias for ORBITAL. If this card is not present, the defaults for *record* are:

RHF	2100	
UHF	2200	(holds both α and β -spin orbitals and related quantities)

These numbers are incremented by one for each subsequent calculation of the same type in the same input. Note that this holds for the sequence number in the input, independently in which order they are executed (see section 2.7).

The default for *file* is 2.

14.3 Starting orbitals

The START directive can be used to specify the initial orbitals used in the SCF iteration. It is either possible to generate an initial orbital guess, or to start with previously optimized orbitals. Alternatively, one can also use a previous density matrix to construct the first fock operator.

If the START card is absent, the program tries to find suitable starting orbitals as follows:

First:	Try to read orbitals from <i>record</i> specified on the ORBITAL or SAVE card or the corresponding default (see ORBITAL). All files are searched.
Second:	Try to find orbitals from a previous SCF or MCSCF calculation. All files are searched.

Third: If no orbitals are found, the starting orbitals are generated using approximate atomic densities or eigenvectors of h (see below).

Since these defaults are usually appropriate, the START card is not required in most cases.

14.3.1 Initial orbital guess

An initial orbital guess can be requested as follows:

START,[TYPE=]*option*

The *option* keyword can be:

H0	Use eigenvectors of h as starting guess.
ATDEN	Use natural orbitals of a diagonal density matrix constructed using atomic occupation numbers.

The atomic density guess works very well with minimal or generally contracted basis sets for which the first contracted basis functions correspond to the atomic $1s$, $2s$, $2p$... orbitals, e.g., Dunning's cc-pVnZ sets, the STO-3G, or the 6-31G bases. For such basis sets ATDEN is used by default. If a segmented basis set with several contractions for each shell is used, ATDEN should not be specified and H0 is used by default. Since eigenvectors of h are often a very poor starting guess, it is recommended to generate the starting orbitals using a small basis like STO-3G (see section 14.3.2 below).

Example:

```
r=1.85,theta=104           !set geometry parameters
geometry={0;               !z-matrix geometry input
  H1,O,r;
  H2,O,r,H1,theta}
basis=STO-3G               !first basis set
hf                         !scf using STO-3G basis
basis=6-311G              !second basis set
hf                         !scf using 6-311G basis set
```

examples/
h2o_sto3gstart1.com

The second calculation uses the optimized orbitals of the STO-3G calculation as starting guess. This is done by default and no START card is necessary. The explicit use of START and SAVE cards is demonstrated in the example in the next section.

The following input is entirely equivalent to the one in the previous section:

```
r=1.85,theta=104           !set geometry parameters
geometry={0;               !z-matrix geometry input
  H1,O,r;
  H2,O,r,H1,theta}
basis=STO-3G               !first basis set
hf                         !scf using STO-3G basis
start,atdens              !use atomic density guess
save,2100.2                !save orbitals to record 2100.2
basis=6-311G              !second basis set
hf                         !scf using 6-311G basis set
start,2100.2              !start with orbitals from the previous STO-3G calculation.
save,2101.2                !save optimized orbitals to record 2101.2
```

examples/
h2o_sto3gstart2.com

14.3.2 Starting with previous orbitals

START,[RECORD=]*record.file*,[*specifications*]

reads previously optimized orbitals from record *record* on file *file*. Optionally, a specific orbital set can be specified as described in section 2.16.

The specified dump record may correspond to a different geometry, basis set, and/or symmetry than used in the present calculation. Using starting orbitals from a different basis set can be useful if no previous orbitals are available and the ATDENS option cannot be used (see above).

The following example shows how to change the symmetry between scf calculations. Of course, this example is quite useless, but sometimes it might be easier first to obtain a solution in higher symmetry and then convert this to lower symmetry for further calculations.

```

r1=1.85,r2=1.85,theta=104           !set geometry parameters
geometry={0;                         !z-matrix geometry input
  H1,O,r1;
  H2,O,r2,H1,theta}
basis=vdz
hf                                   !scf using c2v symmetry
orbital,2100.2                       !save on record 2100.2

set,zsymel=x

hf
start,2100.2                         !start with previous orbitals from c2v symmetry
orbital,2101.2                       !save new orbitals

set,zsymel=[x,y]
hf
start,2101.2                         !start with orbitals from cs symmetry
orbital,2102.2                       save new orbitals

```

examples/
h2o_c2v_cs_start.com

Note, however, that this only works well if the orientation of the molecule does not change. Sometimes it might be helpful to use the `noorient` option.

Note also that a single dump record cannot hold orbitals for different basis dimensions. Using `save=2100.2` in the second calculation would therefore produce an error.

If orbitals from a corresponding SCF calculation at a neighbouring geometry are available, these should be used as starting guess.

14.3.3 Starting with a previous density matrix

START,DENSITY=*record.file*,[*specifications*]

A density matrix is read from the given dump record and used for constructing the first fock matrix. A specific density matrix can be specified as described in section 2.16. It is normally not recommended to use the DENSITY option.

14.4 Rotating pairs of orbitals

ROTATE,*orb₁.sym,orb₂.sym,angle*

Performs a 2×2 rotation of the initial orbitals *orb₁* and *orb₂* in symmetry *sym* by *angle* degrees. With *angle*= 0 the orbitals are exchanged. See MERGE for other possibilities to manipulate orbitals. In UHF, only the β -spin orbitals are rotated.

14.5 Using additional point-group symmetry

Since *MOLPRO* can handle only Abelian point-groups, there may be more symmetry than explicitly used. For instance, if linear molecules are treated in C_{2v} instead of $C_{\infty v}$, the $\delta_{(x^2-y^2)}$ -orbitals appear in symmetry 1 (A_1). In other cases, a linear geometry may occur as a special case of calculations in C_S symmetry, and then one component of the π -orbitals occurs in symmetry 1 (A'). The program is able to detect such hidden “extra” symmetries by blockings in the one-electron hamiltonian h and the overlap matrix S . Within each irreducible representation, an “extra” symmetry number is then assigned to each basis function. These numbers are printed at the end of the integral output. Usually, the extra symmetries are ordered with increasing l -quantum number of the basis functions. This information can be used to determine and fix the extra symmetries of the molecular orbitals by means of the *SYM* command.

SYM, irrep, sym(1), sym(2), ..., sym(n)

sym(i) are the extra symmetries for the first n orbitals in the irreducible representation *irrep*. For instance, if you want that in a linear molecule the orbitals 1.1 to 3.1 are σ and 4.1, 5.1 δ , the *SYM* card would read (calculation done with X,Y as symmetry generators):

SYM, 1, 1, 1, 1, 2, 2

If necessary, the program will reorder the orbitals in each iteration to force this occupation. The symmetries of occupied and virtual orbitals may be specified. By default, symmetry contaminations are not removed. If *irrep* is set negative, however, symmetry contaminations are removed. Note that this may prevent convergence if degenerate orbitals are present.

14.6 Expectation values

EXPEC, oper₁, oper₂, ..., oper_n

Calculates expectation values for one-electron operators *oper₁, oper₂, ..., oper_n*. See section 4.13 for the available operators. By default, the dipole moments are computed. Normally, it is recommended to use the *GEXPEC* directive if expectation values for other operators are of interest. See section 4.13 for details.

14.7 Miscellaneous options

All commands described in this section are optional. Appropriate default values are normally used.

14.7.1 Level shifts

SHIFT, shifta, shiftb, nitord, nitcl

A level shift of *shifta* and *shiftb* hartree for α - and β -spin orbitals, respectively, is applied. This can improve convergence, but has no effect on the solution. *shifta* = -0.2 to -0.3 are typical values. The defaults are *shifta* = 0 and *shifta* = -0.3 in closed and open-shell calculations, respectively, and *shiftb* = 0.

In open-shell calculations, the orbitals are reordered after each iteration to obtain maximum overlap with the orbitals from the previous iteration. This takes only effect after *nitord* iterations. The default is *nitord* = *maxit*/4 if no start card is present and *nitord* = 1 if a *START* card is found.

If the iteration count is smaller than *nitcl*, only the closed-shell part of the Fock matrix is used (default *nitcl* = 0).

14.7.2 Maximum number of iterations

MAXIT,*maxit*

sets the maximum number of iterations to *maxit*. The default is *maxit*= 30.

14.7.3 Convergence threshold

ACCU,*accu*

The convergence threshold is set to $10^{**(-accu)}$. This applies to the square sum of the density matrix element changes. The default is *accu*= 10.

14.7.4 Print options

ORBPRINT,*print,test*

This determines the number of virtual orbitals printed at the end of the calculation. By default, *print*= 0, i.e., only the occupied orbitals are printed. *print*= -1 suppresses printing of orbitals entirely. *test*= 1 has the additional effect of printing the orbitals after each iteration.

14.7.5 Interpolation

IPOL,*iptyp,ipnit,ipstep,maxdis*

This command controls DIIS interpolation. *iptyp* can be:

DIIS	direct inversion of the iterative subspace. This is the default and yields mostly fastest convergence.
DM	obsolete. No effect in MOLPRO98
HFM	obsolete. No effect in MOLPRO98
NONE	No interpolation.

ipnit is the number of the iteration in which the interpolation starts. *ipstep* is the iteration increment between interpolations. *maxdis* is the maximum dimension of the DIIS matrix (default 10).

14.7.6 Reorthonormalization of the orbitals

ORTH,*nitort*

In the RHF program, the orbitals are reorthonormalized after every *nitort* iterations. The default is *nitort*= 8.

14.7.7 Direct SCF

DIRECT,*options*

If this card is present, the calculation is done in direct mode. See section 8 for options. Normally, it is recommended to use the global GDIRECT command to request the direct mode. See section 8 for details.

Table 8: Miscellaneous options for the SCF program

Option	description
SHIFTA= <i>shifta</i>	Level shift for α -spin orbitals (see SHIFT card).
SHIFTB= <i>shiftb</i>	Level shift for β -spin orbitals (see SHIFT card).
ACCURA= <i>accu</i>	Convergence threshold (see ACCU card).
ENERGY= <i>thrden</i>	Energy convergence threshold (default depends on <i>accu</i>).
UNOMIN= <i>unomin</i>	Minimum occupation number for UNO-CAS (default 0.02)
UNOMAX= <i>unomax</i>	Maximum occupation number for UNO-CAS (default 1.98)
POTFAC= <i>potfac</i>	Scale factor for potential energy in first iteration (default 1.0)
IPTYP= <i>iptyp</i>	Interpolation type (see IPOL card).
IPNIT= <i>ipnit</i>	First iteration for DIIS interpolation (see IPOL card).
IPSTEP= <i>ipstep</i>	Iteration increment for DIIS interpolation (see IPOL card).
MAXDIS= <i>maxdis</i>	Max number of Fock matrices used in DIIS interpolation (default 10)
NITORD= <i>nitord</i>	Parameter of reordering orbitals (see SHIFT card).
NITCL= <i>nitcl</i>	Parameter for fock matrix (see SHIFT card).
NITORT= <i>nitort</i>	Parameter for orbital orthonormalization (see ORTH card).
MAXIT= <i>maxit</i>	Maximum number of iterations (see MAXIT card).
NDRU= <i>ndru</i>	Number of virtual orbitals printed, see ORBPRINT card).
NPRT= <i>npert</i>	Test print parameter.
JACOBI= <i>jacobi</i>	If nonzero, use Jacobi diagonalization.

14.7.8 Options

The OPTION directive can be used to set various miscellaneous options, including those described before.

OPTION,*option*₁=*value*₁,*option*₂=*value*₂,*option*₃=*value*₃,...

The following table list the available options:

15 THE DENSITY FUNCTIONAL PROGRAM

Density-functional theory calculations may be performed using one of the following commands:

DFT	calculate functional of a previously computed density.
KS or KS-SCF	calls the closed-shell self-consistent Kohn-Sham program
RKS or RKS-SCF	calls the spin-restricted open-shell Kohn-Sham program
UKS or UKS-SCF	calls the spin-unrestricted open-shell Kohn-Sham program

Each of these commands may be qualified with the key-names of the functional(s) which are to be used:

command, key1, key2, key3, ...

If no functional keyname is given, the contents of the MOLPRO vector variable DF is interpreted as a list of functionals; If DF is empty, DFTNAME is examined; otherwise, the default is LDA (see below). Following this command may appear commands specifying options for the density-functional code, followed, in the Kohn-Sham case, by further SCF options exactly as for the Hartree-Fock programs.

On completion of the functional evaluation, or self-consistent Kohn-Sham calculation, the values of the individual functionals are stored in the MOLPRO vector variable DFTFUNS; the total is in DFTFUN, and the corresponding individual functional names in DFTNAME.

Energy gradients are available for self-consistent Kohn-Sham calculations.

15.1 Density Functionals

In the following, ρ_α and ρ_β are the α and β spin densities; the total spin density is ρ ;

The gradients of the density enter through

$$\sigma_{\alpha\alpha} = \nabla\rho_\alpha \cdot \nabla\rho_\alpha, \sigma_{\beta\beta} = \nabla\rho_\beta \cdot \nabla\rho_\beta, \sigma_{\alpha\beta} = \sigma_{\beta\alpha} = \nabla\rho_\alpha \cdot \nabla\rho_\beta, \sigma = \sigma_{\alpha\alpha} + \sigma_{\beta\beta} + 2\sigma_{\alpha\beta}. \quad (1)$$

$$\chi_\alpha = \frac{\sqrt{\sigma_{\alpha\alpha}}}{\rho_\alpha^{4/3}}, \chi_\beta = \frac{\sqrt{\sigma_{\beta\beta}}}{\rho_\beta^{4/3}}. \quad (2)$$

$$\upsilon_\alpha = \nabla^2\rho_\alpha, \upsilon_\beta = \nabla^2\rho_\beta, \upsilon = \upsilon_\alpha + \upsilon_\beta. \quad (3)$$

Additionally, the kinetic energy density for a set of (Kohn-Sham) orbitals generating the density can be introduced through

$$\tau_\alpha = \sum_i^\alpha |\nabla\phi_i|^2, \tau_\beta = \sum_i^\beta |\nabla\phi_i|^2, \tau = \tau_\alpha + \tau_\beta. \quad (4)$$

All of the available functionals are of the general form

$$F[\rho_s, \rho_{\bar{s}}, \sigma_{ss}, \sigma_{\bar{s}\bar{s}}, \sigma_{s\bar{s}}, \tau_s, \tau_{\bar{s}}, \upsilon_s, \upsilon_{\bar{s}}] = \int d^3\mathbf{r} K(\rho_s, \rho_{\bar{s}}, \sigma_{ss}, \sigma_{\bar{s}\bar{s}}, \sigma_{s\bar{s}}, \tau_s, \tau_{\bar{s}}, \upsilon_s, \upsilon_{\bar{s}}) \quad (5)$$

where \bar{s} is the conjugate spin to s .

15.1.1 B86: $X\alpha\beta\gamma$

A. D. Becke, J. Chem. Phys. 84, 4524 (1986)

Divergence free semiempirical gradient-corrected exchange energy functional.

$$K = \sum_s -\frac{c\rho_s^{4/3}(1+\beta\chi_s^2)}{1+\gamma\chi_s^2}, \quad (6)$$

where

$$c = \frac{3}{2} \left(\frac{3}{4\pi} \right)^{1/3}, \quad (7)$$

$$\beta = 0.0076 \quad (8)$$

and

$$\gamma = 0.004. \quad (9)$$

15.1.2 B86MGC: $X\alpha\beta\gamma$ with Modified Gradient Correction

A. D. Becke, J. Chem. Phys. 85, 7184 (1986)

B86 with modified gradient correction for large density gradients.

$$K = \sum_s -c\rho_s^{4/3} - \frac{\beta\chi_s^2\rho_s^{4/3}}{(1+\lambda\chi_s^2)^{4/5}}, \quad (10)$$

where

$$c = \frac{3}{2} \left(\frac{3}{4\pi} \right)^{1/3}, \quad (11)$$

$$\beta = 0.00375 \quad (12)$$

and

$$\lambda = 0.007. \quad (13)$$

15.1.3 B86R: $X\alpha\beta\gamma$ Re-optimised

A. D. Becke, J. Chem. Phys. 107, 8554 (1997)

Re-optimised β of B86 used in part 3 of Becke's 1997 paper.

$$\beta = 0.00787 \quad (14)$$

15.1.4 B88C:

A. D. Becke, J. Chem. Phys. 88, 1053 (1988)

Correlation functional depending on the B86MGC exchange functional with empirical atomic parameters, t and u . B86MGC should be used in conjunction with B88C.

$$K = -0.8\rho_\alpha\rho_\beta q^2 \left(1 - \frac{\ln(1+q)}{q}\right) - 0.01 \sum_s \rho_s dz^4 \left(1 - 2 \frac{\ln(1+1/2z)}{z}\right), \quad (15)$$

where

$$q = t(x_\alpha + x_\beta), \quad (16)$$

$$x_\alpha = 0.5 \left(c\rho_\alpha^{1/3} + \frac{\beta\chi_\alpha^2\rho_\alpha^{1/3}}{(1+\lambda\chi_\alpha^2)^{4/5}} \right)^{-1}, \quad (17)$$

$$x_\beta = 0.5 \left(c\rho_\beta^{1/3} + \frac{\beta\chi_\beta^2\rho_\beta^{1/3}}{(1+\lambda\chi_\beta^2)^{4/5}} \right)^{-1}, \quad (18)$$

$$t = 0.63, \quad (19)$$

$$z = 2ux_s, \quad (20)$$

$$x_s = 0.5\rho_s \left(c\rho_s^{4/3} + \frac{\beta\chi_s^2\rho_s^{4/3}}{(1+\lambda\chi_s^2)^{4/5}} \right)^{-1}, \quad (21)$$

$$u = 0.96, \quad (22)$$

$$d = \tau_s - \frac{\sigma_{ss}}{4\rho_s}, \quad (23)$$

$$c = \frac{3}{2} \left(\frac{3}{4\pi} \right)^{1/3}, \quad (24)$$

$$\beta = 0.00375 \quad (25)$$

and

$$\lambda = 0.007. \quad (26)$$

15.1.5 B88:

A. D. Becke, Phys. Rev. A. 38, 3098 (1988)

$$K = - \sum_s \rho_s^{4/3} \left(c + \frac{\beta\chi_s^2}{1+6\beta\chi_s \operatorname{arcsinh}(\chi_s)} \right), \quad (27)$$

where

$$c = \frac{3}{2} \left(\frac{3}{4\pi} \right)^{1/3} \quad (28)$$

and

$$\beta = 0.0042. \quad (29)$$

15.1.6 B95: Becke's 1995 Correlation Functional

A. D. Becke, J. Chem. Phys. 104, 1040 (1995)

 τ -dependent dynamical correlation functional.

$$K = \frac{E}{1 + l(\chi_\alpha^2 + \chi_\beta^2)} + \sum_s \frac{F\varepsilon(\rho_s, 0)}{H(1 + \nu\chi_s^2)^2}, \quad (30)$$

where

$$E = \varepsilon(\rho_\alpha, \rho_\beta) - \varepsilon(\rho_\alpha, 0) - \varepsilon(\rho_\beta, 0), \quad (31)$$

$$l = 0.0031, \quad (32)$$

$$F = \tau_s - \frac{\sigma_{ss}}{4\rho_s}, \quad (33)$$

$$H = 3(6\pi^2)^{2/3} \rho_s^{5/3} / 5, \quad (34)$$

$$\nu = 0.038 \quad (35)$$

and $\varepsilon(\alpha, \beta)$ is the correlation energy per particle of the Local Spin Density Approximation(PW92C).**15.1.7 B97: B97=B97DF + 0.1943 Exact Exchange**

A. D. Becke, J. Chem. Phys. 107, 8554 (1997)

B97DF is given by

$$K = (\varepsilon(\rho_\alpha, \rho_\beta) - \varepsilon(\rho_\alpha, 0) - \varepsilon(\rho_\beta, 0)) (A_0 + A_1\eta(d, \lambda_1) + A_2\eta^2(d, \lambda_1)) + \sum_s \varepsilon(\rho_s, 0) (B_0 + B_1\eta(\chi_s^2, \lambda_2) + B_2\eta^2(\chi_s^2, \lambda_2)) \quad (36)$$

$$- \frac{3}{2} \left(\frac{3}{4\pi} \right)^{1/3} \rho_s^{4/3} (C_0 + C_1\eta(\chi_s^2, \lambda_3) + C_2\eta^2(\chi_s^2, \lambda_3)), \quad (37)$$

where

$$d = (\chi_\alpha^2 + \chi_\beta^2)/2, \quad (38)$$

$$\eta(\theta, \mu) = \frac{\mu\theta}{1 + \mu\theta}, \quad (39)$$

$$A = [0.9454, 0.7471, -4.5961], \quad (40)$$

$$B = [0.1737, 2.3487, -2.4868], \quad (41)$$

$$C = [0.8094, 0.5073, 0.7481], \quad (42)$$

$$\lambda = [0.006, 0.2, 0.004] \quad (43)$$

and $\varepsilon(\alpha, \beta)$ is the correlation energy per particle of the Local Spin Density Approximation(PW92C).

15.1.8 B97R: B97 Re-parameterized by Hamprecht *et al.* B97R=B97R + 0.21 Exact Exchange

F. A. Hamprecht, A. J. Cohen, D. J. Tozer and N. C. Handy, J. Chem. Phys. 109, 6264 (1998)

Re-parameterization of the B97 functional in a self-consistent procedure by Hamprecht *et al.* B97RDF is given by

$$A = [0.955689, 0.788552, -5.47869], \quad (44)$$

$$B = [0.0820011, 2.71681, -2.87103], \quad (45)$$

$$C = [0.789518, 0.573805, 0.660975] \quad (46)$$

15.1.9 BR: Becke-Roussel Exchange Functional

A. D. Becke and M. R. Roussel, Phys. Rev. A 39, 3761 (1989)

$$K = \frac{1}{2} \sum_s \rho_s U_s, \quad (47)$$

where

$$U_s = -(1 - e^{-x} - xe^{-x}/2)/b, \quad (48)$$

$$b = \frac{x^3 e^{-x}}{8\pi\rho_s} \quad (49)$$

and x is defined by the nonlinear equation

$$\frac{xe^{-2x/3}}{x-2} = \frac{2\pi^{2/3}\rho_s^{5/3}}{3Q_s}, \quad (50)$$

where

$$Q_s = (v_s - 2\gamma D_s)/6, \quad (51)$$

$$D_s = \tau_s - \frac{\sigma_{ss}}{4\rho_s} \quad (52)$$

and

$$\gamma = 1. \quad (53)$$

15.1.10 BRUEG: Becke-Roussel Exchange Functional — Uniform Electron Gas Limit

A. D. Becke and M. R. Roussel, Phys. Rev. A 39, 3761 (1989)

As for BR but with $\gamma = 0.8$.

15.1.11 BW: Becke-Wigner Exchange-Correlation Functional

P. A. Stewart and P. M. W. Gill, J. Chem. Faraday Trans. 273, 183 (1995)

Hybrid exchange-correlation functional comprising Becke's 1998 exchange and Wigner's spin-polarised correlation functionals.

$$K = -4c \frac{\rho_\alpha \rho_\beta}{\rho} \left(1 + \frac{d}{\rho^{1/3}}\right)^{-1} + \sum_s \alpha \rho_s^{4/3} - \frac{\beta \rho_s^{4/3} \chi_s^2}{1 + 6\beta \chi_s \operatorname{arcsinh}(\chi_s)}, \quad (54)$$

where

$$\alpha = -\frac{3}{2} \left(\frac{3}{4\pi} \right)^{1/3}, \quad (55)$$

$$\beta = 0.0042, \quad (56)$$

$$c = 0.04918 \quad (57)$$

and

$$d = 0.349. \quad (58)$$

15.1.12 CS: Colle-Salvetti correlation functional

R. Colle and O. Salvetti, *Theor. Chim. Acta* 37, 329 (1974); C. Lee, W. Yang and R. G. Parr, *Phys. Rev. B* 37, 785(1988)

CS2 is defined through

$$K = -a \left(\frac{\rho + 2b\rho^{-5/3} [\rho_\alpha t_\alpha + \rho_\beta t_\beta - \rho t_W] e^{-c\rho^{-1/3}}}{1 + d\rho^{-1/3}} \right) \quad (59)$$

where

$$t_\alpha = \frac{\tau_\alpha}{2} - \frac{\nu_\alpha}{8} \quad (60)$$

$$t_\beta = \frac{\tau_\beta}{2} - \frac{\nu_\beta}{8} \quad (61)$$

$$t_W = \frac{1}{8} \frac{\sigma}{\rho} - \frac{1}{2} \nu \quad (62)$$

and the constants are $a = 0.04918, b = 0.132, c = 0.2533, d = 0.349$.

CS1 is formally identical to CS2, except for a reformulation in which the terms involving ν are eliminated by integration by parts. This makes the functional more economical to evaluate. In the limit of exact quadrature, CS1 and CS2 are identical, but small numerical differences appear with finite integration grids.

CS is an alias for CS1.

15.1.13 G96: Gill's 1996 Gradient Corrected Exchange Functional

P. M. W. Gill, *Mol. Phys.* 89, 433 (1996)

$$K = \sum_s \rho_s^{4/3} \left(\alpha - \frac{1}{137} \chi_s^{3/2} \right), \quad (63)$$

where

$$\alpha = -\frac{3}{2} \left(\frac{3}{4\pi} \right)^{1/3}. \quad (64)$$

15.1.14 HCTH93:

F. A. Hamprecht, A. J. Cohen, D. J. Tozer and N. C. Handy, *J. Chem. Phys.* 109, 6264 (1998)

The original HCTH functional with parameters optimized on a set of 93 training systems.

$$\begin{aligned}
 K = & (\epsilon(\rho_\alpha, \rho_\beta) - \epsilon(\rho_\alpha, 0) - \epsilon(\rho_\beta, 0)) (A_0 + A_1\eta(d, \lambda_1) + A_2\eta^2(d, \lambda_1) + A_3\eta^3(d, \lambda_1) + A_4\eta^4(d, \lambda_1)) \\
 & + \sum_s \epsilon(\rho_s, 0) (B_0 + B_1\eta(\chi_s^2, \lambda_2) + B_2\eta^2(\chi_s^2, \lambda_2) + B_3\eta^3(\chi_s^2, \lambda_2) + B_4\eta^4(\chi_s^2, \lambda_2)) \\
 & - \frac{3}{2} \left(\frac{3}{4\pi} \right)^{1/3} \rho_s^{4/3} (C_0 + C_1\eta(\chi_s^2, \lambda_3) + C_2\eta^2(\chi_s^2, \lambda_3) + C_3\eta^3(\chi_s^2, \lambda_3) + C_4\eta^4(\chi_s^2, \lambda_3)), \quad (65)
 \end{aligned}$$

where

$$d = (\chi_\alpha^2 + \chi_\beta^2)/2, \quad (66)$$

$$\eta(\theta, \mu) = \frac{\mu\theta}{1 + \mu\theta}, \quad (67)$$

$$A = [0.72997, 3.35287, -11.543, 8.08564, -4.47857], \quad (68)$$

$$B = [0.222601, -0.0338622, -0.012517, -0.802496, 1.55396], \quad (69)$$

$$C = [1.0932, -0.744056, 5.5992, -6.78549, 4.49357], \quad (70)$$

$$\lambda = [0.006, 0.2, 0.004] \quad (71)$$

and $\epsilon(\alpha, \beta)$ is the correlation energy per particle of the Local Spin Density Approximation(PW92C).

15.1.15 HCTH120:

A. D. Boese, N. L. Doltsinis, N. C. Handy and M. Sprik, *J. Chem. Phys.* 112, 1670 (2000)

HTCH optimized on a set of 120 training systems extending the set of HCTH93 to include Anionic atoms and molecules, 2^{nd} row anions and H-bonded dimers.

$$A = [0.51473, 6.9298, -24.707, 23.110, -11.323], \quad (72)$$

$$B = [0.48951, -0.2607, 0.4329, -1.9925, 2.4853], \quad (73)$$

$$C = [1.09163, -0.7472, 5.0783, -4.1075, 1.1717] \quad (74)$$

and

$$\lambda = [0.006, 0.2, 0.004]. \quad (75)$$

15.1.16 HCTH147:

A. D. Boese, N. L. Doltsinis, N. C. Handy and M. Sprik, *J. Chem. Phys.* 112, 1670 (2000)

HTCH optimized on a further extended set of 147 training systems.

$$A = [0.54235, 7.0146, -28.382, 35.033, -20.428], \quad (76)$$

$$B = [0.56258, -0.0171, -1.3064, 1.0575, 0.8854], \quad (77)$$

$$C = [1.09025, -0.7992, 5.5721, -5.8676, 3.0454] \quad (78)$$

and

$$\lambda = [0.006, 0.2, 0.004]. \quad (79)$$

15.1.17 LTA: Local τ Approximation

J. P. Perdew and Y. Wang, J. Chem. Phys. 111, 911 (1999)

LSDA exchange functional with density represented as a function of τ .

$$K = \frac{1}{2} \sum_s E(2\tau_s), \quad (80)$$

where

$$E(\alpha) = c \left(\frac{5\alpha}{3(3\pi^2)^{2/3}} \right)^{4/5} \quad (81)$$

and

$$c = -\frac{3}{4} \left(\frac{3}{\pi} \right)^{1/3}. \quad (82)$$

15.1.18 LYP: Lee, Yang and Parr Correlation Functional

C. Lee, W. Yang and R. G. Parr, Phys. Rev. B 37, 785(1988); B. Miehlich, A. Savin, H. Stoll and H. Preuss, Chem. Phys. Letters 157, 200 (1989)

$$\begin{aligned} K &= 4 \frac{A\rho_\alpha\rho_\beta Z}{\rho} + AB\omega\sigma(\rho_\alpha\rho_\beta(47-7\delta)/18 - 2\rho^2/3) \\ &+ \sum_s AB\omega \left(\rho_s\rho_{\bar{s}} \left(82^{2/3} e\rho_s^{8/3} - (5/2 - \delta/18)\sigma_{ss} - \frac{(\delta-11)\rho_s\sigma_{ss}}{9\rho} \right) \right. \\ &\left. + (2\rho^2/3 - \rho_s^2)\sigma_{\bar{s}\bar{s}} \right), \end{aligned} \quad (83)$$

where

$$\omega = e^{-\frac{c}{\rho^{1/3}}} Z\rho^{-11/3}, \quad (84)$$

$$\delta = \frac{c+dZ}{\rho^{1/3}}, \quad (85)$$

$$B = 0.04918, \quad (86)$$

$$A = 0.132, \quad (87)$$

$$c = 0.2533, \quad (88)$$

$$d = 0.349, \quad (89)$$

$$e = \frac{3}{10} (3\pi^2)^{2/3} \quad (90)$$

and

$$Z = \left(1 + \frac{d}{\rho^{1/3}} \right)^{-1}. \quad (91)$$

15.1.19 MK00: Exchange Functional for Accurate Virtual Orbital Energies

F. R. Manby and P. J. Knowles, J. Chem. Phys. 112, 7002 (2000)

$$K = -\sum_s \frac{3\pi\rho_s^3}{\tau_s - \nu_s/4}. \quad (92)$$

15.1.20 MK00B: Exchange Functional for Accurate Virtual Orbital Energies

F. R. Manby and P. J. Knowles, J. Chem. Phys. 112, 7002 (2000)

MK00 with gradient correction of the form of B88X.

$$K = \sum_s \frac{3\pi\rho_s^3}{\tau_s - \nu_s/4} - \frac{\beta\rho_s^{4/3}\chi_s^2}{1 + 6\beta\chi_s \operatorname{arcsinh}(\chi_s)}, \quad (93)$$

where

$$\beta = 0.0016. \quad (94)$$

15.1.21 P86:

J. P. Perdew, Phys. Rev. B 33, 8822 (1986)

VWN with gradient correction.

$$K = \rho e + \frac{e^{-\Phi} C(r) \sigma}{d\rho^{4/3}}, \quad (95)$$

where

$$x = \left(\frac{3}{4\pi\rho} \right)^{1/6}, \quad (96)$$

$$\zeta = \frac{\rho\alpha - \rho\beta}{\rho}, \quad (97)$$

$$e = \Lambda + \omega y (1 + h\zeta^4), \quad (98)$$

$$y = \frac{9}{8} (1 + \zeta)^{4/3} + \frac{9}{8} (1 - \zeta)^{4/3} - \frac{9}{4}, \quad (99)$$

$$h = \frac{4(\lambda - \Lambda)}{9(2^{1/3} - 1)\omega} - 1, \quad (100)$$

$$\Lambda = q(k_1, l_1, m_1, n_1), \quad (101)$$

$$\lambda = q(k_2, l_2, m_2, n_2), \quad (102)$$

$$\omega = q(k_3, l_3, m_3, n_3), \quad (103)$$

$$q(A, p, c, d) = A \left(\ln \left(\frac{x^2}{X(x, c, d)} \right) + \frac{2c}{Q(c, d)} \arctan \left(\frac{Q(c, d)}{2x+c} \right) - \frac{cp}{X(p, c, d)} \left(\ln \left(\frac{(x-p)^2}{X(x, c, d)} \right) + 2 \frac{(c+2p)}{Q(c, d)} \arctan \left(\frac{Q(c, d)}{2x+c} \right) \right) \right), \quad (104)$$

$$Q(c, d) = \sqrt{4d - c^2}, \quad (105)$$

$$X(i, c, d) = i^2 + ci + d, \quad (106)$$

$$\Phi = 0.007390075 \frac{z\sqrt{\sigma}}{C(r)\rho^{7/6}}, \quad (107)$$

$$d = 2^{1/3} \sqrt{(1/2 + \zeta/2)^{5/3} + (1/2 - \zeta/2)^{5/3}}, \quad (108)$$

$$C(r) = 0.001667 + \frac{0.002568 + \alpha r + \beta r^2}{1 + \xi r + \delta r^2 + 10000\beta r^3}, \quad (109)$$

$$z = 0.11, \quad (110)$$

$$\alpha = 0.023266, \quad (111)$$

$$\beta = 0.000007389, \quad (112)$$

$$\xi = 8.723, \quad (113)$$

$$\delta = 0.472, \quad (114)$$

$$k = [0.0310907, 0.01554535, -1/(6\pi)^2], \quad (115)$$

$$l = [-0.10498, -0.325, -0.0047584], \quad (116)$$

$$m = [3.72744, 7.06042, 1.13107] \quad (117)$$

and

$$n = [12.9352, 18.0578, 13.0045]. \quad (118)$$

15.1.22 PBE: PBE = PW91C + PBEX

J. P. Perdew, K. Burke and M. Ernzerhof, Phys. Rev. Lett. 77, 3865 (1996)

15.1.23 PBE0: PBE0 = 0.75PBEX + PW91C + 0.25 Exact Exchange

C. Adamo and V. Barone, J. Chem. Phys. 110, 6158 (1999)

15.1.24 PBEX: PBE Exchange Functional

J. P. Perdew, K. Burke and M. Ernzerhof, Phys. Rev. Lett. 77, 3865 (1996)

$$K = \frac{1}{2} \sum_s E(2\rho_s), \quad (119)$$

where

$$E(n) = -\frac{3}{4} \left(\frac{3}{\pi}\right)^{1/3} n^{4/3} F(S), \quad (120)$$

$$S = \frac{\chi_s}{2(6\pi^2)^{1/3}}, \quad (121)$$

$$F(S) = 1 + R - R \left(1 + \frac{\mu S^2}{R}\right)^{-1}, \quad (122)$$

$$R = 0.804, \quad (123)$$

$$\mu = \delta\pi^2/3 \quad (124)$$

and

$$\delta = 0.066725. \quad (125)$$

15.1.25 PW86:

J. P. Perdew and Y. Wang, Phys. Rev. B 33, 8800 (1986)

$$K = \frac{1}{2} \sum_s E(2\rho_s), \quad (126)$$

where

$$E(n) = -\frac{3}{4} \left(\frac{3}{\pi}\right)^{1/3} n^{4/3} F(S), \quad (127)$$

$$F(S) = (1 + 1.296S^2 + 14S^4 + 0.2S^6)^{1/15} \quad (128)$$

and

$$S = \frac{\chi_s}{2(6\pi^2)^{1/3}}. \quad (129)$$

15.1.26 PW91: PW91=PW91X+PW91C

J. P. Perdew, J. A. Chevary, S. H. Vosko, K. A. Jackson, M. R. Pederson and C. Fiolhais, Phys. Rev. B 46, 6671 (1992)

15.1.27 PW91C: Perdew-Wang 1991 GGA Correlation Functional

J. P. Perdew, J. A. Chevary, S. H. Vosko, K. A. Jackson, M. R. Pederson and C. Fiolhais, Phys. Rev. B 46, 6671 (1992)

$$K = \rho (\epsilon(\rho_\alpha, \rho_\beta) + H(d, \rho_\alpha, \rho_\beta)), \quad (130)$$

where

$$d = \frac{\sqrt{\sigma}}{4u(\rho_\alpha, \rho_\beta)} \left(\frac{\pi}{3\rho^7}\right)^{1/6} \quad (131)$$

$$u(\alpha, \beta) = \left\{ (1 + \zeta(\alpha, \beta))^{2/3} + (1 - \zeta(\alpha, \beta))^{2/3} \right\} / 2, \quad (132)$$

$$H(d, \alpha, \beta) = L(d, \alpha, \beta) + J(d, \alpha, \beta), \quad (133)$$

$$L(d, \alpha, \beta) = \frac{u^3(\rho_\alpha, \rho_\beta) \lambda^2}{2\mathfrak{v}} \ln \left(1 + 2 \frac{\mathfrak{v} (d^2 + A(\alpha, \beta) d^4)}{\lambda (1 + A(\alpha, \beta) d^2 + A^2(\alpha, \beta) d^4)} \right), \quad (134)$$

$$J(d, \alpha, \beta) = \mathfrak{v} (\phi(r(\alpha, \beta)) - \kappa - 3Z/7) u^3(\rho_\alpha, \rho_\beta) d^2 e^{-400 \frac{u^4(\rho_\alpha, \rho_\beta) d^2}{(3\pi^5)^{1/3} \rho}}, \quad (135)$$

$$A(\alpha, \beta) = \frac{2\mathfrak{v}}{\lambda} \left(e^{-\frac{2\mathfrak{v}\epsilon(\alpha, \beta)}{u^3(\rho_\alpha, \rho_\beta) \lambda^2}} - 1 \right)^{-1}, \quad (136)$$

$$\mathfrak{v} = 0.09, \quad (137)$$

$$\lambda = \mathfrak{v} \kappa, \quad (138)$$

$$\mathfrak{v} = 16 \left(\frac{3}{\pi}\right)^{1/3}, \quad (139)$$

$$\kappa = 0.004235, \quad (140)$$

$$Z = -0.001667, \quad (141)$$

$$\phi(r) = \theta(r) - Z, \quad (142)$$

$$\theta(r) = \frac{1}{1000} \frac{2.568 + \Xi r + \Phi r^2}{1 + \Lambda r + \Upsilon r^2 + 10\Phi r^3}, \quad (143)$$

$$\Xi = 23.266, \quad (144)$$

$$\Phi = 0.007389, \quad (145)$$

$$\Lambda = 8.723, \quad (146)$$

$$\Upsilon = 0.472 \quad (147)$$

and $\varepsilon(\alpha, \beta)$ is the correlation energy per particle of the Local Spin Density Approximation(PW92C).

15.1.28 PW91X: Perdew-Wang 1991 GGA Exchange Functional

J. P. Perdew, J. A. Chevary, S. H. Vosko, K. A. Jackson, M. R. Pederson and C. Fiolhais, Phys. Rev. B 46, 6671 (1992)

$$K = \frac{1}{2} \sum_s E(2\rho_s), \quad (148)$$

where

$$E(n) = -\frac{3}{4} \left(\frac{3}{\pi}\right)^{1/3} n^{4/3} F(S), \quad (149)$$

$$S = \frac{\chi_s}{2(6\pi^2)^{1/3}} \quad (150)$$

and

$$F(S) = \frac{1 + 0.19645 S \operatorname{arcsinh}(7.7956 S) + (0.2743 - 0.1508 e^{-100 S^2}) S^2}{1 + 0.19645 S \operatorname{arcsinh}(7.7956 S) + 0.004 S^4}. \quad (151)$$

15.1.29 PW92C: Local Spin Density Approximation Correlation Energy

J. P. Perdew and Y. Wang, Phys. Rev. B 45, 13244 (1992)

Electron-gas correlation energy.

$$K = \rho \varepsilon(\rho_\alpha, \rho_\beta), \quad (152)$$

where

$$\begin{aligned} \varepsilon(\alpha, \beta) = & e(r(\alpha, \beta), T_1, U_1, V_1, W_1, X_1, Y_1, P_1) \\ & - e(r(\alpha, \beta), T_3, U_3, V_3, W_3, X_3, Y_3, P_3) \omega(\zeta(\alpha, \beta)) (1 - \zeta^4(\alpha, \beta)) / c \\ & + (e(r(\alpha, \beta), T_2, U_2, V_2, W_2, X_2, Y_2, P_2) \\ & - e(r(\alpha, \beta), T_1, U_1, V_1, W_1, X_1, Y_1, P_1)) \omega(\zeta(\alpha, \beta)) \zeta^4(\alpha, \beta), \end{aligned} \quad (153)$$

$$r(\alpha, \beta) = \left(\frac{3}{4\pi(\alpha + \beta)}\right)^{1/3}, \quad (154)$$

$$\zeta(\alpha, \beta) = \frac{\alpha - \beta}{\alpha + \beta}, \quad (155)$$

$$\omega(z) = \frac{(1+z)^{4/3} + (1-z)^{4/3} - 2}{2^{4/3} - 2}, \quad (156)$$

$$e(r, t, u, v, w, x, y, p) = -2t(1 + ur) \ln \left(1 + \frac{1}{2t(v\sqrt{r} + wr + xr^{3/2} + yr^{p+1})} \right), \quad (157)$$

$$c = 1.709921, \quad (158)$$

$$T = [0.031091, 0.015545, 0.016887], \quad (159)$$

$$U = [0.21370, 0.20548, 0.11125], \quad (160)$$

$$V = [7.5957, 14.1189, 10.357], \quad (161)$$

$$W = [3.5876, 6.1977, 3.6231], \quad (162)$$

$$X = [1.6382, 3.3662, 0.88026], \quad (163)$$

$$Y = [0.49294, 0.62517, 0.49671] \quad (164)$$

and

$$P = [1, 1, 1]. \quad (165)$$

LSDAC and LSDC are aliased to PW92C

15.1.30 S: Slater-Dirac Exchange Energy

J. C. Slater, Phys. Rev. 81, 385 (1951)

$$K = -c \sum_s \rho_s^{4/3}, \quad (166)$$

where

$$c = \frac{3}{2} \left(\frac{3}{4\pi} \right)^{1/3} \quad (167)$$

15.1.31 TH1:

D. J. Tozer and N. C. Handy, J. Chem. Phys. 108, 2545 (1998)

Density and gradient dependent first row exchange-correlation functional.

$$K = \sum_{i=1}^n \omega_i R_i S_i X_i Y_i, \quad (168)$$

where

$$n = 21, \quad (169)$$

$$R_i = \rho_\alpha^{t_i} + \rho_\beta^{t_i}, \quad (170)$$

$$S_i = \left(\frac{\rho_\alpha - \rho_\beta}{\rho} \right)^{2u_i}, \quad (171)$$

$$X_i = \frac{\sigma_{\alpha\alpha}^{v_i/2} + \sigma_{\beta\beta}^{v_i/2}}{2\rho^{4v_i/3}}, \quad (172)$$

$$Y_i = \left(\frac{\sigma_{\alpha\alpha} + \sigma_{\beta\beta} - 2\sqrt{\sigma_{\alpha\alpha}\sigma_{\beta\beta}}}{\rho^{8/3}} \right)^{w_i}, \quad (173)$$

$$t = [7/6, 4/3, 3/2, 5/3, 4/3, 3/2, 5/3, \frac{11}{6}, 3/2, 5/3, \frac{11}{6}, 2, 3/2, 5/3, \frac{11}{6}, 2, 7/6, 4/3, 3/2, 5/3, 1], \quad (174)$$

$$u = [0, 0, 0, 0, 0, 0, 0, 0, 0, 0, 0, 0, 0, 0, 0, 0, 1, 1, 1, 1, 0], \quad (175)$$

$$v = [0, 0, 0, 0, 1, 1, 1, 1, 2, 2, 2, 0, 0, 0, 0, 0, 0, 0, 0, 0], \quad (176)$$

$$w = [0, 0, 0, 0, 0, 0, 0, 0, 0, 0, 0, 0, 1, 1, 1, 1, 0, 0, 0, 0, 0] \quad (177)$$

and

$$\begin{aligned} \omega = & [-0.728255, 0.331699, -1.02946, 0.235703, -0.0876221, 0.140854, \\ & 0.0336982, -0.0353615, 0.00497930, -0.0645900, 0.0461795, \\ & -0.00757191, -0.00242717, 0.0428140, -0.0744891, 0.0386577, \\ & -0.352519, 2.19805, -3.72927, 1.94441, 0.128877]. \end{aligned} \quad (178)$$

15.1.32 TH2:

D. J. Tozer and N. C. Handy, J. Chem. Phys. 102, 3162 (1998)

Density and gradient dependent first row exchange-correlation functional of the form TH1 but with

$$n = 19, \quad (179)$$

$$t = \left[\frac{13}{12}, 7/6, 4/3, 3/2, 5/3, \frac{17}{12}, 3/2, 5/3, \frac{11}{6}, 5/3, \right. \\ \left. \frac{11}{6}, 2, 5/3, \frac{11}{6}, 2, 7/6, 4/3, 3/2, 5/3 \right], \quad (180)$$

$$u = [0, 0, 0, 0, 0, 0, 0, 0, 0, 0, 0, 0, 0, 0, 0, 0, 1, 1, 1, 1], \quad (181)$$

$$v = [0, 0, 0, 0, 0, 1, 1, 1, 1, 2, 2, 2, 0, 0, 0, 0, 0, 0, 0, 0], \quad (182)$$

$$w = [0, 0, 0, 0, 0, 0, 0, 0, 0, 0, 0, 0, 0, 1, 1, 1, 0, 0, 0, 0] \quad (183)$$

and

$$\begin{aligned} \omega = & [0.678831, -1.75821, 1.27676, -1.60789, 0.365610, \\ & -0.181327, 0.146973, 0.147141, -0.0716917, -0.0407167, \\ & 0.0214625, -0.000768156, 0.0310377, -0.0720326, 0.0446562, \\ & -0.266802, 1.50822, -1.94515, 0.679078]. \end{aligned} \quad (184)$$

15.1.33 TH3:

D. J. Tozer and N. C. Handy, Mol. Phys. 94, 70 (1998)

Density and gradient dependent first and second row exchange-correlation functional of the form TH2 but with

$$t = \left[7/6, 4/3, 3/2, 5/3, \frac{17}{12}, 3/2, 5/3, \frac{11}{6}, 5/3, \right. \\ \left. \frac{11}{6}, 2, 5/3, \frac{11}{6}, 2, 7/6, 4/3, 3/2, 5/3, \frac{13}{12} \right], \quad (185)$$

$$u = [0, 0, 0, 0, 0, 0, 0, 0, 0, 0, 0, 0, 0, 0, 0, 0, 1, 1, 1, 1, 0], \quad (186)$$

$$v = [0, 0, 0, 0, 1, 1, 1, 1, 2, 2, 2, 0, 0, 0, 0, 0, 0, 0, 0, 0, 0], \quad (187)$$

$$w = [0, 0, 0, 0, 0, 0, 0, 0, 0, 0, 0, 0, 1, 1, 1, 0, 0, 0, 0, 0] \quad (188)$$

and

$$\begin{aligned} \omega = & [-0.142542, -0.783603, -0.188875, 0.0426830, -0.304953, \\ & 0.430407, -0.0997699, 0.00355789, -0.0344374, 0.0192108, \\ & -0.00230906, 0.0235189, -0.0331157, 0.0121316, 0.441190, \\ & -2.27167, 4.03051, -2.28074, 0.0360204]. \end{aligned} \quad (189)$$

15.1.34 TH4:

D. J. Tozer and N. C. Handy, Mol. Phys. 94, 70 (1998)

Density an gradient dependent first and second row exchange-correlation functional of the form TH2 but with

$$\begin{aligned} t = & [7/6, 4/3, 3/2, 5/3, \frac{17}{12}, 3/2, 5/3, \frac{11}{6}, 5/3, \\ & \frac{11}{6}, 2, 5/3, \frac{11}{6}, 2, 7/6, 4/3, 3/2, 5/3, \frac{13}{12}], \end{aligned} \quad (190)$$

$$u = [0, 0, 0, 0, 0, 0, 0, 0, 0, 0, 0, 0, 0, 0, 1, 1, 1, 1, 0], \quad (191)$$

$$v = [0, 0, 0, 0, 1, 1, 1, 1, 2, 2, 2, 0, 0, 0, 0, 0, 0, 0, 0], \quad (192)$$

$$w = [0, 0, 0, 0, 0, 0, 0, 0, 0, 0, 0, 0, 1, 1, 1, 0, 0, 0, 0, 0] \quad (193)$$

and

$$\begin{aligned} \omega = & [0.0677353, -1.06763, -0.0419018, 0.0226313, -0.222478, \\ & 0.283432, -0.0165089, -0.0167204, -0.0332362, 0.0162254, \\ & -0.000984119, 0.0376713, -0.0653419, 0.0222835, 0.375782, \\ & -1.90675, 3.22494, -1.68698, -0.0235810]. \end{aligned} \quad (194)$$

15.1.35 THGFL:

D. J. Tozer, N. C. Handy and W. H. Green, Chem. Phys. Lett. 273, 183 (1997)

Density dependent first row exchange-correlation functional for closed shell systems.

$$K = \sum_{i=1}^n \omega_i R_i, \quad (195)$$

where

$$n = 4, \quad (196)$$

$$R_i = \rho_{\alpha}^{t_i} + \rho_{\beta}^{t_i}, \quad (197)$$

$$t = [7/6, 4/3, 3/2, 5/3] \quad (198)$$

and

$$\omega = [-1.06141, 0.898203, -1.34439, 0.302369]. \quad (199)$$

15.1.36 THGFC:

D. J. Tozer, N. C. Handy and W. H. Green, Chem. Phys. Lett. 273, 183 (1997)

Density and gradient dependent first row exchange-correlation functional for closed shell systems. Total energies are improved by adding DN , where N is the number of electrons and $D = 0.1863$.

$$K = \sum_{i=1}^n \omega_i R_i X_i, \quad (200)$$

where

$$n = 12, \quad (201)$$

$$R_i = \rho_{\alpha}^{t_i} + \rho_{\beta}^{t_i}, \quad (202)$$

$$X_i = \frac{\sigma_{\alpha\alpha}^{v_i/2} + \sigma_{\beta\beta}^{v_i/2}}{2\rho^{4v_i/3}}, \quad (203)$$

$$t = [7/6, 4/3, 3/2, 5/3, 4/3, 3/2, 5/3, \frac{11}{6}, 3/2, 5/3, \frac{11}{6}, 2], \quad (204)$$

$$v = [0, 0, 0, 0, 1, 1, 1, 1, 2, 2, 2, 2] \quad (205)$$

and

$$\begin{aligned} \omega = & [-0.864448, 0.565130, -1.27306, 0.309681, -0.287658, 0.588767, \\ & -0.252700, 0.0223563, 0.0140131, -0.0826608, 0.0556080, \\ & -0.00936227]. \end{aligned} \quad (206)$$

15.1.37 THGFCFO:

D. J. Tozer, N. C. Handy and W. H. Green, Chem. Phys. Lett. 273, 183 (1997)

Density and gradient dependent first row exchange-correlation functional. The closed- and open-shell parts are fitted to training sets of closed- and open-shell systems independently.

$$K = \sum_{i=1}^n \omega_i R_i S_i X_i Y_i, \quad (207)$$

where

$$n = 20, \quad (208)$$

$$R_i = \rho_{\alpha}^{t_i} + \rho_{\beta}^{t_i}, \quad (209)$$

$$S_i = \left(\frac{\rho_{\alpha} - \rho_{\beta}}{\rho} \right)^{2u_i}, \quad (210)$$

$$X_i = \frac{\sigma_{\alpha\alpha}^{v_i/2} + \sigma_{\beta\beta}^{v_i/2}}{2\rho^{4v_i/3}}, \quad (211)$$

$$Y_i = \left(\frac{\sigma_{\alpha\alpha} + \sigma_{\beta\beta} - 2\sqrt{\sigma_{\alpha\alpha}\sigma_{\beta\beta}}}{\rho^{8/3}} \right)^{w_i}, \quad (212)$$

$$t = [7/6, 4/3, 3/2, 5/3, 4/3, 3/2, 5/3, \frac{11}{6}, 3/2, 5/3, \frac{11}{6}, 2, 3/2, 5/3, \frac{11}{6}, 2, 7/6, 4/3, 3/2, 5/3], \quad (213)$$

$$u = [0, 0, 0, 0, 0, 0, 0, 0, 0, 0, 0, 0, 0, 0, 0, 0, 1, 1, 1, 1], \quad (214)$$

$$v = [0, 0, 0, 0, 1, 1, 1, 1, 2, 2, 2, 2, 0, 0, 0, 0, 0, 0, 0, 0], \quad (215)$$

$$w = [0, 0, 0, 0, 0, 0, 0, 0, 0, 0, 0, 0, 1, 1, 1, 1, 0, 0, 0, 0] \quad (216)$$

and

$$\begin{aligned} \omega = & [-0.864448, 0.565130, -1.27306, 0.309681, -0.287658, \\ & 0.588767, -0.252700, 0.0223563, 0.0140131, -0.0826608, \\ & 0.0556080, -0.00936227, -0.00677146, 0.0515199, -0.0874213, \\ & 0.0423827, 0.431940, -0.691153, -0.637866, 1.07565]. \end{aligned} \quad (217)$$

15.1.38 THGFCO:

D. J. Tozer, N. C. Handy and W. H. Green, Chem. Phys. Lett. 273, 183 (1997)

Density and gradient dependent first row exchange-correlation functional of the form FCFO but fitted to a training both set of open- and closed-shell systems.

$$\begin{aligned} \omega = & [-0.962998, 0.860233, -1.54092, 0.381602, -0.210208, \\ & 0.391496, -0.107660, -0.0105324, 0.00837384, -0.0617859, \\ & 0.0383072, -0.00526905, -0.00381514, 0.0321541, -0.0568280, \\ & 0.0288585, 0.368326, -0.328799, -1.22595, 1.36412]. \end{aligned} \quad (218)$$

15.1.39 VSXC:

T. Van Voorhis and G. E. Scuseria, J. Chem. Phys. 109, 400 (1998)

$$\begin{aligned} K = & F(x, z, p_3, q_3, r_3, t_3, u_3, v_3, \alpha_3) (\epsilon(\rho_\alpha, \rho_\beta) - \epsilon(\rho_\alpha, 0) - \epsilon(\rho_\beta, 0)) \\ & + \sum_s (\rho_s)^{4/3} F(\chi_s, z_s, p_1, q_1, r_1, t_1, u_1, v_1, \alpha_1) \\ & + d_s \epsilon(\rho_s, 0) F(\chi_s, z_s, p_2, q_2, r_2, t_2, u_2, v_2, \alpha_2), \end{aligned} \quad (219)$$

where

$$x = \chi_\alpha^2 + \chi_\beta^2, \quad (220)$$

$$z_s = \frac{\tau_s}{\rho_s^{5/3}} - c_f, \quad (221)$$

$$z = \frac{\tau_\alpha}{\rho_\alpha^{5/3}} + \frac{\tau_\beta}{\rho_\beta^{5/3}} - 2c_f, \quad (222)$$

$$d_s = 1 - \frac{\chi_s^2}{4z_s + 4c_f}, \quad (223)$$

$$F(x, z, p, q, c, d, e, f, \alpha) = \frac{p}{\lambda(x, z, \alpha)} + \frac{qx^2 + cz}{\lambda^2(x, z, \alpha)} + \frac{dx^4 + ex^2z + fz^2}{\lambda^3(x, z, \alpha)}, \quad (224)$$

$$\lambda(x, z, \alpha) = 1 + \alpha(x^2 + z), \quad (225)$$

$$c_f = 3(3\pi^2)^{2/3}/5, \quad (226)$$

$$p = [-0.98, 0.3271, 0.7035], \quad (227)$$

$$q = [-0.003557, -0.03229, 0.007695], \quad (228)$$

$$r = [0.00625, -0.02942, 0.05153], \quad (229)$$

$$t = [-0.00002354, 0.002134, 0.00003394], \quad (230)$$

$$u = [-0.0001283, -0.005452, -0.001269], \quad (231)$$

$$v = [0.0003575, 0.01578, 0.001296], \quad (232)$$

$$\alpha = [0.001867, 0.005151, 0.00305] \quad (233)$$

and $\epsilon(\alpha, \beta)$ is the correlation energy per particle of the Local Spin Density Approximation(PW92C). VS99 is aliased to VSXC.

15.1.40 VWN: Vosko-Wilk-Nusair (1980) local correlation energy

S. H. Vosko, L. Wilk and M. Nusair, Can. J. Phys. 58, 1200 (1980)

$$K = \rho e, \quad (234)$$

where

$$x = \left(\frac{3}{4\pi\rho} \right)^{1/6}, \quad (235)$$

$$\zeta = \frac{\rho_\alpha - \rho_\beta}{\rho}, \quad (236)$$

$$e = \Lambda + \alpha y (1 + h\zeta^4), \quad (237)$$

$$y = \frac{9}{8} (1 + \zeta)^{4/3} + \frac{9}{8} (1 - \zeta)^{4/3} - \frac{9}{4}, \quad (238)$$

$$h = \frac{4(\lambda - \Lambda)}{9(2^{1/3} - 1)\alpha} - 1, \quad (239)$$

$$\Lambda = q(k_1, l_1, m_1, n_1), \quad (240)$$

$$\lambda = q(k_2, l_2, m_2, n_2), \quad (241)$$

$$\alpha = q(k_3, l_3, m_3, n_3), \quad (242)$$

$$q(A, p, c, d) = A \left(\ln \left(\frac{x^2}{X(x, c, d)} \right) + \frac{2c}{Q(c, d)} \arctan \left(\frac{Q(c, d)}{2x+c} \right) - \frac{cp}{X(p, c, d)} \left(\ln \left(\frac{(x-p)^2}{X(x, c, d)} \right) + 2 \frac{(c+2p)}{Q(c, d)} \arctan \left(\frac{Q(c, d)}{2x+c} \right) \right) \right), \quad (243)$$

$$Q(c, d) = \sqrt{4d - c^2}, \quad (244)$$

$$X(i, c, d) = i^2 + ci + d, \quad (245)$$

$$k = [0.0310907, 0.01554535, -1/(6\pi)^2], \quad (246)$$

$$l = [-0.10498, -0.325, -0.0047584], \quad (247)$$

$$m = [3.72744, 7.06042, 1.13107] \quad (248)$$

and

$$n = [12.9352, 18.0578, 13.0045]. \quad (249)$$

15.1.41 Alias functional descriptions

Additional functional keywords are also defined as convenient aliases. The following table gives the translations.

LDA	S+VWN
DIRAC	S
B88	B
B88X	B
LYP88	LYP
B-LYP	B88+LYP88
B3LYP	0.72*B88+0.08*S+0.81*LYP88+0.19*VWN+0.2*exact exchange
B97	B97DF+0.1943*exact exchange
B97R	B97RDF+0.21*exact exchange
PBE0	PBE0DF+0.25*exact exchange
LSDAC	PW92C
LSDC	PW92C
VS99	VSXC

15.2 Options

The following options may be used to control the operation of the DFT modules. In the Kohn-Sham case, these may be *followed* by further options for the SCF program as described in Section 14. Note that DFT and SCF options cannot be intermixed.

15.2.1 Density source (DENSITY, ODENSITY)

DENSITY,*orb.filec*,... ODENSITY,*orbo.fileo*,...

For non-self-consistent DFT calculations, specifies the source of the density matrix. The total density is read from *orb.filec*, with further options specifying density sets in the standard way as described in Section 2.16. ODENSITY can be used to specify the spin density. The defaults are the densities last written by an SCF or MCSCF program.

15.2.2 Thresholds (THR)

THR,*key1=value1,key2=value2*...

Sets various truncation thresholds. *key* can be one of the following.

TOTAL	Overall target accuracy of density functional. Defaults to the value of the global threshold ENERGY. For proper use of this threshold, other thresholds should be left at their default value of zero.
ORBITAL	Orbital truncation threshold.
DENSITY	Density truncation threshold.
FOCK	Fock matrix truncation threshold.

15.2.3 Exact exchange computation (EXCHANGE)EXCHANGE,*factor*

For Kohn-Sham calculations, compute exchange energy according to Hartree-Fock formalism and add the contribution scaled by *factor* to the fock matrix and the energy functional. Otherwise, the default is *factor*=0, i.e., the exchange is assumed to be contained in the functional, and only the Coulomb interaction is calculated explicitly.

FACTOR,*fac1, fac2, ...*

Provide a factor for each functional specified. The functionals will be combined accordingly. By default, all factors are one.

15.2.4 Exchange-correlation potential (POTENTIAL)POTENTIAL,*rec,fil*

For stand-alone DFT calculations, compute exchange-correlation potential pseudo-matrix elements, defined formally as the differential of the sum of all specified functionals with respect to elements of the atomic orbital density matrix. The matrix is written to record *rec* on file *fil*.

15.2.5 Grid blocking factor (BLOCK)BLOCK,*nblock*

Respecify the number of spatial integration points treated together as a block in the DFT integration routines (default 128). Increasing *nblock* may enhance efficiency on, e.g., vector architectures, but leads to increased memory usage.

15.2.6 Dump integrand values(DUMP)DUMP,*file,status*

Write out values of the integrand at grid points to the file *file*. The first line of *file* contains the number of functional components; there then follows a line for each functional giving the input key of the functional. Subsequent lines give the functional number, cartesian coordinates, integrand value and integration weight with Fortran format (I2 , 3F15 . 10 , F23 . 15).

15.3 Examples

The following shows the use of both non-self-consistent and self-consistent DFT.

```

geometry={c;n,c,r}
r=1.1 angstrom
df=[b,lyp]
rhf;method(1)=program
dft;edf(1)=dftfun
uhf;method(2)=program
dft;edf(2)=dftfun
uks;method(3)=program,edf(3)=dftfun
dft;method(4)=program,edf(4)=dftfun
table,dftname,dftfuns
table,method,edf

```

examples/
cndft.com

15.4 Numerical integration grid control (GRID)

Density functionals are evaluated through numerical quadrature on a grid in three-dimensional space. Although the sensible defaults will usually suffice, the parameters that define the grid can be specified by using the GRID top-level command, which should be presented *before* the data for the DFT or KS calculations that will use the grid.

GRID,*orb.file,status*

The integration grid is stored on record *orb.file* (default 1800.2). The information on disk consists of two parts: the parameters necessary to define the grid, and a cache of the evaluated grid points and weights. The latter is flagged as ‘dirty’ whenever any parameters are changed, and whenever the geometry changes; if the cache is dirty, then when an attempt is made to use the grid, it will be recalculated, otherwise the cached values are used.

If *status* is OLD, an attempt to restore the grid from a previous calculation is performed; effectively, the old grid provides a template of parameters which can be adjusted using the parameter commands described below. If *status* is NEW, the grid is always created with default parameters. If *status* is UNKNOWN (the default), a new grid is created either if record *orb.file* does not exist; otherwise the old grid is used.

The GRID command may be followed by a number of parameter-modifying subcommands. The currently implemented default parameters are equivalent to the following input commands.

```
THR,1e-5,0,0
RADIAL,LOG,3,1.0,20,25,25,30
ANGULAR,LEBEDEV,0.0,0.0
LMIN,3,5,5,7
LMAX,53,53,53,53
VORONOI,10
SAVE
SYM
```

15.4.1 Target quadrature accuracy (THR)

THR,*acc,accr,acca*

Specify the target accuracy of integration. Radial and angular grids are generated adaptively, with the aim of integrating the Slater-Dirac functional to the specified accuracy. *acc* is an overall target accuracy, and is the one that should normally be used; radial and angular grid target accuracies are generated algorithmically from it. However, they can be adjusted individually by specifying *accr* and *acca* respectively.

15.4.2 Radial integration grid (RADIAL)

RADIAL,*method,m_r,scale,n₀,n₁,n₂,n₃*

Specify the details of the radial quadrature scheme. Four different radial schemes are available, specified by *method* = EM, BECKE, AHLRICHS or LOG, with the latter being the default.

EM is the Euler-Maclaurin scheme defined by C. W. Murray, N. C. Handy and G. J. Laming, Mol. Phys. 78 (1993) 997. m_r , for which the default value is 2, is defined in equation (6) of the above as

$$r = \alpha \frac{x^{m_r}}{(1-x)^{m_r}} \quad (250)$$

whilst *scale* (default value 1) multiplied by the Bragg-Slater radius of the atom gives the scaling parameter α .

LOG is the scheme described by M. E. Mura and P. J. Knowles, J. Chem. Phys. 104 (1996) 9848. It is based on the transformation

$$r = -\alpha \log_e(1 - x^{m_r}), \quad (251)$$

with $0 \leq x \leq 1$ and simple Gauss quadrature in x -space. The recommended value of m_r is 3 for molecular systems, giving rise to the Log3 grid; $m_r=4$ is more efficient for atoms. α is taken to be *scale* times the recommended value for α given by Mura and Knowles, and *scale* defaults to 1.

BECKE is as defined by A. D. Becke, J. Chem. Phys. 88 (1988) 2547. It is based on the transformation

$$r = \alpha \frac{(1+x)}{(1-x)}, \quad (252)$$

using points in $-1 \leq x \leq +1$ and standard Gauss-Chebyshev quadrature of the second kind for the x -space quadrature. Becke chose his scaling parameters to be half the Bragg-Slater radius except for hydrogen, for which the whole Bragg-Slater radius was used, and setting *scale* to a value other than 1 allows a different α to be used. m_r is not necessary for this radial scheme.

AHLRICHS is the radial scheme defined by O. Treutler and R. Ahlrichs, J. Chem. Phys. 102 (1995) 346. It is based on the transformation their M4 mapping

$$r = \frac{\alpha}{\log_e 2} (1+x)^{0.6} \log_e \left(\frac{2}{1-x} \right), \quad (253)$$

with using standard Gauss-Chebyshev quadrature of the second kind for the x -space integration. m_r is not necessary for this radial scheme.

n_0, n_1, n_2, n_3 are the degrees of quadrature n_r (see equation (3) of Murray et al.), for hydrogen/helium, first row, second row, and other elements respectively.

`accr` as given by the THR command specifies a target accuracy; the number of radial points is chosen according to a model, instead of using an explicit n_i . The stricter of n_i , `accr` is used, unless either is zero, in which case it is ignored.

15.4.3 Angular integration grid (ANGULAR)

ANGULAR,*method,acca,crowd*

LMIN, $l_0^{\min}, l_1^{\min}, l_2^{\min}, l_3^{\min}$

LMAX, $l_0^{\max}, l_1^{\max}, l_2^{\max}, l_3^{\max}$

Specify the details of the angular quadrature scheme. The default choice for *method* is LEBEDEV (ie. as in A. D. Becke, J. Chem. Phys. 88 (1988) 2547) which provides angular grids of octahedral symmetry. The alternative choice for *method* is LEGENDRE which gives Gauss-Legendre quadrature in θ and simple quadrature in ϕ , as defined by C. W. Murray, N. C. Handy and G. J. Laming, Mol. Phys. 78 (1993) 997.

Each type of grid specifies a family of which the various members are characterized by a single quantum number l ; spherical harmonics up to degree l are integrated exactly. l_{\min_i} and l_{\max_i} , $i = 0, 1, 2, 3$ specify allowed ranges of l for hydrogen/helium, first row, second row, and other elements respectively. For the Lebedev grids, if the value of l is not one of the set implemented in MOLPRO (3, 5, 7, 9, 11, 13, 15, 17, 19, 23, 29, 41, 47, 53), then l is increased to give

the next largest angular grid available. In general, different radial points will have different l , and in the absence of any moderation described below, will be taken from l_i^{\max} .

crowd is a parameter to control the reduction of the degree of quadrature close to the nucleus, where points would otherwise be unnecessarily close together; larger values of *crowd* mean less reduction thus larger grids. A very large value of this parameter, or, conventionally, setting it to zero, will switch off this feature.

acca is a target energy accuracy. It is used to reduce l for a given radial point as far as possible below l_i^{\max} but not lower than l_i^{\max} . The implementation uses the error in the angular integral of the kernel of the Slater-Dirac exchange functional using a sum of approximate atomic densities. If *acca* is zero, the global threshold is used instead, or else it is ignored.

15.4.4 Atom partitioning of integration grid (VORONOI)

VORONOI, m_μ

Controls Becke-Voronoi partitioning of space. The algorithm of C. W. Murray, N. C. Handy and G. J. Laming, Mol. Phys. 78 (1993) 997 is used, with m_μ defined by equation (24). The default value is 10.

15.4.5 Grid caching (SAVE, NOSAVE)

NOSAVE

disables the disk caching of the grid, i.e. forces the recalculation of the grid each time it is needed.

SAVE

forces the use of a grid cache where possible.

15.4.6 Grid symmetry (SYM, NOSYM)

NOSYM

switches off the use of symmetry in generating the integration grid, whereas

SYM

forces the use of any point-group symmetry.

15.4.7 Grid printing (PRINT)

PRINT,*key=value*,...

controls printing of the grid, which by default is not done. At present, the only possible value for *key* is GRID, and *value* should be specified as an integer. GRID=0 causes the total number of integration points to be evaluated and reported; GRID=1 additionally shows the number of points on each atom; GRID=2 causes the complete set of grid points and weights to be printed.

16 ORBITAL LOCALIZATION

Localized orbitals are calculated according to the Boys or Pipek-Mezey criteria. Localization takes place within each symmetry species separately. If complete localization is desired, no symmetry should be used. All subcommands can be abbreviated by three characters.

The localization program is invoked by the LOCALI command

LOCALI [,*method*]

The keyword *method* can be either BOYS or PIPEK. By default, the valence orbitals from the last energy calculation are localized using the Boys criterion. Only orbital subsets which leave the energy invariant are transformed. These defaults can be modified using the optional commands described in the following sections.

16.1 Defining the input orbitals (ORBITAL)

ORBITAL,*record.file,specifications*

The orbitals to be localized are read from dump record *record.file*. A state specific orbital set can be selected using *specifications*, as explained in section 2.16. Default are the orbitals calculated last.

16.2 Saving the localized orbitals (SAVE)

SAVE,*record.file*

This specifies the dump record where the localized orbitals are stored. If the dump record already exists, the localized orbitals are added to it. Default is the input record (cf. ORBITAL).

16.3 Choosing the localization method (METHOD)

METHOD,*method*

The localization method *method* can be either BOYS or PIPEK. This can also be specified as argument on the LOCALI card (see above).

16.4 Delocalization of orbitals (DELOCAL)

DELOCAL

If this card is present, the orbitals are delocalized.

16.5 Localizing AOs(LOCAO)

LOCAO

If this card is present, the number of AOs contributing to each MO is minimized. This can be useful to rotate degenerate orbitals (e.g., px, py, pz in an atom) so that pure orbitals (in this case px, py, pz) result.

This implies Pipek-Mezey localization.

16.6 Selecting the orbital space

By default, only the valence orbitals are localized, in order to ensure invariance of subsequent electron correlation treatments. This behaviour can be modified using the OCC and CORE directives.

16.6.1 Defining the occupied space (OCC)

OCC, $o_1, o_2 \dots$

defines the highest orbital o_i in each symmetry i to be localized.

16.6.2 Defining the core orbitals (CORE)

CORE, $c_1, c_2 \dots$

The first c_i orbitals in each symmetry are treated as core orbitals and not localized. Thus, orbitals $c_i + 1$ to o_i are localized in symmetry i .

16.6.3 Defining groups of orbitals (GROUP, OFFDIAG)

GROUP, $orb1, orb2, orb3, \dots$

This card defines groups of orbitals to be localized as follows:

```
GROUP , 1 . 1 , 2 . 1 , 3 . 1      a group of orbitals 1-3 in symmetry 1
GROUP , 1 . 1 , -3 . 1           equivalent to previous example
GROUP , 3 . 1 , 5 . 1 , -8 . 1   this group includes orbitals 3,5,6,7,8 in symmetry 1
```

Orbitals in different groups are localized independently. Orbitals not included in any group are unchanged.

16.6.4 Localization between groups (OFFDIAG)

OFFDIAG

If this card is present, localize between groups instead of within groups.

16.7 Ordering of localized orbitals

ORDER, *type*

If *type*=CHARGE, the orbitals are ordered according to their charge centroids (default).

If *type*=FOCK, the orbitals are ordered according to increasing diagonal elements of the fock operator (PIPEK) or increasing Coulson-additive orbital energies (BOYS). This requires a Fock operator from the preceding energy calculation. For localization of Hartree-Fock orbitals, this operator is stored in the dump record and automatically found. For localization of MCSCF orbitals, an effective fock operator is computed from the MCSCF density matrix (see DENSITY option). Alternatively, a dump record of a previous SCF calculation can be specified on the FOCK card, and then the fock operator is read from this record. For degenerate orbitals, further ordering according to the the coordinates of charge centres is attempted (first according to largest z-coordinates, then according to x, then y).

16.7.1 No reordering (NOORDER)

NOORDER

If this card is present, the localized orbitals are not reordered. This is useful if localized orbitals are used as starting guess, and it is intended that their order remains unchanged.

16.7.2 Defining reference orbitals (REFORB)REFORB,*record.file,specifications*

The localized orbitals are reordered such that the overlap with the reference orbitals read from *record.file* is maximized. This is useful for local correlation treatments for keeping the order of the localized constant for different geometries. A state specific orbital set can be selected using *specifications*, as explained in section 2.16.

16.7.3 Selecting the fock matrix (FOCK)FOCK,*record.file*

This specifies a record holding a Fock operator to be used for ordering the orbitals. Note that only SCF dump records hold fock operators. Default is the Fock operator from the energy calculation which produced the input orbitals.

16.7.4 Selecting a density matrix (DENSITY)DENSITY,*record.file,specifications*

This specifies a record holding a density matrix for construction of a fock operator used for ordering the orbitals. This can be used if no fock operator is available, and has only an effect for MCSCF localizations. By default, the (state averaged) MCSCF density is used. A state specific density matrix can be selected using *specifications* as described in section 2.16.

16.8 Localization thresholds (THRESH)THRESH,*thresh,eorder*

thresh is a threshold for localization (default 1.d-12). If *eorder* is nonzero (default 1.d-4), the orbitals whose energy difference is smaller than *eorder* are considered to be degenerate and reordered according to the position of their charge centres (see section 16.7).

16.9 Printing options (PRINT)PRINT,[ORBITAL=]*pri* [,CHARGE] [,CENTRES] [,TEST] [,TRAN];

If ORB[ITAL] is given, the localized orbitals are printed.

If CHA[RGE] or CEN[TRES] is given, the charge centres of the localized orbitals are printed.

If TRAN is given, the transformation matrix is printed (Boys only).

If TEST is given, intermediate information is printed.

17 THE MCSCF PROGRAM MULTI

MULTI is a general MCSCF/CASSCF program written by
P. J. Knowles and H.-J. Werner (1984).

Bibliography:

H.-J. Werner and P. J. Knowles, *J. Chem. Phys.* 82, 5053 (1985).

P. J. Knowles and H.-J. Werner, *Chem. Phys. Lett.* 115, 259 (1985).

All publications resulting from use of this program must acknowledge the above. See also:

H.-J. Werner and W. Meyer, *J. Chem. Phys.* 73, 2342 (1980).

H.-J. Werner and W. Meyer, *J. Chem. Phys.* 74, 5794 (1981).

H.-J. Werner, *Adv. Chem. Phys.* LXIX, 1 (1987).

This program allows one to perform CASSCF as well as general MCSCF calculations. For CASSCF calculations, one can optionally use Slater determinants or CSFs as a N -electron basis. In most cases, the use of Slater determinants is more efficient. General MCSCF calculations must use CSFs as a basis.

A quite sophisticated optimization method is used. The algorithm is second-order in the orbital and CI coefficient changes and is therefore quadratically convergent. Since important higher order terms in the independent orbital parameters are included, almost cubic convergence is often observed. For simple cases, convergence is usually achieved in 2-3 iterations. However, convergence problems can still occur in certain applications, and usually indicate that the active space is not adequately chosen. For instance, if two weakly occupied orbitals are of similar importance to the energy, but only one of them is included in the active set, the program might alternate between them. In such cases either reduction or enlargement of the active orbital space can solve the problem. In other cases difficulties can occur if two electronic states in the same symmetry are almost or exactly degenerate, since then the program can switch from one state to the other. This might happen near avoided crossings or near an asymptote. Problems of this sort can be avoided by optimizing the energy average of the particular states. It is also possible to force convergence to specific states by choosing a subset of configurations as primary space (PSPACE). The hamiltonian is constructed and diagonalized explicitly in this space; the coefficients of the remaining configurations are optimized iteratively using the P-space wavefunction as zeroth order approximation. For linear molecules, another possibility is to use the LQUANT option, which makes it possible to force convergence to states with definite Λ quantum number, i.e., Σ , Π , Δ , etc. states.

17.1 Structure of the input

All sub-commands known to *MULTI* may be abbreviated by four letters. The input commands fall into several logical groups; within each group commands may appear in any order, but the groups must come in correct order.

- a) The program is invoked by the command MULTI or MCSCF
- b) cards defining partitioning of orbitals spaces – OCC , CORE , CLOSED
- c) general options (most commands not otherwise specified here)
- d) a WF card defining a state symmetry
- e) options pertaining to that state symmetry – WEIGHT , STATE , LQUANT

- f) configuration specification for that state symmetry – SELECT , CON , RESTRICT
- g) definition of the primary configurations for that state symmetry - PSPACE
- h) further general options

Stages d) through to h) may be repeated several times; this is the way in which you specify an average energy of several states of different symmetry.

17.2 Defining the orbital subspaces

17.2.1 Occupied orbitals

OCC, n_1, n_2, \dots, n_8 ;

n_i specifies numbers of occupied orbitals (including CORE and CLOSED) in irreducible representation number i . In the absence of an OCC card, the information from the most recent MCSCF calculation is used, or, if there is none, those orbitals corresponding to a minimal valence set, i.e., full valence space, are used.

17.2.2 Frozen-core orbitals

CORE, $n_1, n_2, \dots, record.file$;

n_i is the number of frozen-core orbitals in irrep number i . These orbitals are doubly occupied in all configurations and not optimized.

record.file is the record name for frozen core orbitals; if not supplied, taken from *orb* on START card. *record.file* can be specified in any field after the last nonzero n_i . It should always be given if the orbital guess is from a neighbouring geometry and should then specify the SCF orbitals calculated at the present geometry. If a subsequent gradient calculation is performed with this wavefunction, *record.file* is mandatory and must specify closed-shell SCF orbitals at the present geometry. Note that *record* must be larger than 2000.

If the CORE card is omitted, then the numbers of core orbitals are taken from the most recent MCSCF calculation, or otherwise no orbitals are frozen. If the CORE card is given as CORE,*record.file*, then the orbitals corresponding to atomic inner shells are taken, i.e., 1s for Li–Ne, 1s2s2p for Na–Ar, etc. A CORE card without any specification resets the number of frozen core orbitals to zero.

17.2.3 Closed-shell orbitals

CLOSED, n_1, n_2, \dots, n_8

n_i is the number of closed-shell orbitals in irrep number i , inclusive of any CORE orbitals. These orbitals do not form part of the active space, i.e., they are doubly occupied in all CSFs. In contrast to the core orbitals (see CORE), these orbitals are fully optimized. Please note that the program output sometimes says *CORE* when it means *CLOSED*, and *FROZEN* when it means *CORE* (historical reasons).

If the CLOSED card is omitted, then the data defaults to that of the most recent MCSCF calculation, or else the atomic inner shells as described above for CORE.

17.2.4 Freezing orbitals

FREEZE,*orb.sym*;

The specified orbital will not be optimized and will remain identical to the starting guess. *orb.sym* should be an active or closed-shell orbital. If *orb.sym* is a frozen core orbital, this card has no effect.

17.3 Defining the optimized states

Each state symmetry to be optimized is specified by one WF card, which may optionally be followed by STATE, WEIGHT, RESTRICT, SELECT, CON, and/or PSPACE cards. All cards belonging to a particular state symmetry as defined on the WF card must form a block which comes directly after the WF card. The cards can be in any order, however.

17.3.1 Defining the state symmetry

The number of electrons and the total symmetry of the wavefunction are specified on the WF card:

WF,*elec,sym,spin*

where

<i>elec</i>	is the number of electrons
<i>sym</i>	is the number of the irreducible representation
<i>spin</i>	defines the spin symmetry, $spin = 2S$ (singlet=0, doublet=1, triplet=2, etc.)

Note that these values take sensible defaults if any or all are not specified (see section 2.13).

The input directives STATE, WEIGHT, LQUANT, SELECT, PUNCSF always refer to the state symmetry as defined on the previous WF card. If such a directive is found before a WF card has been given, the current state symmetry is assumed, either from a previous calculation or from variables [MC]SYMMETRY(1) and [MC]SPIN(1) (if these are defined). If any of these cards or a WF card is given, the variables STATE, WEIGHT, LQUANT, SELECT are *not* used, and the number of state symmetries defaults to one, regardless of how many symmetries are specified in variable [MC]SYMMETRY.

17.3.2 Defining the number of states in the present symmetry

STATE,*nstate*;

nstate is the number of states in the present symmetry. By default, all states are optimized with weight 1 (see WEIGHT card).

17.3.3 Specifying weights in state-averaged calculations

WEIGHT, $w(1), w(2), \dots, w(nstate)$;

$w(i)$ is the weight for the state i in the present symmetry. By default, all weights are 1.0. See also STATE card. If you want to optimize the second state of a particular state symmetry alone, specify

```
STATE, 2; WEIGHT, 0, 1;
```

Note, however, that this might lead to root-flipping problems.

17.4 Defining the configuration space

By default, the program generates a complete configuration set (CAS) in the active space. The full space may be restricted to a certain occupation pattern using the RESTRICT option. Alternatively, configurations may be selected from the wavefunction of a previous calculation using SELECT, or explicitly specified on CON cards. Note that this program only allows to select or specify orbital configurations. For each orbital configuration, all spin couplings are always included. Possible RESTRICT, SELECT and CON cards must immediately follow the WF card which defines the corresponding state symmetry.

17.4.1 Occupation restrictions

```
RESTRICT, nmin, nmax, orb1, orb2, ... orbn;
```

This card can be used to restrict the occupation patterns. Only configurations containing between $nmin$ and $nmax$ electrons in the specified orbitals $orb_1, orb_2, \dots, orb_n$ are included in the wavefunction. If $nmin$ and $nmax$ are negative, configurations with exactly $abs(nmin)$ and $abs(nmax)$ electrons in the specified orbitals are deleted. This can be used, for instance, to omit singly excited configurations. The orbitals are specified in the form *number.sym*, where *number* is the number of the orbital in irrep *sym*. Several RESTRICT cards may follow each other. RESTRICT only works if a CONFIG card is specified before the first WF card.

RESTRICT cards given before the first WF cards are global, i.e., are active for all state symmetries. If such a global restrict card is given, variable [MC]RESTRICT is *not* used.

Additional state-specific RESTRICT cards may be given after a WF card. These are used in addition to the global orbital restrictions.

If neither state-specific nor global RESTRICT cards are found, the values from the variable [MC]RESTRICT are used.

17.4.2 Selecting configurations

```
SELECT, ref1, ref2, refthr, refstat, mxshrf;
```

This card is used to specify a configuration set other than a CAS, which is the default. This option automatically triggers the CONFIG option, which selects CSFs rather than determinants. Configurations can be defined using CON cards, which must follow immediately the SELECT card. Alternatively, if *ref1* is an existing MOLPRO record name, the configurations are read in from that record and may be selected according to a given threshold.

ref1=rec1.file

(*rec1* > 2000) The configurations are read in from the specified record. If *ref1* is not specified, the program assumes that the configurations are read from subsequent CON cards (see CON).

<i>ref2=rec2.file</i>	(<i>rec2</i> > 2000) Additional configurations are read from the specified record. If <i>rec2</i> is negative, all records between <i>rec1</i> and <i>abs(rec2)</i> are read. All configurations found in this way are merged.
<i>refthr</i>	Selection threshold for configurations read from disc (records <i>rec1</i> – <i>rec2</i>). This applies to the norm of all CSFs for each orbital configuration.
<i>refstat</i>	Specifies from which state vector the configurations are selected. This only applies to the case that the configurations were saved in a state-averaged calculation. If <i>refstat</i> is not specified, the configurations are selected from all states.
<i>mxshrf</i>	max. number of open shells in the selected or generated configurations.

17.4.3 Specifying orbital configurations

CON,*n*₁,*n*₂,*n*₃,*n*₄,...

Specifies an orbital configuration to be included in the present symmetry. The first CON card must be preceded by a SELECT card. *n*₁, *n*₂ etc. are the occupation numbers of the active orbitals (0,1,or 2). For example, for

```
OCC,5,2,2;CLOSED,2,1,1;
```

*n*₁ is the occupation of orbital 3.1 (number.sym), *n*₂ is the occupation of orbital 4.1, *n*₃ of 5.1, *n*₄ of 2.2, and *n*₅ of 2.3 Any number of CON cards may follow each other.

Example for the BH molecule:

```
OCC,4,1,1;           ! four sigma, one pi orbitals are occupied
CORE,1;             ! first sigma orbital is doubly occupied and frozen
WF,6,1;            ! 6 electrons, singlet Sigma+ state
SELECT              ! triggers configuration input
CON,2,2             ! 2sigma**2, 3sigma**2
CON,2,1,1          ! 2sigma**2, 3sigma, 4sigma
CON,2,0,2          ! 2sigma**2, 4sigma**2
CON,2,0,0,2        ! 2sigma**2, 1pi_x**2
CON,2,0,0,0,2      ! 2sigma**2, 1pi_y**2
```

17.4.4 Selecting the primary configuration set

PSPACE,*thresh*

The hamiltonian is constructed and diagonalized explicitly in the primary configuration space, which can be selected with the PSPACE card. The coefficients of the remaining configurations (Q-space) are optimized iteratively using the P-space wavefunction as zeroth order approximation.

If *thresh* is nonzero, it is a threshold for automatically selecting all configurations as P-space configurations which have energies less than *emin* + *thresh*, where *emin* is the lowest energy of all configurations. Further P-space configurations can be specified using CON cards, which must follow immediately after the PSPACE card. These are merged with the ones selected according to the threshold. Automatic selection can be avoided by specifying a very small threshold. There is a sensible default value for *thresh* (0.4), so you usually don't need a pspace card in your input. Furthermore, if the number of configurations in the MCSCF is less than 20, all configurations go into the P-space unless you give a PSPACE card in the input.

A P-space threshold defined on a PSPACE card before the first WF (or STATE, WEIGHT, SELECT, PUNCSF if WF is not given) card is global, i.e., valid for all state symmetries. State-specific thresholds can be defined by placing a PSPACE card after the corresponding WF card. In the latter case the PSPACE card can be followed by CON cards, which define state-specific P-space configurations.

17.4.5 Projection to specific Λ states in linear molecules

Since MOLPRO can only use Abelian point groups (e.g. C_{2v} instead of $C_{\infty v}$ for linear molecules), $\Delta_{x^2-y^2}$ states as well as Σ^+ states occur in the irreducible representation number 1, for example. Sometimes it is not possible to predict in advance to which state(s) the program will converge. In such cases the LQUANT option can be used to specify which states are desired.

LQUANT,*lam(1),lam(2),... ,lam(nstate)*;

lam(i) is the Λ quantum number of state *i*, i.e., 0 for Σ states, 1 for Π states, 2 for Δ states, etc. The matrix over Λ^2 will be constructed and diagonalized in the P-space configuration basis. The eigenvectors are used to transform the P-space hamiltonian into a symmetry adapted basis, and the program then selects the eigenvectors of the correct symmetry. The states will be ordered by symmetry as specified on the LQUANT card; within each symmetry, the states will be ordered according to increasing energy.

17.5 Restoring and saving the orbitals and CI vectors

MULTI normally requires a starting orbital guess. In this section we describe how to define these orbitals, and how to save the optimized orbitals. In a CASSCF calculation, one has the choice of transforming the final orbitals to natural orbitals (the first order density matrix is diagonalized), to pseudo-canonical orbitals (an effective Fock-operator is diagonalized), or of localizing the orbitals.

17.5.1 Defining the starting guess

START,*record,[options]*;

record: dump record containing starting orbitals. As usual, *record* has the form *irec.ifil*, where *irec* is the record number (e.g., 2140), and *ifil* the file number (usually 2). The *options* can be used to select orbitals of a specific type; for details, see section 2.16.

If this card is missing, the program tries to find suitable starting orbitals as follows:

- | | |
|---------|---|
| First: | Try to read orbitals from the record specified on the ORBITAL card (or the corresponding default, see ORBITAL). All files are searched. |
| Second: | Try to find orbitals from the most recent MCSCF calculation. All files are searched. |
| Third: | Try to find orbitals from the most recent SCF calculation. All files are searched. |

If no orbitals are found, a starting orbital guess is generated.

It is often useful to employ MCSCF orbitals from a neighbouring geometry as starting guess (this will happen automatically if orbitals are found, see the above defaults). Note, however,

that frozen-core orbitals should always be taken from an SCF or MCSCF calculation at the present geometry and must be specified separately on the CORE card. Otherwise the program is likely to stop with error “non-orthogonal core orbitals”. The program remembers where to take the core orbitals from if these have been specified on a CORE card in a previous MCSCF calculation.

17.5.2 Rotating pairs of initial orbitals

ROTATE,*orb1.sym,orb2.sym,angle*

Performs a 2×2 rotation of the initial orbitals *orb1* and *orb2* in symmetry *sym* by *angle* degrees. With *angle*=0 the orbitals are exchanged. ROTATE is meaningful only after the START card. See MERGE for other possibilities to manipulate orbitals.

17.5.3 Saving the final orbitals

ORBITAL,*record.file*

The orbitals are dumped to record *record.file*. Default for *record* is 2140 and *file*=2. This default record number is incremented by one for each subsequent MCSCF calculation in the same job (see section 2.16). Therefore, if several different MCSCF calculations at several geometries are performed in one job, each MCSCF will normally start with appropriate orbitals even if no ORBITAL or START card is present.

The ORBITAL card can be omitted if a NATORB, CANORB or LOCORB card is present, since *orb* can also be specified on these cards (the same defaults for *orb* as above apply in these cases).

17.5.4 Saving the CI vectors and information for a gradient calculation

Old form (obsolete):

SAVE,*cidump,refsav,grdsav*;

New form:

SAVE,[CI=*cidump*,] [REF=*refsav*,] [GRD=*grdsav*];

This directive must be placed before any WF or STATE cards. The options can be given in any order.

cidump: record name for saving the CI vectors. By default, the vectors are only written to a scratch file. If NATORB, CANORB or LOCORB cards are present, *cidump* should be specified on these cards. At present, there is hardly any use of saved CI vectors, and therefore this option is rarely needed.

refsav: record name for saving the orbital configurations and their weights for use in subsequent MULTI or CI calculations using the SELECT directive. If wavefunctions for more than one state symmetry are optimized in a state-averaged calculation, the weights for each state symmetry are saved separately on records $refsav + (istsym - 1) * 100$, where *istsym* is the sequence number of the WF card in the input. If several NATORB, CANORB, or LOCORB cards are present, the record number is increased by 1000 for each subsequent orbital set. Note that this option implies the use of CSFs, even if no CONFIG card (see section 17.6.1) is present.

grdsav: record name for saving the information which is needed in a subsequent gradient calculation. This save is done automatically to record 5000.1 if the input contains a FORCE or OPTG card, and therefore the GRD option is normally not required.

17.5.5 Natural orbitals

NATORB,[*record*,] [*options*]

Request to calculate final natural orbitals and write them to record *record*. The default for *record* is 2140.2, or what else has been specified on an ORBITAL card, if present. By default, the orbitals are not printed and the hamiltonian is not diagonalized for the new orbitals. The following *options* can be specified (in any order):

CI	Diagonalize the hamiltonian in the basis of the computed natural orbitals and print the configurations and their associated coefficients. This has the same effect as the GPRINT , CIVECTOR directive (see section 4.12). By default, only configurations with coefficients larger than 0.05 are printed. This threshold can be modified using the THRESH (see section 17.8.2) or GTHRESH (see section 4.11) options.
STATE= <i>state</i>	Compute natural orbitals for the specified state. <i>state</i> has the form <i>istate.isym</i> , e.g., 3.2 for the third state in symmetry 2. In contrast to earlier versions, <i>isym</i> refers to the number of the irreducible representation, and not the sequence number of the state symmetry. It is therefore independent of the order in which WF cards are given. The specified state must have been optimized. If STATE is not given and two or more states are averaged, the natural orbitals are calculated with the state-averaged density matrix (default).
SPIN= <i>ms2</i>	Compute natural orbitals for states with the specified spin. <i>ms2</i> equals $2 * S$, i.e., 0 for singlet, 1 for doublet etc. This can be used together with STATE to select a specific state in case that states of different spin are averaged. If STATE is not specified, the state-averaged density for all states of the given spin is used.
SAVE= <i>record</i>	Request to save the civector(s) to the specified record.
ORBITAL= <i>record</i>	Request to save the orbitals to the specified record (same effect as specifying <i>record</i> as first argument (see above)).
PRINT= <i>nvirt</i>	Request to print <i>nvirt</i> virtual orbitals in each symmetry. By default, the orbitals are not printed unless the ORBPRINT option (see section 17.8.1) is present or the global GPRINT , ORBITALS (see section 4.12) directive has been given before. The PRINT option on this card applies only to the current orbitals.

Several NATORB , CANORB, and LOCORB cards (for different states) may follow each other. In contrast to earlier versions of MOLPRO the different orbital sets can all be stored in one dump record (but different records still work). See section 2.16 for information about dump records and how specific orbital sets can be requested in a later calculation.

17.5.6 Pseudo-canonical orbitals

CANORB,[*record*,] [*options*]

or

CANONICAL,[*record*,][*options*]

Request to canonicalize the final orbitals, and writing them to record *record*. All options have the same effect as described for NATORB.

17.5.7 Localized orbitals

LOCORB,[*record*,][*options*]

or

LOCAL,[*record*,][*options*]

Request to localize the final orbitals, and writing them to record *record*. All options have the same effect as described for NATORB.

Note: LOCAL is interpreted by MULTI, but LOCALI is a separate command which calls the localization program and not recognized by MULTI. In order to avoid confusion, it is recommended to use LOCORB rather than LOCAL as subcommand within MULTI.

17.5.8 Diabatic orbitals

In order to construct diabatic states, it is necessary to determine the mixing of the diabatic states in the adiabatic wavefunctions. In principle, this mixing can be obtained by integration of the non-adiabatic coupling matrix elements. Often, it is much easier to use an approximate method, in which the mixing is determined by inspection of the CI coefficients of the MCSCF or CI wavefunctions. This method is applicable only if the orbital mixing is negligible. For CASSCF wavefunctions this can be achieved by maximizing the overlap of the active orbitals with those of a reference geometry, at which the wavefunctions are assumed to be diabatic (e.g. for symmetry reasons). The orbital overlap is maximized using the new DIAB command in the MCSCF program. Only the active orbitals are transformed.

This procedure works as follows: first, the orbitals are determined at the reference geometry. Then, the calculations are performed at displaced geometries, and the "diabatic" active orbitals, which have maximum overlap with the active orbitals at the reference geometry, are obtained by adding a DIAB directive to the input:

Old form (Molpro96, obsolete):

DIAB,*orbref*, *orbsav*, *orb1*,*orb2*,*pri*

New form:

DIAB,*orbref*[,TYPE=*orbtype*][,STATE=*state*][,SPIN=*spin*][,MS2=*ms2*][,SAVE=*orbsav*][,ORB1=*orb1*,ORB2=*orb2*][,PRINT=*pri*][,METHOD=*method*]

Here *orbref* is the record holding the orbitals of the reference geometry, and *orbsav* is the record on which the new orbitals are stored. If *orbsav* is not given (recommended!) the new orbitals are stored in the default dump record (2140.2) or the one given on the ORBITAL directive (see section 17.5.3). In contrast to earlier versions of MOLPRO it is possible that *orbref* and *orbsav* are the same. The specifications TYPE, STATE, SPIN can be used to select specific sets of reference orbitals, as described in section 2.16. *orb1*, *orb2* is a pair of orbitals for which the overlap is to be maximized. These orbitals are specified in the form *number.sym*, e.g. 3.1 means the third orbital in symmetry 1. If *orb1*, *orb2* are not given, the overlap of all active orbitals is maximized. *pri* is a print parameter. If this is set to 1, the transformation angles for each orbital are printed for each Jacobi iteration. *method* determines the diabatization method.

method=1 (default): use Jacobi rotations; *method=2*: use block diagonalization. Both methods yield very similar results. *method=2* must only be used for CASSCF wavefunctions. *method=-1* and *method=-2*: as the positive values, but AO overlap matrix of the current geometry is used. This minimizes the change of the MO coefficients, rather than maximizing the overlap to the neighbouring orbitals.

Using the defaults described above, the following input is sufficient in most cases:

```
DIAB,orbref
```

Using `MOLPRO98` is not necessary any more to give any GEOM and DISPL cards. The displacements and overlap matrices are computed automatically (the geometries are stored in the dump records, along with the orbitals).

The diabatic orbitals have the property that the sum of orbital and overlap contributions in the non-adiabatic coupling matrix elements become approximately zero, such that the adiabatic mixing occurs only through changes of the CI coefficients. This allows to determine the mixing angle directly from the CI coefficients, either in a simple way as described for instance in J. Chem. Phys. **89**, 3139 (1988), or in a more advanced manner as described by Pacher, Cederbaum, and Köppel in J. Chem. Phys. **89**, 7367 (1988). Recently, an automatic procedure, as described in J. Chem. Phys. **102**, 0000, (1999) has been implemented into MOLPRO. This is available in Version 99.1 and later and is described in section 28.

Below we present an example for the first two excited states of H₂S, which have B_1 and A_2 symmetry in C_{2v} , and A'' symmetry in C_s . We first perform a reference calculation in C_{2v} symmetry, and then determine the diabatic orbitals for displaced geometries in C_s symmetry. Each subsequent calculation uses the previous orbitals as reference. One could also use the orbitals of the C_{2v} calculation as reference for all other calculations. In this case one would have to take out the second-last input card, which sets `reforb=2141.2`.


```

***,H2S diabatic A" states

basis=VDZ                                !use cc-pVDZ basis set
geometry={x;                               !use Cs symmetry
          planeyz;                         !fix orientation of the molecule
          noorient                         !dont allow automatic reorientation
          s;h1,s,r1;h2,s,r2,h1,theta}     !Z-matrix geometry input

gprint,orbitals,civector                 !global print options

text,reference calculation for C2V
theta=92.12,r1=2.3,r2=2.3                !reference geometry

hf;occ,7,2;wf,18,1;                       !scf calculation for ground state

multi;occ,9,2;closed,4,1;                 !define active and inactive spaces
wf,18,2;state,2;                          !two A" states (1B1 and 1A2 in C2v)
orbital,2140.2                             !save orbitals to 2140.2
reforb=2140.2                               examples/
                                             h2s_diab.com

text,calculations at displaced geometries

rd=[2.4,2.5,2.6]                          !define a range of bond distances

do i=1,#rd                                 !loop over displaced geometries

r2=rd(i)                                   !set r2 to current distance

multi;occ,9,2;closed,4,1;                 !same wavefunction definition as at reference geom.
wf,18,2;state,2;                          !save new orbitals to record
orbital,2141.2                             !compute diabatic orbitals using reference orbitals
diab,reforb                                !stored on record reforb
reforb=2141.2                              !set variable reforb to the new orbitals.
enddo

```

See section 28 for the automatic generation of diabatic energies.

17.6 Selecting the optimization methods

By default, MULTI uses the non-linear optimization method developed by Werner, Meyer, and Knowles. Other methods, such as the Newton-Raphson procedure or the Augmented Hessian procedure, are also implemented and can be selected using the ITERATIONS directive (for state-averaged calculations, only the non-linear optimization method can be used). For CASSCF calculations, the CI problem is solved in a basis of Slater determinants, unless a CONFIG card is given. Some procedures may be disabled using the DONT directive.

17.6.1 Selecting the CI method

CONFIG,*key*;

key may be DET or CSF, and defaults to CSF. If no CONFIG or SELECT card is given, the default is determinants (CASSCF).

17.6.2 Selecting the orbital optimization method

The ITERATIONS directive can be used to modify the defaults for the optimization method. It consists of a sequence of several cards, ending with an END card.

```
ITERATIONS;
DO,method1,iter1[,TO,iter2];
DONT,method2,iter3[,TO,iter4];
...
END;
```

method can be one of the following:

DIAGCI	Diagonalize hamiltonian in the beginning of the specified iterations. This is the default for iteration 1.
INTERNAL	Optimize internal orbitals at the beginning of the specified iterations. This is default for second and subsequent iterations.
WERNER	use Werner-Meyer-Knowles non-linear optimization method for the specified iterations. This is the default for all iterations.
AUGMENT	Use step-restricted Augmented Hessian method for the specified iterations.
NEWTON	Use Newton-Raphson method for specified iterations.
UNCOUPLE	Do not optimize orbitals and CI coefficients simultaneously in the specified iterations. This option will set DIAGCI for these iterations.
NULL	No orbital optimization.

17.6.3 Disabling the optimization

In addition to the ITERATIONS directive described above, some procedures can be disabled more simply using the DONT directive. DONT,*code*

code may be

ORBITAL	Do initial CI but don't optimize orbitals.
WAVEFUNC	Do not optimize the orbitals and CI coefficients (i.e. do only wavefunction analysis, provided the orbitals and CI coefficients are supplied (see START card)).
WVFN	Alias for WAVEFUNC.
ANAL	Do no wavefunction analysis.

17.6.4 Disabling the extra symmetry mechanism

NOEXTRA

This card disables the search for extra symmetries. By default, if extra symmetries are present, each orbital is assigned to such an extra symmetry and rotations between orbitals of different extra symmetry are not performed.

17.6.5 Saving the density matrix

The first order density matrix in AO basis is written automatically to the dump record specified on the ORBITAL card (default 2140.2). If no ORBITAL card is present, but a record is specified on a NATORB, CANORB, or LOCORB card, the densities are saved to the first record occurring in the input. In a state-averaged the SA-density, as well the individual state densities are saved. See section 2.16 for information about how to recover any of these densities for use in later programs.

17.7 Calculating expectation values

By default, the program calculates the dipole expectation and transition moments. Further expectation values or transition properties can be computed using the TRAN, TRAN2 and EXPEC, EXPEC2 directives.

17.7.1 Matrix elements over one-electron operators

```
EXPEC,oper1,oper2,...,opern
TRAN,oper1,oper2,...,opern
```

Calculate expectation values and transition matrix elements for the given one-electron operators. With EXPEC only expectation values are calculated. *oper_i* is a codeword for the operator. The available operators and their associated keywords are given in section 4.13.

17.7.2 Matrix elements over two-electron operators

```
EXPEC2,oper1,oper2,...,opern
TRAN2,oper1,oper2,...,opern
```

Calculate transition matrix elements for two-electron operators. This is presently only useful for angular momentum operators. With EXPEC2 only diagonal matrix elements will be computed. For instance

TRAN2 , LXX	calculates matrix elements for L_x^2
TRAN2 , LYY	calculates matrix elements for L_y^2
TRAN2 , LXZ	calculates matrix elements for $\frac{1}{2}(L_x L_z + L_z L_x)$
TRAN2 , LXX , LYY , LZZ	calculates matrix elements for L_x^2 , L_y^2 , and L_z^2 . The matrix elements for the sum L^2 are also printed.

17.8 Miscellaneous options

All commands described in this section are optional. Appropriate default values are normally used. Note that printing of the orbitals and civectors can also be requested using the global GPRINT command, or by giving NATORB or CANORB options.

17.8.1 Print options

ORBPRINT[,*nvirt*]

requests the occupied and *nvirt* virtual orbitals in each symmetry to be printed (default *nvirt*=0). By default, the program does not print the orbitals, unless the ORBPRINT directive or a global GPRINT, ORBITALS (see section 4.12) command is present. Specific orbital sets can be printed using the PRINT option on a NATORB, CANORB, or LOCORB card (see section 17.5.5). To print additional information at the end of the calculation, use

PRINT,*key1,key2,...*;

Printing is switched on for *key1, key2, . . .*. To print information in each iteration, use

IPRINT,*key1,key2,...*;

Possible print keys are:

MICRO	print details of “microiterations” — useful for finding out what’s going wrong if no convergence
REF	print summary of configuration set (CSFs only)
REF1	print list of configuration set (CSFs only)
COR	print summary of intermediate spaces used in CSF calculation
COR1	print list of intermediate configuration sets (CSFs only)
PSPACE	print list of configurations making up the “primary” space
ORBITALS	print orbitals (see also ORBPRINT)
NATORB	print natural orbitals (see also ORBPRINT)
VIRTUALS	print virtual orbitals (see also ORBPRINT)
CIVECTOR	print CI vector (better use CANORB or NATORB)
INTEGRAL	print transformed integrals (for testing only!)
DENSITY	print density matrices
HESSIAN	print hessian
DIAGONAL	print diagonal elements of hessian
GRADIENT	print gradient
LAGRANGI	print Lagrangian
STEP	print update vector
ADDRESS	print addressing information (for testing only!)
DEBUG	print debugging information
CI2	print debugging information in routine ci2 (Warning: may be long!!)
IO	print debugging information in I/O routines

17.8.2 Convergence thresholds

Convergence thresholds can be modified using

ACCURACY,[GRADIENT=*conv*] [,STEP=*sconv*] [,ENERGY=*econv*]

where

<i>conv</i>	Threshold for orbital gradient (default 10^{-2} .)
<i>econv</i>	Threshold for change of total energy (default 10^{-6}).
<i>sconv</i>	Threshold for size of step (default 10^{-3}).

The default values can be modified using the global GTHRESH command (see section 4.11). Normally, the above default values are appropriate.

17.8.3 Maximum number of iterations

MAXITER,*maxit*;

maxit is maximum number of iterations (default 6). If the calculation does not converge in the default number of iterations, you should first think about the reason before increasing the limit. In most cases the choice of active orbitals or of the optimized states is not appropriate (see introduction of MULTI)

17.8.4 Test options

TEST,*i1, i2, i3, ...*;

Activate testing options numbered *i1, i2, ...*. Please do not use unless you know what you are doing!

17.8.5 Special optimization parameters

STEP,*radius, trust1, tfac1, trust2, tfac2*;

Special parameters for augmented hessian method. For experts only!

GOPER,*igop*;

Use G-operator technique in microiterations (Default). If *igop.lt.0* do not use G-operators.

COPT,*ciacc, copvar, maxci, cishft, icimax, icimx1, icimx2, icstrt, icstep*;

Special parameters for the direct CI method. For experts only!

<i>ciacc</i>	grad threshold for CI diagonalization
<i>copvar</i>	start threshold for CI-optimization
<i>maxci</i>	max. number of CI-optimizations per microiteration
<i>cishft</i>	denominator shift for q-space
<i>icimax</i>	max. number of CI-optimizations in first macroiteration
<i>icimx1</i>	max. number of CI-optimizations in second and subsequent iterations
<i>icimx2</i>	max. number of CI-optimizations in internal absorption step
<i>icstrt</i>	first microiteration with CI-optimization
<i>icstep</i>	microiteration increment between CI-optimizations

INTOPT,*maxito,maxitc,maxrep,nitrep,iuprod*;

Special parameters for internal optimization scheme. For experts only!

NONLINEAR,*itmaxr,ipri,drmax,drdamp,gfak1,gfak2,gfak3,irdamp,ntexp*

Special parameters for non-linear optimization scheme. For experts only!

Old form (obsolete):

THRESH,*thrpri,thrpun,varmin,varmax,thrdiv,thrdoub*

New form:

THRESH [,THRPRI=*thrpri*] [,THRPUN=*thrpun*] [,VARMIN=*varmin*]
[,VARMAX=*varmax*] [,THRDIV=*thrdiv*] [,THRDOU=*thrdoub*]

<i>thrpri</i>	threshold for printing CI coefficients (default 0.04)
<i>thrpun</i>	threshold for writing CI coefficients to the punch file. Default is no write to the punch file
<i>varmin,varmax</i>	thresholds for non-linear optimization scheme. For experts only!
<i>thrdoub</i>	threshold for detecting almost doubly occupied orbitals for inclusion into the pseudo canonical set (default 1.d-7).

DIIS,*disvar,augvar,maxdis,maxaug,idsci,igwgt,igvec,idstrt,idstep*;

Special parameters for DIIS convergence acceleration. For experts only!

17.8.6 Saving wavefunction information for CASVB

VBDUMP[,*vbdump*];

For users of the valence bond program *CASVB*, all wavefunction information that may subsequently be required is saved to the record *vbdump*. The default is not to write this information. If the keyword is specified without a value for *vbdump*, then record 4299.2 is used. This keyword is not needed prior to variational *CASVB* calculations.

17.8.7 Saving transformed integrals

TRNINT,*trnint*;

trnint specifies the record name for integrals in the basis of active CASSCF MOs. These are used for example by *CASVB* (see section 29.5). The default value for *trnint* is 1900.1.

17.9 Coupled-perturbed MCSCF

The coupled-perturbed MCSCF is required for computing gradients with state-averaged orbitals, non-adiabatic couplings, difference gradients or polarizabilities. We note that the present implementation is somewhat preliminary and not very efficient.

17.9.1 Gradients for SA-MCSCF

For computing state-averaged gradients, use

```
CPMCSCF , GRAD,state,[SPIN=spin],[MS2=ms2],[ACCU=thresh],[RECORD=record]
```

where *state* specifies the state (e.g., 2.1 for the second state in symmetry 1) for which the gradients will be computed. *spin* specifies the spin of the state: this is half the value used in the corresponding WF card (e.g., 0=Singlet, 0.5=Doublet, 1=Triplet). Alternatively, MS2 can be used, where $ms2 = 2 * spin$, i.e., the same as specified on WF cards. The specification of SPIN or MS2 is only necessary if states with different spin are state-averaged. *record* specifies a record on which the gradient information is stored (the default is 5101.1). *thresh* is a threshold for the accuracy of the CP-MCSCF solution. The default is 1.d-7.

The gradients are computed by a subsequent call to FORCES or OPTG.

Note: if for some reason the gradients are to be computed numerically from finite energy differences, it is in state-averaged calculations necessary to give, instead of the CPMCSCF input, the following:

```
SAVE , GRAD=-1
```

Otherwise the program will stop with an error message.

17.9.2 Difference gradients for SA-MCSCF

For computing difference gradients, use

```
CPMCSCF , DGRAD,state1,state2,[ACCU=thresh],[RECORD=record]
```

where *state1* and *state2* specify the two states considered. (e.g., 2.1,3.1 for the second and third states in symmetry 1) The gradient of the energy difference will be computed. Both states must have the same symmetry. *record* specifies a record on which the gradient information is stored (the default is 5101.1). *thresh* is a threshold for the accuracy of the CP-MCSCF solution. The default is 1.d-7.

The gradients are computed by a subsequent call to FORCES or OPTG.

17.9.3 Non-adiabatic coupling matrix elements for SA-MCSCF

For computing non-adiabatic coupling matrix elements analytically, use

```
CPMCSCF , NACM,state1,state2,[ACCU=thresh],[RECORD=record]
```

where *state1* and *state2* specify the two states considered. (e.g., 2.1,3.1 for the second and third states in symmetry 1) Both states must have the same symmetry. *record* specifies a record on which the gradient information is stored (the default is 5101.1). This will be read in the subsequent gradient calculation. *thresh* is a threshold for the accuracy of the CP-MCSCF solution. The default is 1.d-7.

NADC and NADK are an aliases for NADC, and SAVE is an alias for RECORD.

The matrix elements for each atom are computed by a subsequent call to FORCES.

Note: this program is not yet extensively tested and should be used with care!

17.10 Optimizing valence bond wavefunctions

VB={ . . . }

Using this keyword, the optimization of the CI coefficients is carried out by *CASVB*. The VB keyword can be followed by any of the directives described in section 29. Energy-based optimization of the VB parameters is the default, and the output level for the main *CASVB* iterations is reduced to -1.

17.11 Hints and strategies

MCSCF is not a “black box” procedure like SCF! For simple cases, for example a simple CASSCF with no CLOSED orbitals, this program will converge in two or three iterations. For more complicated cases, you may have more trouble. In that case, consider the following:

- Always start from neighbouring geometry orbitals when available (this is the default).
- The convergence algorithm is more stable when there are no CLOSED orbitals, i.e., orbitals doubly occupied in all configurations, but fully optimized. Thus a reasonable approach is to make an initial calculation with CLOSED replaced by CORE (all doubly occ. frozen).
- If still no success, you can switch off the coupling between CI coefficients and orbital rotations for a few iterations, e.g.:

```
ITERATIONS; UNCOUPLE, 1, TO, 2; END;
```

and/or disable the simultaneous optimization of internal orbitals & CI, e.g.:

```
ITERATIONS; DONT, INTERNAL, 1, TO, 2; END;
```

You can often get a clue about where the program starts to diverge if you include:

```
IPRINT, MICRO;
```

in the data. Also consider the general remarks at the beginning of this chapter. For the details of the algorithms used, see J. Chem. Phys 82, 5053 (1985); Chem. Phys. Letters 115, 259 (1985); Advan. Chem. Phys. 59, 1 (1987);

17.12 Examples

The simplest input for a CASSCF calculation for H₂O, C_{2v} symmetry, is simply:

```
geometry={o;h1,o,r;h2,o,r,h1,theta}    !Z-matrix geometry input
r=1 ang                                !bond length
theta=104                               !bond angle
hf                                       !do scf calculation
multi                                   !do full valence casscf
```

examples/
h2o_casscf.com

This could be extended, for instance, by the following input cards

```
OCC,4,1,2;          ! specify occupied space
CLOSED,2            ! specify closed-shell (inactive) orbitals
CORE,1;             ! specify frozen core orbitals
WF,10,1;           ! define wavefunction symmetry
```



```

START,2100.2;      ! read guess orbitals from record 2100, file 2
ORBITAL,2140.2;   ! save final orbitals to record 2140, file 2
NATORB,PRINT,CI  ! print natural orbitals and diagonalize the hamiltonian
                  ! for the natural orbitals. The largest CI coefficients
                  ! are printed.

```

Example for a state-averaged calculation for CN, X and $B^2\Sigma^+$ states, and $A^2\Pi_x$, $^2\Pi_y$ states averaged. A full valence CASSCF calculation is performed

```

***,cn
r=2.2                !define bond length
geometry={c;n,c,r}
rhf;occ,5,1,1;wf,13,1,1;          !RHF calculation for sigma state
orbital,2100.2          !save orbitals to record 2100.2 (default)

multi;occ,6,2,2;closed,2;        !Define active and inactive orbitals
start,2100.2;                !Start with RHF orbitals from above
save,ref=4000.2              !Save configuration weights for CI in record 4000.2
wf,13,1,1;state,2;wf,13,2,1;wf,13,3,1; !Define the four states
natorb,ci,print;            !Print natural orbitals and associated ci-coefficients
tran,lz                      !Compute matrix elements over LZ
expec2,lzz                   !compute expectation values for LZZ

```

Example for an RASSCF (restricted active space) calculation for N_2 , including SCF determinant plus all double excitations into valence orbitals. The single excitations are excluded. D_{2h} symmetry, CSF method used:

```

***,N2
geometry={N1;N2,N1,r}          !geometry input
r=2.2                          !bond length
hf;occ,3,1,1,,2;wf,14,1;save,2100.2 !scf calculation

multi;occ,3,1,1,,3,1,1;        !Define occupied orbitals
core,1,,,,1,2100.2;            !Define frozen core scf orbitals
config;                          !Use CSF method
wf,14,1;                        !Define state symmetry
restrict,0,2,3.5,1.6,1.7;      !Restriction to singles and doubles
restrict,-1,-1,3.5,1.6,1.7;    !Take out singles
print,ref1                      !Print configurations
natorb,ci,print;                !Print natural orbitals and CI coeffs

```

18 THE CI PROGRAM

Multiconfiguration reference internally contracted configuration interaction

Bibliography:

H.-J. Werner and P.J. Knowles, J. Chem. Phys. 89, 5803 (1988).

P.J. Knowles and H.-J. Werner, Chem. Phys. Lett. 145, 514 (1988).

All publications resulting from use of this program must acknowledge the above. See also:

H.-J. Werner and E.A. Reinsch, J. Chem. Phys. 76, 3144 (1982).

H.-J. Werner, Adv. Chem. Phys. 59, 1 (1987).

The command `CI` or `CI-PRO` calls the program. The command `CISD` calls fast closed-shell CISD program. The command `QCI` calls closed-shell quadratic CI program. The command `CCSD` calls closed-shell coupled-cluster program.

18.1 Introduction

The internally contracted MRCI program is called by the `CI` command. This includes as special cases single reference CI, CEPA, ACPF, MR-ACPF and MR-AQCC. For closed-shell reference functions, a special faster code exists, which can be called using the `CISD`, `QCI`, or `CCSD` commands. This also allows to calculate Brueckner orbitals for all three cases (`QCI` and `CCSD` are identical in this case).

With no further input cards, the wavefunction definition (core, closed, and active orbital spaces, symmetry) corresponds to the one used in the most recently done SCF or MCSCF calculation. By default, a CASSCF reference space is generated. Other choices can be made using the `OCC`, `CORE`, `CLOSED`, `WF`, `SELECT`, `CON`, and `RESTRICT` cards. The orbitals are taken from the corresponding SCF or MCSCF calculation unless an `ORBITAL` directive is given. The wavefunction may be saved using the `SAVE` directive, and restarted using `START`. The `EXPEC` directive allows to compute expectation values over one-electron operators, and the `TRAN` directive can be used to compute transition matrix elements for one-electron properties. Natural orbitals can be printed and saved using the `NATORB` directive.

For excited state calculations see `STATE`, `REFSTATE`, and `PROJECT`.

18.2 Specifying the wavefunction

18.2.1 Occupied orbitals

`OCC,n1,n2,...,n8;`

n_i specifies numbers of occupied orbitals (including `CORE` and `CLOSED`) in irreducible representation number i . If not given, the information defaults to that from the most recent SCF, MCSCF or CI calculation.

18.2.2 Frozen-core orbitals

`CORE,n1,n2,...,n8;`

n_i is the number of frozen-core orbitals in irrep number i . These orbitals are doubly occupied in all configurations, i.e., not correlated. If no CORE card is given, the program uses the same core orbitals as the last CI calculation; if there was none, then the atomic inner shells are taken as core. To avoid this behaviour and correlate all electrons, specify

CORE

18.2.3 Closed-shell orbitals

CLOSED, n_1, n_2, \dots, n_8

n_i is the number of closed-shell orbitals in irrep number i , inclusive of any core orbitals. These orbitals do not form part of the active space, i.e., they are doubly occupied in all reference CSFs; however, in contrast to the core orbitals (see CORE), these orbitals are correlated through single and double excitations. If not given, the information defaults to that from the most recent SCF, MCSCF or CI calculation. For calculations with closed-shell reference function (closed=occ), see CISD, QCI, and CCSD.

18.2.4 Defining the orbitals

ORBIT,*name.file*,[*specifications*];

name.file specifies the record from which orbitals are read. Optionally, various *specifications* can be given to select specific orbitals if *name.file* contains more than one orbital set. For details see section 2.16.

The default is the set of orbitals from the last SCF, MCSCF or CI calculation.

18.2.5 Defining the state symmetry

The number of electrons and the total symmetry of the wavefunction are specified on the WF card:

WF,*elec,sym,spin*

where

<i>elec</i> :	is the number of electrons
<i>sym</i> :	is the number of the irreducible representation
<i>spin</i> :	defines the spin symmetry, $spin = 2S$ (singlet=0, doublet=1, triplet=2, etc.)

The WF card must be placed after any cards defining the orbital spaces (*OCC*, *CORE*, *CLOSED*).

The *REF* card can be used to define further reference symmetries used for generating the configuration space, see *REF*.

18.3 Additional reference symmetries

REF,*sym*;

This card, which must come after the WF directive, defines an additional reference symmetry used for generating the uncontracted internal and singly external configuration spaces. This is sometimes useful in order to obtain the same configuration spaces when different point group symmetries are used. For instance, if a calculation is done in C_s symmetry, it may happen that the two components of a Π state, one of which appears in A' and the other in A'' , come out not exactly degenerate. This problem can be avoided as in the following example:

for a doublet A' state:

```
WF,15,1,1;      !define wavefunction symmetry (1)
REF,2;          !define additional reference symmetry (2)
```

and for the doublet A'' state:

```
WF,15,2,1;      !define wavefunction symmetry (2)
REF,1;          !define additional reference symmetry (1)
```

For linear geometries the same results can be obtained more cheaply using C_{2v} symmetry,

```
WF,15,2,1;      !define wavefunction symmetry (2)
REF,1;          !define additional reference symmetry (1)
REF,3;          !define additional reference symmetry (3)
```

or

```
WF,15,3,1;      !define wavefunction symmetry (2)
REF,1;          !define additional reference symmetry (1)
REF,2;          !define additional reference symmetry (2)
```

Each REF card may be followed by RESTRICT, SELECT, and CON cards, in the given order.

18.3.1 Selecting configurations

SELECT,*ref1,ref2,refthr,refstat,mxshrf*;

This card is used to specify a reference configuration set other than a CAS, which is the default. Configurations can be defined using CON cards, which must appear after the *SELECT* card. Alternatively, if *ref1* is an existing MOLPRO record name, the configurations are read in from that record and may be selected according to a given threshold.

The select card should normally be placed directly after the *WF* or *REF* card(s), or, if present, the RESTRICT cards. The general order of these cards is

```
WF (or REF)
RESTRICT (optional)
SELECT (optional)
CON (optional)
```

<i>ref1=rec1.file:</i>	(<i>rec1</i> >2000) The configurations are read in from the specified record. See section 17.5.4 about how to save the configurations in the MCSCF calculation. If <i>ref1</i> is not specified, the program assumes that the configurations are read from subsequent CON cards (see CON).
<i>ref2=rec2.file:</i>	(<i>rec2</i> >2000) additional configurations are read from the specified record. If <i>rec2</i> is negative, all records between <i>rec1</i> and <i>abs(rec2)</i> are read. All configurations found in this way are merged.
<i>refthr:</i>	Selection threshold for configurations read from disc (records <i>rec1</i> – <i>rec2</i>). This applies to the norm of all CSFs for each orbital configuration.
<i>refstat:</i>	Specifies from which state vector the configurations are selected. This only applies to the case that the configurations were saved in a state-averaged calculation. If <i>refstat</i> is zero or not specified, the configurations are selected from all states. If <i>refstat</i> is greater than zero, then the specified reference state is used. If <i>refstat</i> is less than zero, then all appropriate reference states are used. Lastly, if <i>refstat</i> is of the form <i>istat1.istat2</i> , states <i>istat1</i> through <i>istat2</i> are used.
<i>mxshrf:</i>	maximum number of open shells in the selected or generated configurations.

18.3.2 Occupation restrictions

RESTRICT,*nmin,nmax,orb₁,orb₂,...orb_n*;

This card can be used to restrict the occupation patterns in the reference configurations. Only configurations containing between *nmin* and *nmax* electrons in the specified orbitals *orb₁, orb₂, ..., orb_n* are included in the reference function. If *nmin* and *nmax* are negative, configurations with exactly *abs(nmin)* and *abs(nmax)* electrons in the specified orbitals are deleted. This can be used, for instance, to omit singly excited configurations. The orbitals are specified in the form *number.sym*, where *number* is the number of the orbital in *irrep sym*. Several RESTRICT cards may follow each other.

The RESTRICT cards must follow the WF or REF cards to which they apply. The general order of these cards is

```
WF (or REF)
RESTRICT (optional)
SELECT (optional)
CON (optional)
```

If a RESTRICT cards precedes the WF card, it applies to all reference symmetries. Note that RESTRICT also affects the spaces generated by SELECT and/or CON cards.

18.3.3 Explicitly specifying reference configurations

CON,*n₁,n₂,n₃,n₄,...*

Specifies an orbital configuration to be included in the reference function. n_1, n_2 etc. are the occupation numbers of the active orbitals (0,1,or 2). Any number of CON cards may follow each other, but they must all appear directly after a SELECT card.

18.3.4 Defining state numbers

```
STATE,nstate,nroot(1),nroot(2),...,nroot(nstate);
```

nstate is the number of states treated simultaneously; *nroot(i)* are the root numbers to be calculated. These apply to the order of the states in the initial internal CI. If not specified, *nroot(i)=i*. Note that it is possible to leave out states, i.e.,

```
STATE,1,2;    ! calculates second state
STATE,2,1,3;  ! calculates first and third state
```

All states specified must be reasonably described by the internal configuration space. It is possible to have different convergence thresholds for each state (see ACCU card). It is also possible not to converge some lower roots which are included in the list *nroot(i)* (see REFSTATE card). For examples, see REFSTATE card.

18.3.5 Defining reference state numbers

```
REFSTATE,nstatr,nrootr(1),nrootr(2),...,nrootr(nstatr);
```

nstatr is the number of reference states for generating contracted pairs. This may be larger or smaller than *nstate*. If this card is not present, *nstatr=nstate* and *nrootr(i)=nroot(i)*. Roots for which no reference states are specified but which are specified on the STATE card (or included by default if the *nroot(i)* are not specified explicitly on the STATE card) will not be converged, since the result will be bad anyway. However, it is often useful to include these states in the list *nroot(i)*, since it helps to avoid root flipping problems. Examples:

```
state,2;
```

will calculate two states with two reference states.

```
state,2;refstate,1,2;
```

will optimize second state with one reference state. One external expansion vector will be generated for the ground state in order to avoid root flipping. The results printed for state 1 are bad and should not be used (unless the pair space is complete, which might happen in very small calculations).

```
state,1,2;refstate,1,2;
```

As the second example, but no external expansion vectors will be generated for the ground state. This should give exactly the same energy for state 2 as before if there is no root flipping (which, however, frequently occurs).

```
state,2;accu,1,1,1;
```

Will calculate second state with two reference states. The ground state will not be converged (only one iteration is done for state 1) This should give exactly the same energy for state 2 as the first example.

18.3.6 Specifying correlation of orbital pairs

PAIR,*iorb1.isy1,iorb2.isy2,np*;

is a request to correlate a given orbital pair.

np=1: singlet pair
np=-1: triplet pair
np=0: singlet and triplet pair (if possible)

Default is to correlate all electron pairs in active and closed orbitals. See also PAIRS card.

PAIRS,*iorb1.isy,iorb2.isy,np*;

Correlate all pairs which can be formed from orbitals *iorb1.isy1* through *iorb2.isy2*. Core orbitals are excluded. Either *iorb2* must be larger than *iorb1* or *isy2* larger than *isy1*. If *iorb1.isy1=iorb2.isy2* the PAIRS card has the same effect as a PAIR card. PAIR and PAIRS cards may be combined.

If no PAIR and no PAIRS card is specified, all valence orbitals are correlated. The created pair list restricts not only the doubly external configurations, but also the all internal and semi internals.

18.3.7 Restriction of classes of excitations

NOPAIR ;

No doubly external configurations are included.

NOSINGLE ;

No singly external configurations are included.

NOEXC ;

Perform CI with the reference configurations only.

18.4 Options

18.4.1 Coupled Electron Pair Approximation

CEPA,*ncepa*;

Instead of diagonalizing the hamiltonian, perform CEPA calculation, CEPA type *ncepa*. This is currently available only for single configuration reference functions.

18.4.2 Coupled Pair Functional (ACPF, AQCC)

CPF,*ncpf,gacpfi,gacpfe*;

ACPF,*ncpf,gacpfi,gacpfe*;

AQCC,*ncpf,gacpfi,gacpfe*;

Instead of diagonalizing the hamiltonian, perform CPF calculation (*ncpf=2*) (not yet implemented) ACPF calculation (*ncpf=0*) or AQCC calculation (*ncpf=1*). For ACPF and AQCC, the internal and external normalization factors *gacpfi*, *gacpfe* may be reset from their default values of 1, $2/nelec$ and 1, $1-(nelec-2)(nelec-3)/nelec(nelec-1)$, respectively.

18.4.3 Projected excited state calculations

PROJECT,*record*,*nprojc*;

Initiate or continue a projected excited state calculation, with information stored on *record*. If *nprojc* > 0, the internal CI vectors of *nprojc* previous calculations are used to make a projection operator. If *nprojc* = -1, this calculation is forced to be the first, i.e. ground state, with no projection. If *nprojc* = 0, then if *record* does not exist, the effect is the same as *nprojc* = -1; otherwise *nprojc* is recovered from the dump in *record*. Thus for the start up calculation, it is best to use `project,record,-1`; for the following excited calculations, use `project,record`; At the end of the calculation, the wavefunction is saved, and the information in the dump *record* updated. The project card also sets the `tranh` option, so by default, transition hamiltonian matrices are calculated.

For example, to do successive calculations for three states, use

```
ci;...;project,3000.3,-1;
ci;...;project,3000.3;
ci;...;project,3000.3;
```

18.4.4 Transition matrix element options

TRANH,*option*;

If *option* > -1, this forces calculation of transition hamiltonian matrix elements in a TRANS or PROJECT calculation. If *option* < 1, this forces calculation of one electron transition properties.

18.4.5 Convergence thresholds

ACCU,*istate*,*energy*,*coeff*;

Convergence thresholds for state *istate*. The actual thresholds for the energy and the CI coefficients are $10^{**}(-energy)$ and $10^{**}(-coeff)$. If this card is not present, the thresholds for all states are the default values or those specified on the THRESH card.

18.4.6 Level shifts

SHIFT,*shiftp*,*shifts*,*shifti*;

Denominator shifts for pairs, singles, and internals, respectively.

18.4.7 Maximum number of iterations

MAXITER,*maxit*,*maxiti*;

maxit: maximum number of macroiterations;
maxiti: maximum number of microiterations (internal CI).

18.4.8 Restricting numbers of expansion vectors

MAXDAV,*maxdav*,*maxvi*;

maxdav: maximum number of external expansion vectors in macroiterations;

maxvi: maximum number of internal expansion vectors in internal CI.

18.4.9 Selecting the primary configuration set

PSPACE,*select*,*npspac*;

select: energy criterion for selecting p-space configurations. If negative, a test for p-space H is performed.

npspac: minimum number of p-space configurations. Further configurations are added if either required by *select* or if configurations are found which are degenerate to the last p-space configuration. A minimum number of *npspace* is automatically determined from the state specifications.

18.4.10 Canonicalizing external orbitals

FOCK,*n*₁,*n*₂,...;

External orbitals are obtained as eigenfunctions of a Fock operator with the specified occupation numbers *n_i*. Occupation numbers must be provided for all valence orbitals.

18.4.11 Saving the wavefunction

SAVE,*savecp*,*saveco*,*idelcg*;

or

SAVE [,CIVEC=*savecp*] [,CONFIG=*saveco*] [,DENSITY=*dumprec*] [,NATORB=*dumprec*] [,FILES]

savecp: record name for save of wavefunction. If negative the wavefunction is saved after each iteration, else at the end of the job. In case of coupled cluster methods (CCSD, QCISD, BCCD), the wavefunction is saved in each iteration in any case (presently only implemented for the closed-shell case).

saveco: record name for save of internal configurations and their maximum weight over all states for subsequent use as reference input (see SELECT card). If the record already exists, the record name is incremented by one until a new record is created.

idelcg: if nonzero or FILES is specified, don't erase icfil and igfil (holding CI and residual vectors) at the end of the calculation.

dumprec: Dump record for saving density matrix and natural orbitals. Only one dump record must be given. In any case the density matrix and the natural orbitals are saved. See also DM or NATORB cards.

18.4.12 Starting wavefunction

START,*readc1*,*irest*;

readc1: record name from which the wavefunction is restored for a restart. In the case of coupled cluster methods (CCSD, QCISD, BCCD), the amplitudes are read from record *readc1* and used for restart (presently only implemented for closed-shell methods)

irest: If nonzero, the CI coefficients are read and used for the restart; otherwise, only the wavefunction definition is read in.

18.4.13 One electron properties

EXPEC,*oper1*,*oper2*,*oper3*,...;

After the wavefunction determination, calculate expectation values for one-electron operators *oper_i*. See section 4.13 for the available operators and their keywords. In multi-state calculations or in projected calculations, also the transition matrix elements are calculated.

18.4.14 Transition moment calculations

TRANS,*readc1*,*readc2*,*oper1*,*oper2*,*oper3*,...;

Instead of performing an energy calculation, only calculate transition matrix elements between wavefunctions saved on records *readc1* and *readc2*. See section 4.13 for a list of available operators and their corresponding keywords. If no operator names are specified, the dipole transition moments are calculated.

18.4.15 Saving the density matrix

DM,*record.ifil*,[*idip*];

The first order density matrices for all computed states are stored in record *record* on file *ifil*. If *idip* is not zero, the dipole moments are printed starting at iteration *idip*. See also NATORB. In case of transition moment calculation, the transition densities are also stored, provided both states involved have the same symmetry.

18.4.16 Natural orbitals

NATORB,[RECORD=]*record.ifil*,[PRINT=*nprint*],[CORE[=*natcor*]];

Calculate natural orbitals. The number of printed external orbitals in any given symmetry is *nprint* (default 2). *nprint=-1* suppressed the printing. If *record* is nonzero, the natural orbitals and density matrices for all states are saved in a dump record *record* on file *ifil*. If *record.ifil* is specified on a DM card (see above), this record is used. If different records are specified on the DM and NATORB cards, an error will result. The record can also be given on the SAVE card. If CORE is specified, core orbitals are not printed.

Note: The dump record must not be the same as *savecp* or *saveco* on the SAVE card, or the record given on the PROJECT.

18.4.17 Miscellaneous options

OPTION,*code1=value,code2=value,...*

Can be used to specify program parameters and options. If no codes and values are specified, active values are displayed. The equal signs may be omitted. The following codes are allowed (max 7 per card):

NSTATE :	see state card
NSTATI :	number of states calculated in internal CI
NSTATR :	see refstat card
NCEPA :	see CEPA card
NOKOP :	if nonzero, skip integral transformation
ITRDM :	if .ge. 0 transition moments are calculated
ITRANS :	if nonzero, perform full integral transformation (not yet implemented)
IDIP :	Print dipole moments from iteration number <i>value</i>
REFOPT :	if nonzero, optimize reference coefficients; otherwise extract reference coefficients from internal CI
IAVDEN :	average HII and HSS denominators over spin couplings if nonzero
IDELCG :	if.ne.0 then destroy files icfil,igfil at end
IREST :	if nonzero, restart
NATORB :	if nonzero, natural orbitals are calculated and printed. The number of printed external orbitals per symmetry is $\min(natorb,2)$
WFNAT :	if nonzero, natural orbitals are saved to this record
IPUNRF :	if nonzero, punch coefficients of reference configurations
NPUPD :	if nonzero, update pairs in nonorthogonal basis, otherwise in orthogonal basis.
MAXIT :	see maxiter card
MAXITI :	see maxiter card
MAXDAV :	see maxdav card
MAXVI :	see maxdav card
NOSING :	see nosing card
NOPAIR :	see nopair card
MXSHRF :	see select card
IKCPS=0 :	In CIKEXT, only K(CP) is calculated; this option taken when and only when no singles.
IKCPS=1 :	only K(CP') is calculated. Implies that modified coupling coefficients are used.
IKCPS=2 :	K(CP) and K(CP') are calculated. Default is IKCPS=2 except when single reference configuration, when IKCPS=1.
IOPTGM :	Option for density matrix routines.
IOPTGM=0 :	all quantities in density matrix routines are recalculated for each intermediate symmetry (max. CPU, min. core).

IOPTGM=1 :	quantities precalculated and stored on disk (max. I/O, min. core).
IOPTGM=2 :	quantities precalculated and kept in core (min. CPU, max. core).
IOPTOR :	If nonzero, calculate intermediate orbitals for each pair. Might improve convergence in some cases, in particular if localized orbitals are used.

18.4.18 Miscellaneous parameters

PARAM,*code1=value,code2=value...*

Redefine system parameters. If no codes are specified, the default values are displayed. The following codes are allowed:

LSEG :	disc sector length
INTREL :	number of integers per REAL*8 word
IVECT=0 :	scalar machine
IVECT=1 :	vector machine
MINVEC :	call MXMB in coupling coefficient routines if vector length larger than this value.
IBANK :	number of memory banks for vector machines. If IBANK>1, vector strides which are multiples of IBANK are avoided where appropriate.
LTRACK :	number of REAL*8 words per track or block (for file allocation)
LTR :	determines how matrices are stored on disc. If LTR=LSEG, all matrices start at sector boundaries (which optimizes I/O), but unused space is between matrices (both on disc and in core). With LTR=1 all matrices are stored dense. This might increase I/O if much paging is necessary, but reduce I/O if everything fits in core.
NCPUS :	Maximum number of CPUs to be used in multitasking.

18.5 Miscellaneous thresholds

THRESH,*code1=value,code2=value...*

If *value=0*, the corresponding threshold is set to zero, otherwise $10^{*(-value)}$. The equal signs may be omitted. If no codes are specified, the default values are printed. The following codes are allowed (max 7 per card):

ZERO :	numerical zero
THRDLP :	delete pairs if eigenvalue of overlap matrix is smaller than this threshold.
PNORM :	delete pair if its norm is smaller than this threshold (all pairs are normalized to one for a closed shell case).

PRINT :	print CI coefficients which are larger than this value.
INTEG :	omit two-electron integrals which are smaller than this value.
ENERGY :	convergence threshold for energy; see also: ACCU card.
COEFF :	convergence threshold for coefficients; see also: ACCU card.
SPARSE :	omit coefficient changes which are smaller than this value.
EQUAL :	set values in the internal vector and the diagonal elements equal if they differ by less than this value. Useful for keeping track of symmetry.

18.6 Print options

PRINT,*code1=value,code2=value,...*

Print options. Generally, the value determines how much intermediate information is printed. *value=-1* means no print (default for all codes). In some of the cases listed below the specification of higher values will generate even more output than described. The equal signs and zeros may be omitted. All codes may be truncated to three characters. The following codes are allowed (max 7 per card):

ORBITALS :	print orbitals
JOP=0 :	print operator list
JOP=1 :	print coulomb operators in MO basis
JOP=2 :	print coulomb operators in AO and MO basis
KOP :	as JOP for internal exchange operators
KCP=0 :	print paging information for CIKEXT
KCP=1 :	print external exchange operators in MO basis
KCP=2 :	print operators in AO and MO basis
DM=0 :	print paging information for CIDIMA
DM=1 :	print density matrix in MO basis
DM=2 :	print density matrix in AO and MO basis
FPP=0 :	print energy denominators for pairs
FPP=1 :	in addition, print diagonal coupling coefficients in orthogonal basis.
FPP=2 :	print operators FPP
CP=0 :	print update information for pairs in each iteration
CP=1 :	print pair matrix updates (MO basis)
CP=2 :	in addition print pair matrices (MO basis)
CP=3 :	print CP in AO basis (in CIKEXT)
CI=0 :	print convergence information for internal CI
CI=1 :	print internal CI coefficients and external expansion coefficients
CS :	as CP for singles
CPS=0 :	print paging information for CICPS
CPS=1 :	print matrices CPS in MO basis

GPP=0 : print paging information for CIGPQ
 GPP=1 : print matrices GP at exit of CIGPQ
 GPS=0 : print paging information for CIGPS
 GPS=1 : print vectors GS at exit CIGPS
 GSP=1 : print matrices GP at exit CIGPS
 GPI=0 : print paging information for CIGPI
 GPI=1 : print total GP in orthogonal basis
 GPI=2 : print matrices GP and TP
 GIP=0 : print paging information for CIGIP
 GIP=1 : print GI at exit CIGIP
 GSS=0 : print paging information for CIGSS
 GSS=1 : print vectors GS at exit CIGSS
 GSI=0 : print paging information for CIGSI
 GSI=1 : print GS at exit CIGSI
 GIS=0 : print paging information for CIGIS
 GIS=1 : print GI at exit CIGIS
 GII : print intermediate information in internal CI
 DPQ : print coupling coefficients $\alpha(P, Q)$
 EPQ : print coupling coefficients $\beta(P, Q)$
 HPQ : print coupling coefficients $\gamma(P, Q)$
 DPI : print coupling coefficients for pair-internal interactions
 DSS : not yet used
 DSI : not yet used
 LOG : At end of each iteration, write summary to log file. Delete at
 end of job if LOG=0
 CC=0 : print address lists for coupling coefficients
 CC=1 : print coupling coefficients
 DEN=1 : print internal first order density
 DEN=2 : print internal second order density
 DEN=3 : print internal third order density
 DEN=4 : print first, second and third order densities
 GAM=1 : print first order transition densities
 GAM=2 : print second order transition densities
 GAM=3 : print first and second order transition densities
 PAIRS=0 : print list of non redundant pairs
 PAIRS=1 : print list of all pairs
 CORE=0 : print summary of internal configurations ($N, N-1$ and $N-2$
 electron)
 CORE=1 : print internal configurations ($N, N-1, N-2$)
 REF=0 : print summary of reference configurations

```

REF=1:          print reference configurations and their coefficients
PSPACE:        print p-space configurations
HII:           print diagonal elements for internals
HSS:           print diagonal elements for singles
SPQ:           various levels of intermediate information in pair orthogonal-
                ization routine.

TEST=0:        print information at each subroutine call
TEST=1:        print in addition information about I/O in LESW, SREIBW
TEST=2:        print also information about I/O in FREAD, FWRITE
CPU:           print analysis of CPU and I/O times
ALL:           print everything at given level (be careful!)

```

18.7 Examples

```

***,Single reference CISD and CEPA-1 for water

```

```

r=0.957,angstrom
theta=104.6,degree;
geometry={0;          !z-matrix geometry input
          H1,O,r;
          H2,O,r,H1,theta}
hf;wf,10,1;          !TOTAL SCF ENERGY      -76.02680642
ci;occ,3,1,1;core,1;wf,10,1;          !TOTAL CI (SD) ENERGY -76.22994348
ci;occ,3,1,1;core,1;wf,10,1;cepa,1; !TOTAL CEPA(1) ENERGY -76.23799334

```

examples/
h2o_cep1.com

```

***,Valence multireference CI for X and A states of H2O

```

```

gthresh,energy=1.d-8
r=0.957,angstrom,theta=104.6,degree;
geometry={0;          !z-matrix geometry input
          H1,O,r;
          H2,O,r,H1,theta}
hf;wf,10,1;          !TOTAL SCF ENERGY      -76.02680642
multi;occ,4,1,2;closed,2;core,1;wf,9,2,1;wf,9,1,1;tran,ly;
          !MCSCF ENERGY      -75.66755631
          !MCSCF ENERGY      -75.56605896
ci;occ,4,1,2;closed,2;core,1;wf,9,2,1;save,7300.1;
          !TOTAL MRCI ENERGY      -75.79831209
ci;occ,4,1,2;closed,2;core,1;wf,9,1,1;save,7100.1;
          !TOTAL MRCI ENERGY      -75.71309879
ci;trans,7300.1,7100.1,ly;
          !Transition moment <1.3|X|1.1> = -0.14659810 a.u.
          !Transition moment <1.3|LY|1.1> = 0.96200488i a.u.

```

examples/
h2o_mrci_trans.com

```
***,BH singlet Sigma and Delta states
r=2.1
geometry={b;h,b,r}
hf;occ,3;wf,6,1;
multi;
occ,3,1,1;core,1;wf,6,1;state,3;lquant,0,2,0;wf,6,4;lquant,2;
tran,lz;
expec2,lz;lz;
! Sigma states -- energies -25.20509620 -24.94085861
ci;occ,3,1,1;core,1;wf,6,1;state,2,1,3;
! Delta states -- energies -24.98625171
ci;occ,3,1,1;core,1;wf,6,1;state,1,2;
! Delta state -- xy component
ci;occ,3,1,1;core,1;wf,6,4;
```

examples/
bh_mrci_sigma_delta.com

19 MULTIREFERENCE RAYLEIGH SCHRÖDINGER PERTURBATION THEORY

Bibliography:

Original RS2/RS3:

H.-J. Werner, Mol. Phys. 89, 645-661 (1996)

New internally contracted RS2C:

P. Celani and H.-J. Werner, J. Chem. Phys. 112, 5546 (2000)

All publications resulting from use of this program must acknowledge the above.

The commands

RS2

RS2C

RS3

are used to perform second or third-order perturbation calculations. RS3 always includes RS2 as a first step. For closed-shell single-reference cases, this is equivalent to MP2 or MP3 (but a different program is used). RS2C calls a new more efficient second-order program (see below), which should normally be used if third-order is not required (note that RS3C is not available).

19.1 Introduction

Multireference perturbation calculations are performed by the MRCI program as a special case. For RS2 (CASPT2,RASPT2) only matrix elements over a one-electron operator need to be computed, and therefore the computational effort is much smaller than for a corresponding MRCI. For RS3 (CASPT3) the energy expectation value for the first-order wavefunction must be computed and the computational effort is about the same as for one MRCI iteration. The RS2 and RS3 programs use the same configuration spaces as the MRCI, i.e., only the doubly external configurations are internally contracted.

A new version of the program has been implemented in which also subspaces of the singly external and internal configuration spaces are internally contracted (see reference given above). This program, which is called using the keyword RS2C, is more efficient than RS2, in particular for large molecules with many closed-shell (inactive) orbitals. It is recommended to use this program for normal applications of second-order multireference perturbation theory (CASPT2, RASPT2). Note that it gives slightly different results than RS2 due to the different contraction scheme. It should also be noted that neither RS2 or RS2C are identical with the CASPT2 of Roos et al. [J. Chem. Phys. **96**, 1218 (1992)], since certain configuration subspaces are left uncontracted. However, the differences are normally very small. The last point that should be mentioned is that the calculation of CASPT2/RASPT2 density matrices (and therefore molecular properties) is presently possible only with the RS2 command and *not* with RS2C.

The results of multireference perturbation theory may be sensitive to the choice of the zeroth-order Hamiltonian. This dependence is more pronounced in second-order than in third-order. Several options are available, which will be described in the following sections. It may also happen that $(\hat{H}^{(0)} - E^{(0)})$ in the basis of the configuration state functions becomes (nearly) singular. This is known as "intruder state problem" and can cause convergence problems or lead to a blow-up of the wavefunction. Often, such problems can be eliminated by including

more orbitals into the reference wavefunction, but of course this leads to an increase of the CPU time. The use of modified Fock operators (see below) or level shifts, as proposed by Roos and Andersson [Chem. Phys. Lett. **245**, 215 (1995)] may also be helpful. Presently, only "real" level shifts have been implemented.

With no further input cards, the wavefunction definition (core, closed, and active orbital spaces, symmetry) corresponds to the one used in the most recently done SCF or MCSCF calculation. By default, a CASSCF reference space is generated. Other choices can be made using the OCC, CORE, CLOSED, WF, SELECT, CON, and RESTRICT cards, as described for the CI program. The orbitals are taken from the corresponding SCF or MCSCF calculation unless an ORBITAL directive is given.

For a CASPT2 calculation, the zeroth-order hamiltonian can be brought to a block-diagonal form when (pseudo)canonical orbitals are used. This leads to fastest convergence. It is therefore recommended that in the preceding MULTI calculation the orbitals are saved using the CANONICAL directive (note that the default is NATORB).

Most options for MRCI calculations (like STATE, REFSTATE etc.) apply also for RS2(C) and RS3 and are not described here again. Some additional options which specific for CASPT2/3 and are described below.

19.2 Coupling MRCI and MRPT2

For particularly difficult cases with strong intruder problems, or in which second-order perturbation theory fails to predict reliable results, a new method that couples MRCI and CASPT2 has been developed. This variant is invoked using the CIPT2 directive:

```
CIPT2
```

In this case all excitations solely from active orbitals are treated by MRCI, while the remaining excitations involving inactive (closed-shell) orbitals are treated by second-order perturbation theory. Both methods are coupled by minimizing an appropriate energy functional. Of course, this method is much more expensive than MRPT2. The cost is comparable to the cost for an MRCI without correlating the inactive orbitals.

19.3 Excited state calculations

There are two possibilities to perform excited state calculations:

1) One can calculate each state separately. This is done using the card

```
STATE,1,root
```

where *root* is the desired root (i.e., 2 for the first excited state). In this case the Fock operator used in the zeroth-order hamiltonian is computed using the density for the given state.

2) Alternatively, two or more states can be computed simultaneously, using

```
STATE, n [root1, root2, ..., rootn]
```

where *n* is the number of states to be computed. The default is to compute the lowest *n* roots. Optionally, this default can be modified by specifying the desired roots *rooti* as shown. One should note that this *does not* correspond to the multi-state CASPT2 presented in [Chem. Phys. Lett. **288**, 299 (1998)].

In the case that several states are computed simultaneously, the fock operator employed in the zeroth-order Hamiltonian is computed from a state-averaged density matrix, and the zeroth-order hamiltonians for all states are constructed from the same fock operator. By default, equal weights for all states are used. This default can be modified using the WEIGHT directive

WEIGHT,*w1*, *w2*, ..., *wn*.

If a REFSTATE card is given (see section 18.3.5), the state-averaged fock operator is made for all reference states, and the WEIGHT card refers to the corresponding states.

19.4 Modified Fock-operators in the zeroth-order Hamiltonian.

The g_1 , g_2 , and g_3 operators proposed by Andersson [Theor. Chim. Acta **91**, 31 (1995)] as well as a further g_4 operator may be used. g_4 makes CASPT2 calculations size extensive for cases in which a molecule dissociates to high-spin open-shell (RHF) atoms.

The index n of the operator to be used is specified on the RS2, RS2C, or RS3 card:

RS2,*n*

RS2C,*n*

RS3,*n*

where n can take the values 1 to 4. Instead of the value n one can also specify G1, G2 etc.

19.5 Level shifts

Level shifts are often useful to avoid intruder state problems in excited state calculations. MOLPRO allows the use of shifts as described by Roos and Andersson, [Chem. Phys. Lett. **245**, 215 (1995)]. The shift can be specified on the RS2 or RS2C card

RS2 [*Gn*] [,SHIFT=*shift*]

RS2C [*Gn*] [,SHIFT=*shift*]

Typical choices for the shift is are 0.1 – 0.3. Only two figures after the decimal point are considered. The shift affects the results, the printed energies as well as the ENERGY variable include the energy correction for the shift as proposed by Roos and Andersson. At convergence, also the uncorrected energies are printed for comparison.

19.6 Integral direct calculations

RS2, RS2C, and RS3 calculations with very large basis sets can be performed in integral-direct mode. The calculation will be direct if a global DIRECT or GDIRECT card appears earlier in the input. Alternatively, (mainly for testing) DIRECT can be specified as an option on the RS*n*[C] card:

RS2 [*Gn*] [,SHIFT=*shift*] [,DIRECT]

RS2C [*Gn*] [,SHIFT=*shift*] [,DIRECT]

19.7 Options for CASPT2 and CASPT3

Other options can be set using the OPTION command. These options are mainly used for testing purposes and should be used with care. It should be noted that the only option that can be modified in the RS2C program is IFDIA: all others only work with RS2/RS3.

OPTION,*code1=value,code2=value,...*

Of relevance for the CASPT2/3 program are the following options:

IPOCS=0	(Default). Calculation uses uncontracted singles with RS2.
IPOCS=1	Non-interacting singles are projected out during update. This is an approximate procedure which should be used with care.
IPOCS=2	The singles are fully internally contracted in RS2. This is achieved via a projection operator during the coefficient update and may be inefficient.
IPOCS=3	Only singles with one or two holes in the closed-shells are internally contracted in RS2 using a projection operator.
IPROCI=0	(Default). Calculation uses uncontracted internals with RS2.
IPROCI=1	Internals with two holes in the inactive space are internally contracted in RS2 using a projection operator.
IPOCS=3 , IPROCI=1	This combination of options reproduces with RS2 the RS2C result using projection operators. This requires lot of memory and disk space and it is feasible only for small molecules.
IFDIA=0	(Default). All off-diagonal elements of the effective Fock matrix are included.
IFDIA=1	The internal-external block of the Fock-matrix is neglected. This eliminates the single-pair coupling.
IFDIA=2	All off-diagonal elements of the Fock matrix are neglected. This corresponds to CASPT2D of Andersson et al. Note: in this case the result is not invariant to rotations among active orbitals!
IHINT=0	(Default). Only one-electron integrals are used in the zeroth-order Hamiltonian for all interactions.
IHINT=1	The all-internal two-electron integrals are used in the zeroth-order Hamiltonian for the internal-internal and single-single interactions.
IHINT=2	The all-internal two-electron integrals in the zeroth-order Hamiltonian are used for the internal-internal, single-single, and pair-pair interactions. Using IHINT=2 and IDFIA=1 corresponds to Dyal's CAS/A method for the case that CASSCF references with no closed-shells (inactive orbitals) are used. Note that this requires more CPU time than a standard CASPT2 calculation. Moreover, convergence of the CAS/A method is often slow (denominator shifts specified on a SHIFT card may be helpful in such cases). In general, we do not recommend the use of IHINT with nonzero values.
NOREF=1	(Default). Interactions between reference configurations and singles are omitted.
NOREF=0	Interactions between reference configurations and singles are included. This causes a relaxation of the reference coefficients but may lead to intruder-state problems.

IMP3=2

After CASPT2 do variational CI using all internal configurations and the first-order wavefunctions of all states as a basis. In this case the second-order energy will correspond to the variational energy, and the third-order energy approximately to a Davidson-corrected energy. This is useful in excited state calculations with near-degeneracy situations.

20 MØLLER PLESSET PERTURBATION THEORY

Closed-shell Møller-Plesset perturbation theory up to full fourth order [MP4(SDTQ)] is part of the coupled-cluster program.

The commands MP2, MP3, MP4 perform the MP calculations up to the specified order (lower orders are included).

MP4;NOTRIPL; performs MP4(SDQ) calculations.

Normally, no further input is needed if the MPn card directly follows the corresponding HF-SCF. Otherwise, occupancies and orbitals can be specified as in the CI program. The resulting energies are stored in variables as explained in section 6.7.

20.1 Expectation values for MP2

One-electron properties can be computed as analytical energy derivatives for MP2. This calculation is much more expensive than a simple MP2, and therefore only done if an EXPEC card follows the MP2 card (the GEXPEC directive has no effect in this case). The syntax of the EXPEC card is explained in section 4.13. For an example, see section 21.6.1. The density matrix can be saved using

DM,record.ifil];

See also sections 21.7 and 21.8.

20.2 Density-fitting MP2 (DF-MP2, RI-MP2)

Density-fitting MP2 (RI-MP2) can be performed with standard density or and Poisson fitting basis sets. The present implementation works only without symmetry. The input is as follows:

DF-MP2;

Optionally, a card DFIT can follow on which the following options can be specified:

BASIS_MP2= <i>string</i> :	Fitting basis sets, e.g., JKFIT (default) for standard density fitting or DENSITY:POISSON for mixed density/Poisson fitting. These basis sets must have been defined in a previous BASIS block.
THRAO= <i>value</i> :	Screening threshold for 3-index integrals in the AO basis
THRMO= <i>value</i> :	Screening threshold for 3-index integrals in the MO basis
THROV= <i>value</i> :	Screening threshold for 2-index integrals of fitting basis.
THRPROD= <i>value</i> :	Screening product threshold for first half transformation.
SPARSE= <i>value</i> :	If Non-zero, use sparse algorithm in second-half transformation (default).

RI-MP2 is an alias for the command DF-MP2. At present, expectation values and gradients cannot be computed with DF-MP2.

21 THE CLOSED SHELL CCSD PROGRAM

Bibliography:

C. Hampel, K. Peterson, and H.-J. Werner, Chem. Phys. Lett. 190, 1 (1992)

All publications resulting from use of this program must acknowledge the above.

The CCSD program is called by the CISD, CCSD, BCCD, or QCI directives. CID or CCD can be done as special cases using the NOSINGL directive. The code also allows to calculate Brueckner orbitals (QCI and CCSD are identical in this case). Normally, no further input is needed if the CCSD card follows the corresponding HF-SCF. Optional ORBITAL, OCC, CLOSED, CORE, SAVE, START, PRINT options work as described for the MRCI program in section 18. The only special input directives for this code are BRUECKNER and DIIS, as described below.

The convergence thresholds can be modified using

```
THRESH,ENERGY=thrden,COEFF=thrvar
```

Convergence is reached if the energy change is smaller than *thrden* (default 1.d-6) and the square sum of the amplitude changes is smaller than *thrvar* (default (1.d-10)). The THRESH card must follow the command for the method (e.g., CCSD) and then overwrites the corresponding global options (see GTHRESH, sec. 4.11).

The computed energies are stored in variables as explained in section 6.7. As well as the energy, the T_1 diagnostic (T. J. Lee and P. R. Taylor, Int. J. Quant. Chem. S23 (1989) 199) is printed and stored in the variable T1DIAG for later analysis.

21.1 Coupled-cluster, CCSD

The command CCSD performs a closed-shell coupled-cluster calculation. Using the CCSD(T) command, the perturbative contributions of connected triple excitations are also computed. For further information on triples corrections see under RCCSD.

21.2 Quadratic configuration interaction, QCI

QCI or QCISD performs quadratic configuration interaction, QCISD. Using the QCI(T) or QCISD(T) commands, the contributions of connected triples are also computed by perturbation theory. Normally, no further input is needed if the QCI card follows the corresponding HF-SCF. Otherwise, occupancies and orbitals can be specified as in the CI program. For modifying DIIS directives, see section 21.5

21.3 Brueckner coupled-cluster calculations, BCCD

```
BCCD,[SAVE=record],[PRINT],[TYPE=type]
```

BCCD performs a Brueckner coupled-cluster calculation and computes Brueckner orbitals. With these orbitals, the amplitudes of the singles vanish at convergence. Using the BCCD(T) command, the contributions of connected triples are also computed by perturbation theory. Normally, no further input is needed if the BCCD card follows the corresponding HF-SCF. Otherwise, occupancies and orbitals can be specified as in the CI program. BRUECKNER parameters can be modified using the BRUECKNER directive.

The Brueckner orbitals and approximate density matrix can be saved on a MOLPRO dump record using the SAVE option. The orbitals are printed if the PRINT option is given. TYPE can be used to specify the type of the approximate density to be computed:

TYPE=REF	Compute and store density of reference determinant only (default). This corresponds to the BOX (Brueckner orbital expectation value) method of Chem. Phys. Lett. 315 , 248 (1999).
TYPE=TOT	Compute and store density with contribution of pair amplitudes (linear terms). Normally, this does not seem to lead to an improvement.
TYPE=ALL	Compute and store both densities

Note: The expectation variables are stored in variables as usual. In the case that both densities are made, the variables contain two values, the first corresponding to REF and the second to TOT (e.g., DMZ(1) and DMZ(2)). If TYPE=REF or TYPE=TOT is given, only the corresponding values are stored.

21.3.1 The BRUECKNER directive

BRUECKNER,*orbbrk,ibrstr,ibrueck,brsfak*;

This directive allows the modification of options for Brueckner calculations. Normally, none of the options has to be specified, and the BCCD command can be used to perform a Brueckner CCD calculation.

<i>orbbrk</i> :	if nonzero, the Brueckner orbitals are saved on this record.
<i>ibrstr</i> :	First iteration in which orbitals are modified (default=3).
<i>ibrueck</i> :	Iteration increment between orbital updates (default=1).
<i>brsfak</i> :	Scaling factor for singles in orbital updates (default=1).

21.4 Singles-doubles configuration interaction, CISD

Performs closed-shell configuration interaction, CISD. The same results as with the CI program are obtained, but this code is somewhat faster. Normally, no further input is needed. For specifying DIIS directives, see section 21.5

21.5 The DIIS directive

DIIS,*itedis,incdis,maxdis,itydis*;

This directive allows to modify the DIIS parameters for CCSD, QCISD, or BCCD calculations.

<i>itedis</i> :	First iteration in which DIIS extrapolation may be performed (default=4).
<i>incdis</i> :	Increment between DIIS iterations (default=1).
<i>maxdis</i> :	Maximum number of expansion vectors to be used (default=6).
<i>itydis</i> :	DIIS extrapolation type. <i>itedis</i> =1 (default): residual is minimized. <i>itedis</i> =2: ΔT is minimized.

In addition, there is a threshold *THRDIS* which may be modified with the THRESH directive. DIIS extrapolation is only done if the variance is smaller than *THRDIS*.

21.6 Examples

21.6.1 Single-reference correlation treatments for H₂O

```

***,h2o test
memory,1,m                !allocate 1 MW dynamic memory
geometry={o;h1,o,r;h2,o,r,h1,theta} !Z-matrix geometry input
basis=vtz                 !cc-pVTZ basis set
r=1 ang                   !bond length
theta=104                 !bond angle
hf                         !do scf calculation

text,examples for single-reference correlation treatments

ci                         !CISD using MRCI code
ci;cepa,1                 !cepa-1 using MRCI code                examples/
mp2                        !Second-order Moeller-Plesset          h2o_ccsd.com
mp3                        !Second and third-order MP
mp4                        !Second, third, and fourth-order MP4(SDTQ)
mp4;notripl               !MP4(SDQ)
cisd                       !CISD using special closed-shell code
ccsd(t)                   !coupled-cluster CCSD(T)
qci(t)                    !quadratic configuration interaction QCISD(T)
bccd(t)                   !Brueckner CCD(T) calculation
---
```

21.6.2 Single-reference correlation treatments for N₂F₂

```

***,N2F2 CIS GEOMETRY (C2h)
rnn=1.223,ang             !define N-N distance
rnf=1.398,ang            !define N-F distance
alpha=114.5;             !define FNN angle
geometry={N1
  N2,N1,rnn
  F1,N1,rnf,N2,alpha
  F2,N2,rnf,N1,alpha,F1,180}
basis=vtz                 !cc-pVTZ basis set                examples/
method=[hf,cisd,ccsd(t),qcisd(t),bccd(t)] !all methods to use          n2f2_ccsd.com
do i=1,#method           !loop over requested methods
$method(i)               !perform calculation for given methods
e(i)=energy              !save energy in variable e
enddo                    !end loop over methods
table,method,e           !print a table with results
title,Results for n2f2, basis=$basis !title of table
```

This calculation produces the following table:

Results for n2f2, basis=VTZ

METHOD	E	E-ESCF
CISD	-308.4634948	-0.78283137
BCCD(T)	-308.6251173	-0.94445391
CCSD(T)	-308.6257931	-0.94512967
QCISD(T)	-308.6274755	-0.94681207

21.7 Saving the density matrix

DM,*record.ifil*];

The effective first order density matrix is computed and stored in record *record* on file *ifil*. This currently works for closed-shell MP2 and QCISD. See also NATORB.

21.8 Natural orbitals

NATORB,[RECORD=]*record.ifil*,[PRINT=*nprint*],[CORE[=*natcor*]];

Calculate natural orbitals. This currently only works for closed-shell MP2 and QCISD. The number of printed external orbitals in any given symmetry is *nprint* (default 2). *nprint=-1* suppressed the printing. The natural orbitals and the density matrix are saved in a dump record *record* on file *ifil*. If *record.ifil* is specified on a DM card (see above), this record is used. If different records are specified on the DM and NATORB cards, an error will result. The record can also be given on the SAVE card. Note that the effective density matrix of non-variational methods like MP2 or QCISD does not strictly behave as a density matrix. For instance, it has non-zero matrix elements between core and valence orbitals, and therefore core orbitals are affected by the natural orbital transformation. Also, occupation numbers of core orbitals can be larger than 2.0. If CORE is given (*natcor=1*), the core orbitals are frozen by excluding them from the natural orbital transformation.

21.9 Excited states using linear response (CCSD-LR, EOM-CCSD)

Excitation energies can be computed using linear response (LR) theory (also called equation of motion (EOM) approach). Accurate results can only be expected for singly excited states. The states to be computed are specified on an EOM input card, which is a subcommand of CCSD. The following input forms are possible

EOM, *state1*, *state2*, *state3*, ...

Computes the given states. Each state is specified in the form *number.sym*, e.g., 5.3 means the fifth state in symmetry 3. Note that state 1.1 corresponds to the ground state CCSD wavefunction and is ignored if given.

EOM, $-n1.sym1$, $-n2.sym2$, ...

Computes the first *n1* states in symmetry *sym1*, *n2* in *sym2* etc.

EOM, $n1.sym1$, $-n2.sym1$, ...

Computes states *n1* through *n2* in symmetry *sym1*.

The different forms can be combined, e.g.,

EOM, -3.1 , 2.2, 2.3, -5.3

computes states 1-3 in symmetry 1, the second excited state in symmetry 2, and the second through fifth excited states in symmetry 3. Note that state 1.1 is the ground-state CCSD wavefunction.

21.9.1 Parameters for EOM-CCSD (EOMPAR)

Normally, no further input is needed. However, some defaults can be changed using the EOMPAR directive:

EOMPAR, key1=value1, key2=value2,...

where the following keywords key are possible:

MAXDAV= <i>nv</i>	Maximum value of expansion vectors per state in Davidson procedure (default 10).
INISINGL= <i>ns</i>	Number of singly excited configurations to be included in initial Hamiltonian (default 20; the configurations are ordered according to their energy). Sometimes INISINGL should be put to zero in order to catch states dominated by double excitations.
INIDOUBL= <i>nd</i>	Number of doubly excited configurations to be included in initial Hamiltonian (default 10).
INIMAX= <i>nmax</i>	Maximum number of excited configurations to be included in initial Hamiltonian. By default, $nmax = ns + nd$.
MAXITER= <i>itmax</i>	Maximum number of iterations in EOM-CCSD (default 20).
MAXEXTRA= <i>maxex</i>	Maximum number of extra configurations allowed to be included in initial Hamiltonian (default 0). In the case of near degeneracy it is better to include a few extra configurations to avoid a slow convergence.
DIFOCK= <i>difo</i>	If set to 0, the program uses an approximate diagonal of \bar{H} for looking for the initial configurations (corresponding to the $nmax + maxex$ lowest diagonal elements) and for the vector update in the Davidson procedure. If set to 1, the diagonal of the Fock matrix is used instead. If set to 2, the update procedures from CCSD program are used. DIFOCK=1 and 2 should give exactly the same results. Default for non-local EOM-CCSD (even based on local CCSD) is 0. For the local EOM-CCSD DIFOCK is always set to 2.
EOMLOCAL= <i>eoml</i>	If set to 0, non-local calculation (default). If set to 1, it simulates the local EOM-CCSD using the standard EOM program. EOMLOCAL=1 is a test option and shouldn't be used for the time being.

INIMAX is used only if INISINGL and INIDOUBL are both zero.

All keywords can be abbreviated by at least four characters.

21.9.2 Print options for EOM-CCSD (EOMPRINT)

The following print options are for testing purposes and for looking for the convergence problems.

EOMPRINT, key1=value1, key2=value2,...

where the following keywords key are possible:

DAVIDSON= <i>ipr</i>	Information about Davidson procedure: <i>ipr=1</i> print results of each "small diagonalization"
----------------------	---

ipr=2 also print warning information about complex eigenvalues
ipr=3 also print hamiltonian and overlap matrix in trial space.

DIAGONAL Information about configurations:
ipr=1 print the lowest approximate diagonal elements of the transformed hamiltonian
ipr=2 print orbital labels of important configurations
ipr=3 print all approximate diagonal elements
ipr=4 also print the long form of above.

PSPACE Print information about the initial approximate hamiltonian:
ipr=2 print the approximate hamiltonian used to find the first approximation.

HEFF Print information about effective Hamiltonian:
ipr=2 print columns of effective hamiltonian and overlap matrix in each iteration

RESIDUUM Print information about residuum vectors:
ipr=-1 no print in iteration
ipr=0 print energy values + residuum norm (squared) for each iteration (default)
ipr=1 also print warning about complex eigenvalue, and a warning when no new vectors is added to the trial space due to the too small norm of the residuum vector.
ipr=2 also print how many vectors are left

22 OPEN-SHELL COUPLED CLUSTER THEORIES

Spin unrestricted (RHF-UCCSD) and partially spin restricted (RHF-RCCSD) open-shell coupled cluster theories as described in J. Chem. Phys. **99** (1993) 5219 (see also erratum, J. Chem. Phys., submitted for publication) are available in MOLPRO. In both cases a high-spin RHF reference wavefunction is used. In the description that follows, the acronyms RCCSD and UCCSD are used, but the theories should normally be referred to as RHF-RCCSD, RHF-UCCSD, in order to distinguish them from alternative ansätze based on spin-unrestricted orbitals. The program will accept either the full or abbreviated acronyms as input commands.

In the RCCSD theory certain restrictions among the amplitudes are introduced, such that the linear part of the wavefunction becomes a spin eigenfunction (this is not the case in the UCCSD method, even if an RHF reference function is used). At present, the implementation of RCCSD is only preliminary, and no CPU time is saved by as compared to UCCSD. However, improved algorithms, as described in the above publication, are currently being implemented, and will be available in the near future.

The input is exactly the same as for closed-shell CCSD, except that RCCSD or UCCSD are used as keywords. By default, the open-shell orbitals are the same as used in the RHF reference function, but this can be modified using OCC, CLOSED, and WF cards.

Perturbative triples corrections are computed as follows:

- RCCSD(T), UCCSD(T) triples corrections are computed as defined by J. D. Watts, J. Gauss and R. J. Bartlett, J. Chem. Phys. **98** 8718 (1993).
- RCCSD[T], UCCSD[T] corrections are computed without contributions of single excitations (sometimes called CCSD+T(CCSD)).
- RCCSD-T, UCCSD-T triples corrections are computed as defined by M. J. O. Deegan and P. J. Knowles, Chem. Phys. Letters 227 (1994) 321.

In fact, all three contributions are always computed and printed. The following variables are used to store the results (here CCSD stands for either UCCSD or RCCSD):

ENERGY	total energy for method specified in the input.
ENERGC	total CCSD energy without triples.
ENERGT(1)	total CCSD(T) energy.
ENERGT(2)	total CCSD[T] energy.
ENERGT(3)	total CCSD-T energy.

We note that the present implementation performs the (partial) integral transformation using an algorithm which is not optimal in its memory use, and therefore needs quite a large memory if extensive basis sets are used. Future improvements will remove this bottleneck.

23 LOCAL CORRELATION TREATMENTS

23.1 Introduction

The local correlation program of MOLPRO can currently perform closed-shell MP2(D), MP3(D), MP4(SDQ), CISD, QCISD, and CCSD calculations. So far, only MP2 is publicly available. Linear scaling versions of the other methods are still experimental and will be released later. Perturbative energy corrections for triple excitations are under development and will also be available in the near future. Note that you have to install the ‘local’ module in order to use local correlation methods. If you want linear scaling, you also need the ‘direct’ module.

References:

General, local CCSD:

[1] C. Hampel and H.-J. Werner, *Local Treatment of electron correlation in coupled cluster (CCSD) theory*, J. Chem. Phys. **104**, 6286 (1996).

Multipole treatment of distant pairs:

[2] G. Hetzer, P. Pulay, H.-J. Werner, *Multipole approximation of distant pair energies in local MP2 calculations*, Chem. Phys. Lett. **290**, 143 (1998).

Linear scaling local MP2:

[3] M. Schütz, G. Hetzer and H.-J. Werner, *Low-order scaling local electron correlation methods. I. Linear scaling local MP2*, J. Chem. Phys. **111**, 5691 (1999).

LMP2 Gradients and geometry optimization:

[4] A. El Azhary, G. Rauhut, P. Pulay and H.-J. Werner, *Analytical energy gradients for local second-order Møller-Plesset perturbation theory*, J. Chem. Phys. **108**, 5185 (1998).

LMP2 vibrational frequencies:

[5] G. Rauhut, A. El Azhary, F. Eckert, U. Schumann and H.-J. Werner, *Impact of Local Approximations on MP2 Vibrational Frequencies*, Spectrochimica Acta **55**, 651 (1999).

Intermolecular interactions and the BSSE problem:

[6] M. Schütz, G. Rauhut and H.-J. Werner, *Local Treatment of Electron Correlation in Molecular Clusters: Structures and Stabilities of $(H_2O)_n$, $n = 2 - 4$* , J. Phys. Chem. **102**, 5997 (1998).

[7] N. Runeberg, M. Schütz and H.-J. Werner, *The aurophilic attraction as interpreted by local correlation methods*, J. Chem. Phys. **110**, 7210 (1999).

23.2 Getting started

The local correlation treatment is switched on by preceding the command name by an L, i.e., by using the LMP2, LMP3, LMP4, LQCI [[SD] (T)], LCCSD [(t)], or LCISD directives. Further options can be given on the same input card or on subsequent LOCAL or MULTP cards. Alternatively, one can also give the LOCAL or the MULTP directive after the normal MP2, MP3, MP4, QCI [[SD] (T)], CCSD [(T)], or CISD directives.

Thus, the two input cards

```
METHOD;
LOCAL,[key1=value],[key2=value2], ...
```

are equivalent to

```
LMETHOD,[key1=value],[key2=value2], ...
```

where METHOD is one of MP2, MP3, MP4, QCI [[SD] (T)], CCSD [(T)], or CISD.

Similarly,

METHOD;

MULTP,[key1=value],[key2=value2], ...

is equivalent to

LMETHOD,MULTP,[key1=value],[key2=value2], ...

(The full set of options is described in section 23.6, and summarized in Table 9.) The LOCAL and MULTP directives only differ in the defaults that they assume for the input keys.

The LOCAL directive requests a traditional local correlation calculation, where all pairs (of occupied orbitals) that are correlated by MP2 are treated equal, regardless of their distance. The MULTP directive turns on additional approximations that depend on the distance between the orbitals: The *distant pairs* are treated by a multipole approximation as described in Ref. [2], and *very distant pairs* are neglected. This is a prerequisite to obtain linear scaling for large molecules. Using LOCAL (without choosing appropriate settings manually) will result in $O(N^3)$ scaling. Be sure to read the applicable parts of the next section before starting your own calculations.

23.3 Doing it right

23.3.1 Always...

Turn off symmetry! Otherwise, you won't get appropriately localized orbitals (local orbitals will tend to be symmetry equivalent instead of symmetry adapted). Symmetry is in principle OK only if all atoms are symmetry unique. This allows the treatment of planar molecules in C_s symmetry when using the LOCAL directive. But note that the multipole program does not support symmetry at all, so choose always C_1 symmetry with the MULTP directive.

To turn off symmetry, specify nosym as the first line of your geometry input, e.g.

```
geometry={
  nosym
  O1
  H1,O1,roh
  H2,O1,roh,h1,hoh
}
```

Use NOORIENT! Under certain circumstances it may happen that the domains and correlation energies are not rotationally invariant. We therefore recommend to use the NOORIENT option in the geometry input, to avoid unintended rotations of the molecule when the geometry changes. This is particularly important for geometry optimizations and for calculations of interaction energies (see section 23.4.1).

Check your orbital domains! Local correlation methods are less 'black box' than the canonical ones. It is therefore recommended always to check the orbital domains, which are printed in the beginning of each local calculation. For checking, the option DOMONLY=1 can be used to stop the calculation after the domain generation. The orbital domains consist of all basis functions for a subset of atoms. These atoms are selected so that the domain spans the corresponding localized orbital with a preset accuracy (alterable with key DOMSEL). A typical domain output, here for water, looks like this:

Orbital domains

Orb.	Atom	Charge	Crit.
2.1	1 O1	1.17	0.00
	3 H2	0.84	1.00
3.1	1 O1	2.02	1.00
4.1	1 O1	1.96	1.00
5.1	1 O1	1.17	0.00
	2 H1	0.84	1.00

This tells you that the domains for orbitals 2.1 and 5.1 comprise the basis functions of the oxygen atom and one hydrogen atom, while the domains for orbitals 3.1 and 4.1 consist of the basis function on oxygen only. The latter ones correspond to the oxygen lone pairs, the former to the two OH bonds, and so this is exactly what one would expect.

Improper domains could result from improperly localized orbitals or forgotten NOSYM directive. This does not only negatively affect performance and memory requirements, but can also lead to bogus results. Poor localization is sometimes an intrinsic problem, in particular for strongly conjugated systems. In rare cases it might also happen that the localization procedure does not converge.

The default for the selection criterion DOMSEL is 0.98. This works usually well for small basis sets like cc-pVDZ. For larger basis sets like cc-pVTZ we recommend to use a slightly larger value of 0.985 to ensure that enough atoms are included in each domain.

There are some other options which affect the domain selection:

CHGMIN= <i>value</i>	determines the minimum allowed Mulliken charge for an atom (except H) in a domain, i.e., atoms with a smaller (absolute) charge are not included, even if the DOMSEL criterion is not fulfilled (default 0.01).
CHGMINH= <i>value</i>	as CHGMIN, but used for H-atoms (default 0.03).
CHGMAX= <i>value</i>	If Mulliken charge is larger than this value, the atom is included independent of any ranking.
MAXBP= <i>maxbp</i>	If <i>maxbp</i> =1, the atoms are ranked according to their contribution to the Boughton-Pulay overlap. If <i>maxbp</i> =0 (default), the atoms are ranked according to atomic charges. In both cases atoms with charges greater than CHGMAX are always included, and atoms with the same charges are added as groups.
MULLIKEN= <i>option</i>	Determines method to determine atomic charges. <i>option</i> =0 (default): squares of diagonal elements of $\mathbf{S}^{-\frac{1}{2}}\mathbf{C}$ are used. <i>option</i> =1: Mulliken gross charges. It appears that the first choice works better with diffuse basis sets.
MERGEDOM= <i>number</i>	If this option is given, all orbital domains containing <i>number</i> or more atoms in common are merged (<i>number</i> =1 is treated as <i>number</i> =2, default 0). This is particularly useful for geometry optimizations of conjugated or aromatic systems like, e.g., benzene. In the latter case, MERGEDOM=1 causes the generation of full π -domains, i.e., the domains for all three π -orbitals comprise all carbon basis functions. Note that the merged domains are generated after the above print of orbital domains, and information about merged domains is printed separately. See section 23.4.2 for further discussion of geometry optimizations.

These options can be disabled by setting their values to zero.

23.3.2 Linear scaling MP2

Linear scaling of the CPU time as well as memory and disk requirements with molecular size (for a fixed basis set) can be achieved for very extended systems if the calculation is performed in integral-direct manner (i.e., the 2-electron integrals are never stored on disk) [3]. This requires the DIRECT module. In order to achieve low-order scaling the MULTP directive must be used. Normally, the program uses appropriate defaults and no further options must be set. Thus, the typical input structure looks as follows:

```
memory,64,m      !specify enough memory if you try a very large calculation
file,2,name.wfu  !save orbitals and other info for later restart
gdirect         !enable integral-direct mode
basis={...}     !basis specification
geometry={nosym;noorient;...} !geometry specification
hf              !scf calculation (this often takes most of the time, since at
               !present this is not linearly scaling)
locali,pipek    !Orbital localization (this card is optional). If not given
               !Pipek-Mezey localization is done automatically within the LMP2 program
mp2;multp      !Local MP2 with multipole approximation for distant pairs.
```

Notice that for small systems a local MP2 with these default options might take more time than a conventional MP2. There are two reasons for this: firstly, the default options in the LMP2 case are chosen so that the memory and disk requirements are minimized. This requires the evaluation of each unique integral twice in the transformation. Secondly, the LMP2 equations have to be solved iteratively, which takes some additional time. It is possible to speed up local MP2 calculations for small or medium size molecules using the following option on the GDIRECT card:

```
GDIRECT, PAGE=1
```

In this case each unique integral is computed only once, but intermediate storage of the half transformed integrals on disk is required. The disk requirements scale cubically in this case, and therefore this option cannot be used for very large cases.

The CPU-time depends sensitively on the prescreening thresholds used in the transformation. Their choice, in turn, depends on the required accuracy. The default thresholds (1.d-7) normally ensure that absolute numerical errors in the energy are around 1 μ H. For many applications, lower accuracy is sufficient, and a change of the default values can be achieved by requesting a less accurate energy:

```
GTHRESH, ENERGY=1 .d-5
```

This card must be given before the MP2 directive and will then increase the prescreening thresholds to 1.d-6 (note that it might also influence the accuracy of the SCF if given before the HF card!). On the other hand, requesting a more accurate energy will tighten the thresholds. Tighter thresholds will also be chosen automatically if the AO overlap matrix has very small eigenvalues, which can happen for large and diffuse basis sets.

The prescreening thresholds for LMP2 can also be changed using the specific options THRAO_LMP2, THRQ1_LMP2, and THRQ2_LMP2 on the DIRECT card (see section 8), but this is only recommended for experienced users.

23.4 Density-fitting LMP2 (DF-LMP2, RI-LMP2)

Density-fitting LMP2 can be performed with standard density or and Poisson fitting basis sets. The present implementation works only without symmetry. The input is as follows:

DF-LMP2,[*lmp2 options*]

Optionally, a card DFIT can follow on which the following options can be specified (appropriate default values are available):

BASIS_MP2= <i>string</i>	Fitting basis sets, e.g., JKFIT (default) for standard density fitting or DENSITY:POISSON for mixed density/Poisson fitting. These basis sets must have been defined in a previous BASIS block.
THROV= <i>value</i> :	Screening threshold for 2-index integrals of fitting basis.
THRAO= <i>value</i> :	Screening threshold for 3-index integrals in the AO basis.
THRMO= <i>value</i> :	Screening threshold for half-transformed 3-index integrals.
THRSW= <i>value</i> :	Threshold for Schwarz screening.
THRPROD= <i>value</i> :	Product threshold for screening in first half transformation.
SPARSE= <i>value</i> :	If <i>value</i> is on-zero, use sparse algorithm in second-half transformation (default).
LOCFIT= <i>value</i> :	If <i>value</i> =1 use united orbital fitting domains (default). If <i>value</i> =2 use united pair fit domains.
RDOMAUX= <i>value</i> :	Radius for extending the pair fitting domains. Should be nonzero if LOCFIT=2.
KSCEEN= <i>value</i> :	If <i>value</i> =1 use Schwarz screening (linear scaling algorithm).
MINBLK= <i>value</i> :	Minimum AO blocking size.
MAXBLK= <i>value</i> :	Maximum AO blocking size.
MAXFIT= <i>value</i> :	Maximum block size for fitting functions.
MAXBATCH= <i>value</i> :	Blocking size for first half transformation.

RI-MP2 is an alias for the command DF-MP2. At present, expectation values and gradients cannot be computed with DF-MP2.

23.4.1 Intermolecular interactions

For calculations of interaction potentials of weakly interacting systems, the domains of the sub-systems should be determined at a very large distance and saved using the SAVE=*record* option on the LOCAL or MULTP cards, or the SAVE command (see section 23.5.1). If the asymptotic energy is not needed it is sufficient to do this initial calculation using option DOMONLY=1). These domains should then be reused in the subsequent calculations at all other intermolecular distances by using the START=*record* option or the START command (see section 23.5.2). Only in this way the basis set superposition error is minimized and normally negligible (of course, this does not affect the BSSE for the SCF, and therefore the basis set should be sufficiently large to make the SCF BSSE negligible). Usually, diffuse basis functions are important for obtaining accurate intermolecular interactions. Unfortunately, these spoil the efficiency of prescreening and therefore make direct calculations much more expensive.

For examples and discussions of these aspects see Refs. [6,7]

23.4.2 Gradients and frequency calculations

Geometry optimizations [4,5] and numerical frequency calculations [5] can be performed using analytical energy gradients [4] for local MP2. LMP2 geometry optimizations are particularly attractive for weakly bound systems, since virtually BSSE free structures are obtained (see section 23.4.1 and Refs. [6,7]). It should be noted, however, that the current implementation is not particularly efficient, and nothing has been done so far to achieve low-order scaling for large systems. Analytical energy gradients are not yet available for the multipole approximation of distant pairs, and therefore MULTP cannot be used in geometry optimizations or frequency calculations.

A particular problem in LMP2 gradient calculations is the elimination of redundant basis functions in the domains (see Refs. [1,4]). If redundancies are present, gradient calculations require the elimination of individual basis functions (option DELBAS=1), which is less unique than the elimination of eigenvectors corresponding to small eigenvalues of the overlap matrix (option DELBAS=0). If the redundancies are exact, i.e., if the overlap matrix for a domain has zero eigenvalues, it is in principle irrelevant which function is deleted. In practice, however, the selection sometimes influences the numerical stability. On the other hand, if the overlap matrix has very small but nonzero eigenvalues, the computed energy slightly depends on which basis functions are eliminated. We tried very hard to make the selection algorithm as robust as possible, but pitfalls in certain cases cannot be fully excluded. Problems with the redundancy elimination normally occur only for very small molecules (2 or 3 atoms) with very small basis sets. If difficulties are encountered, it is recommended to use PRINT , DOMAINS to obtain more detailed information about domains and redundant functions. The default behaviour can be changed using the DELBAS, DELEIG, DELSHL, TYPECHECK, and THRLOC options.

Whenever possible, the domains should be made *rotationally invariant*, which can be achieved by eliminating shells of basis functions (see DELSHL option). Sometimes it may be necessary to modify the threshold THRLOC to obtain the desired result. In order to avoid problems when rotational invariance is not strictly fulfilled, we recommend always to use of the NOORIENT option in the geometry input for optimizations or frequency calculations.

The SAVE and START options should always be used to keep the domains fixed during geometry optimizations, frequency calculations, or whenever smooth potential energy functions are required. In optimizations with very large geometry changes, it may be useful to determine new domains at the optimized geometry and repeat the geometry optimization with these domains.

Particular care must be taken in optimizations of highly symmetric aromatic systems, like, e.g., benzene. In D_{6h} symmetry, the localization of the π -orbitals is not unique, i.e., the localized orbitals can be rotated around the C_6 axis without changing the localization criterion. This redundancy is lost if the symmetry is slightly distorted, which can lead to large changes of the localized orbitals. If now the domains are kept fixed using the SAVE and START options, a large error in the energy might result. On the other hand, if the domains are not kept fixed, their size and quality might change during the optimization, again leading to spurious energy changes and divergence of the optimization. The best way to avoid this problem is to use the MERGEDOM=1 option (see section 23.3.1). If this option is given, the domains for the π orbitals will comprise the basis functions of all six carbon atoms, and the energy will be invariant with respect to unitary transformations among the three π orbitals. Note that this problem does not occur if the symmetry of the aromatic system is lowered by a substituent.

Finally, we note that the LMP2 gradients are quite sensitive to the accuracy of the SCF convergence (as is also the case for MP2). If very accurate structures are required, or if numerical frequencies are computed from the gradients, the default SCF accuracy might be insufficient. We recommend in such cases to add an ACCU , 14 directive (possibly even ACCU , 16) after the

HF card. Indicative of insufficient SCF accuracy are small positive energy changes near the end of the geometry optimization.

23.4.3 Basis sets

For numerical reasons, it is useful to eliminate projected core orbitals, since these may have a very small norm. By default, projected core orbitals are eliminated if their norm is smaller than 0.1 (this behaviour can be changed using the DELCOR and THRCOR options). For local calculations we recommend the use of generally contracted basis sets, e.g., the correlation consistent cc-pVnZ sets of Dunning and coworkers. For these basis sets the core basis functions are uniquely defined, and will always be eliminated if the defaults for DELCOR and THRCOR are used.

23.5 Further commands

23.5.1 Saving the wavefunction SAVE

The wavefunction can be saved for later restart using

`SAVE,record`

where *record* has the usual form, e.g., 4000.2 means record 4000 on file 2. If this command is given, the domain information as well as the amplitudes are saved (for MPn the amplitudes are not saved). If just the domain information should be stored, the SAVE option on the LOCAL card can also be used (cf. section 23.6).

23.5.2 Restarting a calculation START

Local CCSD or QCISD calculations can be restarted using

`START,record`

The record given must have been saved in a previous local calculation using the SAVE directive. If the START command is given, the domain information as well as the amplitudes of the previous calculation are used for restart. It is possible, for instance, to start a local CCSD calculation with the amplitudes previously saved for a local QCISD calculation (but of course it is not possible to use a record saved for a non-local CCSD or QCISD calculation). If it is intended only to use the domain information but not the amplitudes for a restart, the START option on the LOCAL card can be used (cf. section 23.6).

23.5.3 Defining orbital domains DOMAIN

Normally, the orbital domains are determined automatically using the procedure of Boughton and Pulay, J. Comput. Chem., **14**, 736 (1993) and J. Chem. Phys. **104**, 6286 (1996). The selection criterion can be modified by the DOMSEL key (see section 23.6). It is also possible to define the domains “by hand”, using the DOMAIN directive:

`DOMAIN,orbital,atom1,atom2...`

where *orbital* has the form *i orb. isym*, e.g., 3.1 for the third orbital in symmetry 1, and *atom_i* are the atomic labels as given in the Z-matrix geometry input, or, alternatively, the Z-matrix row numbers. All basis functions centred at the given atoms are included into the domain. For instance

`DOMAIN,3.1,C1,C2`

defines a domain for a bicentric bond between the carbon atoms C1 and C2. The DOMAIN cards must be given after any OCC, CLOSED, or CORE directives. Note that the order of the localized orbitals depends on the localization procedure, and could even change as function of geometry, and therefore manual DOMAIN input should be used with great care. The domains of all orbitals which are not explicitly defined using DOMAIN cards are determined automatically as usual.

23.5.4 Correlating subsets of electrons ATOMLIST

In large molecules, it may be sufficient to correlate only the electrons in the vicinity of an *active group*, and to treat the rest of the molecule only at the SCF level. The ATOMLIST directive allows the specification of a subset of atoms:

ATOMLIST,*atom1*, *atom2* ...

The program will then correlate only electrons in orbitals whose domains are exclusively covered by the given atoms. Electrons in a bonding orbital from one of the given atoms to one which is not part of the list are not correlated. This may significantly reduce the computation time, and, provided the active atoms are sensibly chosen, may give still sufficiently accurate results for the *active group*, e.g. bond lengths and bond angles.

23.5.5 Energy partitioning for molecular cluster calculations ENEPART

The local character of occupied and virtual orbitals in the local correlation treatment also offers the appealing possibility to decompose the intermolecular interaction energy of molecular clusters into individual contributions of different excitation classes. This allows to distinguish between intramolecular-, dispersive-, and ionic components of the correlation contribution to the interaction energy (cf. M. Schütz, G. Rauhut and H.J. Werner, *J. Phys. Chem.* **102**, 5197 (1998)). The energy partitioning algorithm is activated either by supplying the global ENEPART card:

```
ENEPART,[epart],[iepart]
```

The *epart* parameter determines the cutoff distance for (intramolecular) bond lengths (in a.u., default 3 a.u.) and is used to automatically determine the individual monomer subunits of the cluster. The *iepart* parameter enables the energy partitioning, if set to a value larger than zero (default 1). Additionally, if *iepart* is set to 2, a list of all intermolecular pair energies and their components is printed.

The output section produced by the energy partitioning algorithm will look similar to the following example:

```
energy partitioning enabled !
centre groups formed for cutoff [au] = 3.00
 1  :O1  H11 H12
 2  :O2  H21 H22
energy partitioning relative to centre groups:
intramolecular correlation:    -.43752663
exchange dispersion           :    .00000037
dispersion energy             :    -.00022425
ionic contributions           :    -.00007637
```

The centre groups correspond to the individual monomers determined for *epart*=3. In the present example, two water monomers were found. The correlation energy is partitioned into the four components shown above. The exchange dispersion, dispersion and ionic components reflect directly the related intermolecular components of the complex, while the intramolecular correlation contribution to the interaction energy has to be determined by a supermolecular calculation, i.e. by subtracting the (two) corresponding monomer correlation energies from the intramolecular correlation component of the complex given in the output.

23.5.6 Split Coulomb operator treatment of weak and strong pairs ATTENUATE

This method is still in development and has to be considered experimental. We document it here because we hate ‘undocumented features’, but please, do yourself a favour and don’t use it for the time being. If you get in trouble with it, we won’t help you.

The method relies on the partitioning of the Coulomb operator into a rapidly decaying short range part containing the singularity and a smooth long range part. The integrals over both parts

of the Coulomb operator are then treated separately. The short range integrals are obtained by transformation of the short range integrals in the AO basis, which is faster than the conventional transformation as more efficient screening is possible. The long range integrals are treated by a multipole expansion. In contrast to conventional multipole expansions, this expansion has an infinite radius of convergence. The method is available by replacing the LOCAL or MULTP card by the ATTENUATE card.

ATTENUATE,[key1=value],[key2=value2], ...

It does everything the MULTP card does (i. e., distant pairs are still treated by ordinary multipole expansion), plus it will enable the split Coulomb operator treatment of weak and strong pairs and select reasonable defaults. See section 23.6 for details. If you don't want distant pairs to be treated by ordinary multipole expansion, simply specify DISTPAIR=0 on the ATTENUATE card. Note that this method will only work in the context of integral-direct calculations.

23.6 Options

Various options can be specified using key/value pairs qualifying the LOCAL or MULTP command. For all options appropriate default values are set, and so these options must usually be modified only for special purposes. For convenience and historical reasons, alias names are available for various options, which correspond to the variable name used in the program. Table 9 summarizes the keys, aliases and default values. In the following, the parameters are described in more detail.

General Parameters:

LOCAL= <i>local</i>	<p>Determines which method is used:</p> <p>LOCAL=0: Conventional (non-local) calculation.</p> <p>LOCAL=1: Local method is simulated using canonical MOs. The local basis is used only at an intermediate stage to update the amplitudes in each iteration (only for testing).</p> <p>LOCAL=2: Calculation is done in local basis, but without using local blocking (i.e. full matrices are used). This is the most expensive method and only for testing.</p> <p>LOCAL=3: Fully local calculation. This is the fastest method for local calculations with no weak pairs.</p> <p>LOCAL=4: Fully local calculation (Default). This is the fastest method for large molecules with many weak pairs and requires minimum memory.</p>
SAVE= <i>record</i>	<p>Allows the domain information to be saved on <i>record=name.ifil</i> for later restart using START. This can be used to freeze the domains as function of geometry. Note that the domain information is automatically stored if a SAVE directive is given (see above), and in this case the record given on the SAVE card will overwrite any record given as SAVE option.</p>
START= <i>record</i>	<p>Retrieves domain information previously saved using SAVE. Note that the domain information is automatically restored if a START directive is given (see above), and in this case the record given on the START card will overwrite any record given as START option.</p>
PIPEK= <i>option</i>	<p>If this option is given and <i>option</i> > 0, the orbitals are localized using the Pipek-Mezey technique. If this option is not given or <i>option</i>=0 (default), the orbitals are localized unless localized orbitals are found</p>

in the orbital record (cf. ORBITAL and LOCALIZE directives). In the latter case, the most recent localized orbitals are used. Setting *option*=-1 switches the localization off. If *option*> 1 the localized orbitals are printed. Note: Boys localization can only be performed using the LOCALIZE command. The program will use the Boys orbitals if they are found in the orbital record and the PIPEK option is absent or *option*≤ 0.

- SAVORB=*record* Allows the localized and projected orbitals to be saved on *record=name.ifil* for later use (e.g. plotting). The two orbital sets are stored in the same dump record and can be restored at later stages using ORBITAL,*record*,[TYPE=]LOCAL or ORBITAL,*record*,[TYPE=]PROJECTED, respectively.
- DOMONLY=*value* If *value*> 0 only domains are made, but no energy is computed. This can be used to check and save the domains for later use.

Parameters for selection of weak and distant pairs:

- WEAKPAIR=*distance* If all atoms of orbital domain [i] are separated by at least *distance* [a.u.] from any atom of orbital domain [j], pair (ij) is treated by MP2. The default is *distance*=1, which means that all pairs for which [i] and [j] have no atom in common are treated as weak pairs. Setting *distance*=0 eliminates weak pairs, i.e. all pairs are fully included in the calculation. This option has no effect for local MP2 calculations.
- DISTPAIR=*distance* If all atoms of orbital domain [i] are separated by at least *distance* [a.u.] from any atom of orbital domain [j], pair (ij) is treated approximately by MP2, provided the multipole approximation is activated. Setting *distance*=0 eliminates distant pairs, i.e. no pairs are treated approximately. Default is 0 (MULTP card: 8).
- VERYDIST=*distance* If all atoms of orbital domain [i] are separated by at least *distance* [a.u.] from any atom of orbital domain [j], pair (ij) is neglected. Setting *distance*=0 (default) eliminates very distant pairs, i.e. no pairs are neglected. Reasonable values for *distance* would be 12–15 [a.u.] Default is 0 (MULTP card: 15).

Parameters to define domains:

- DOMSEL=*value* Threshold for selecting the atoms contributing to orbital domains using the method of Boughton and Pulay. The default is *value*=0.98. *value*=1.0 would include all atoms into each orbital domain. The criterion is somewhat basis dependent. The larger the basis, the fewer functions will be selected with a given threshold. The default value usually works well for double-zeta basis sets. For larger basis sets (e.g., cc-pVTZ) it is recommended to use *value*=0.985. In most cases, the domain selection is uncritical for saturated molecules. However, for delocalized systems it is recommended to check the printed orbital domains! In cases of doubt, compare the domains you get with a smaller basis (e.g., cc-pVDZ). See also the MAXANG option below.
- DELCOR=*nshell* Activates elimination of basis functions corresponding to core orbitals. If *nshell*=1, only 1s-functions are eliminated from projected space. If *nshell*=2 (default) 1s functions on first-row atoms, and 1s, 2s, and 2p-functions are eliminated on second-row atoms. Nothing is eliminated on H or He atoms. If effective core potentials are used, nothing is deleted at the corresponding atom. Also, functions are only

deleted if the norm of the projected function is below THRCOR (default 0.1)

DELBAS=*ibaso*

This parameter determines the method for eliminating redundant functions of pair domains.

ibaso=0: The space of normalized eigenvectors of $\tilde{\mathbf{S}}^{ij}$, which correspond to small eigenvalues, is eliminated (default if no gradients are computed).

ibaso > 0: individual basis functions are eliminated. The value of *ibaso* affects details of the method to determine redundant functions.

ibaso=1: Redundant functions eliminated from pair domains, using Jacobi method for diagonalization of overlap matrices. This is the default if properties or gradients are computed.

ibaso=2: Redundant functions are eliminated from pair domains, using Householder method for diagonalization of overlap matrices.

ibaso=3: Redundant functions are eliminated from orbital and pair domains, using Jacobi method for diagonalization of overlap matrices.

ibaso=4: Redundant functions are eliminated from orbital and pair domains, using Householder method for diagonalization of overlap matrices.

The diagonalization method has only an effect for DELEIG=1 if degenerate eigenvalues are present. *ibaso > 2* has only an effect if NONORM=0.

Parameters for selection of redundant functions if DELBAS > 0:

DELSHL=*idlshl*

This parameter determines if whole shells of basis functions (i.e., all *p*-functions for a given exponent at one atom) should be simultaneously eliminated. This may be useful in order to guarantee rotational invariance in geometry optimizations and frequency calculations.

idlshl=1: eliminate as many functions of a shell simultaneously as possible, but never more than determined by small eigenvalue of the overlap matrix (default).

idlshl=2: as *idlshl=1*, but also eliminate functions with identical norm simultaneously.

idlshl=3: eliminate all functions of a shell simultaneously, even if a larger number of functions is deleted than determined by small eigenvalues of the overlap matrix. This must be used with care, since very poor energies may sometimes result!

idlshl=4: as *idlshl=3*, but also eliminate functions with identical norm simultaneously.

idlshl > 4: as *idlshl=4*, but equivalent functions centred at all symmetry equivalent atoms are considered to form a shell (not recommended!).

TYPECHECK=*typechk*

If nonzero, activates basis function type restrictions in redundancy check. For a given atom, only basis functions corresponding to occupied atomic orbitals are allowed to be deleted. For instance, on first row atoms at most two *s*-functions and one *p*-shell will be deleted. No functions are deleted from hydrogen or He atoms.

DELEIG=*idleig*

This option determines how redundant basis functions are selected.

idleig=0: functions corresponding to the smallest diagonal elements of projected orbital matrix are eliminated.

idleig=1: Functions corresponding to the largest coefficients in the

eigenvectors of \tilde{S}^{ij} are deleted (default). Since degenerate eigenvectors can arbitrarily mix, the selection may not be unique and depend on the diagonalization method (see DELBAS).

DELCMIN=*value* Only effective with DELEIG=1. Only basis functions with coefficients larger than *value* in the eigenvectors of small eigenvalues can be deleted (default 0.1).

Parameters for multipole treatment of exchange operators

MULTMETHOD=*option* Used internally by the MULTP card - don't mess with it.

DSTMLT=*level* Determines the expansion level of the multipole expansion of distant pairs (e.g. 1 means dipole approximation, 2 quadrupole approximation and so on). Default is 0 (MULTP card: 3).

Parameters for energy partitioning:

IEPART=*value* enables/disables energy partitioning.
iepart=0: energy partitioning is disabled.
iepart=1: energy partitioning is enabled.
iepart=2: energy partitioning is enabled. Additionally, a list of all pair energies and their components is printed.

EPART=*cutoff* cutoff parameter to determine individual monomers in a cluster (i.e. centre groups). Should be somewhat larger than the largest intramolecular bond length (given in a.u.).

Miscellaneous options:

SKIPDIST=*skipdist* Test-parameter. Its value should only affect the efficiency but not influence the results.
skipdist=-1: weak and distant pairs are set to zero after MP2 but are not eliminated from the pair list and not skipped in any loop.
skipdist=0: No pairs are deleted from pair list, but weak and distant pairs are skipped in the loops were appropriate.
skipdist=1: Very distant pairs are neglected from the beginning. Distant pairs are eliminated after MP2.
skipdist=2: As skipdist=1, but also weak pairs are eliminated after MP2.
skipdist=3: As skipdist=2, but distant pairs are eliminated from the operator list in case of LMP2 with multipole approximations for distant pairs. This is the default.

ASYDOM=*jiterm* Enables the use of asymmetric domains for distant pairs. The asymmetric domain approximation supplements the multipole approximation for distant pairs, as it suppresses the treatment of configurations for which no integrals can be computed by multipole expansion. This leads to computational savings and improved numerical stability.
jiterm=0: Disable asymmetric domains.
jiterm=-1: Enable asymmetric domains (default).
jiterm=-2: Enable a variation of the asymmetric domain formalism: Exchange operators will initially be projected to the asymmetric domain instead of simply packed.

LOCSING=*locsing* If *locsing.ne.0*, the single excitations use the full space, i.e., they are not treated locally. This is only works for LOCAL=1.

MAXANG=*lmax* The purpose of this experimental option is to reduce the basis set sensitivity of the Boughton-Pulay (BP) method for domain selection.

Only basis functions with angular momentum up to $l_{max}-1$ are included when computing the overlap of the approximate and exact orbitals. For example, $MAXANG=2$ means to omit all contributions of d , f and higher angular momentum functions. To obtain reasonable domains, the value of $DOMSEL$ must often be reduced (to 0.97 or so). This option should only be used with care!

$MAXBP=maxbp$	If $maxbp=1$, the atoms are ranked according to their contribution to the Boughton-Pulay overlap (default); this should normally give the smallest and best orbital domains. If $maxbp=0$, the atoms are ranked according to Mulliken charges. In both cases atoms with Mulliken charges greater than 0.6 are always included, and atoms with the same Mulliken charges are added as groups.
$MULLIKEN=option$	Determines method to determine atomic charges. $option=0$: diagonal elements of $\mathbf{S}^{\frac{1}{2}}\mathbf{C}$ are used. $option=1$: Mulliken gross
$PIPEKAO=option$	If $option \geq 0$, the orbitals are localized by maximizing the coefficients of basis functions of a given type at a given atom. Normally, this is only useful to uniquely define degenerate orbitals in atoms. For instance, when this option is used to localize the orbitals for a dimer like $(Ar)_2$ at a very long distance, clean s , p_x , p_y , and p_z atomic orbitals will be obtained. It is not recommended to use this option for molecular calculations!
$NONORM=value$	Determines if projected functions are normalized (not recommended). $value=-1$: projected orbitals are normalized before redundancy check. $value=0$: projected orbitals are normalized after redundancy check (default). $value=1$: projected orbitals are normalized in redundancy check, afterwards unnormalized. $value=2$: projected orbitals are never normalized.
$LMP2ALGO=value$	If nonzero, use low-order scaling method in LMP2 iterations. This may require more CPU time in calculations for smaller molecules.
$OLDDEF=value$	For compatibility with older versions: if nonzero, revert to old defaults. Options set before this may be overwritten.

Thresholds:

$THRPIP=thresh$	Threshold for Pipek-Mezey localization. The localization is assumed to be converged if all 2×2 rotation angles are smaller than $thresh$. The default is $1.d-12$. It can also be modified globally using $GTHRESH$, $LOCALI=thresh$. Note that $GTHRESH$ is not an input command of the local program and must be given before the <code>METHOD</code> card.
$THRORB=thresh$	Threshold for eliminating functions from pair domains whose norm is smaller than $thresh$ after projecting out the occupied space. The default is $thorb=1.d-6$.
$THRLOC=thresh$	Threshold for eliminating redundant basis functions from pair domains. For each eigenvalue of $\tilde{\mathbf{S}}^{ij} < thresh$ one function is deleted. The default is $1.d-6$. The method used for deleting functions depends on the parameters <code>IDLEIG</code> and <code>IBASO</code> .
$THRMP2=thresh$	Threshold for neglecting small fock matrix couplings in the LMP2 iterations (default $1.d-8$). Specifying a larger threshold speeds up the iterations but may lead to small errors in the energy. In the initial

iterations, a larger threshold is chosen automatically. It is gradually reduced to the specified final value during the iterations.

THRCOR=*thresh* Threshold for deleting projected core orbitals. The functions are only deleted if their norm is smaller than *thresh* (default 0.1)

The thresholds can also be specified on the THRESH card.

23.7 Additional options available on the ATTENUATE card

The defaults reported for the following keys are likely to change in the future.

Most important options

DECAY= ω This is the decay parameter that determines the splitting of the Coulomb operator in the split approach. Larger values of ω put more weight to the long range part of the operator, which means that the multipole correction will have more difficulties to converge but the transformation of the short range part will be faster. Default: 0.20

SHORTMLT=*level* Determines the expansion level of monopolar multipole expansions in the context of the split Coulomb operator approach. Default: 15

LONGMLT=*level* Determines the expansion level of bipolar multipole expansions in the context of the split Coulomb operator approach. Default: 13

Specifying which integrals to treat by which multipole expansion type

RMAIN=*distance* When the distance between two orbitals is closer than the absolute value of *distance*, multipole corrections (in the context of the split approximation) will be carried out as monopolar expansions, otherwise a more sophisticated approach will be used involving four expansions for each pair, one for the ordinary dispersion block of the exchange matrix for the given pair, two for the two ionic exchange blocks and one for the exchange-dispersion block (see the section on energy partitioning for explanation of these terms). If *distance* is a positive value, the distance between the orbitals will be taken to be the distance between the closest atoms of the two orbital domains (as in WEAKPAIR, DISTPAIR etc.), if it is a negative value, the distance of the centroids of the two orbitals will be considered. Default: 1

RIONIC=*distance* When the distance between two orbitals is closer than *distance* [a.u.], the multipole correction of ionic blocks of the exchange operators will be carried out as monopolar expansion, otherwise a bipolar expansion will be performed. The distance is understood as the distance of the centroids of the two orbitals. Default: 0

SUPPRESS=*distance* When the distance between two orbitals (distance of centroids) is higher than the absolute value of *distance* [a.u.], the multipole correction of the exchange-dispersion blocks of the exchange operators is suppressed. If a positive value is given, the multipole correction of the ionic exchange blocks is also suppressed. A zero value disables the suppression of multipole corrections. Default: 0

Options for least squares fit generation of interaction coefficients

<i>FITMLTP=option</i>	Specifies how the coefficients for the multipole expansion of long range integrals are calculated. <i>option=0</i> : Taylor expansion <i>option=1</i> : Least squares fit Default: 1
<i>F1DGRID=value</i>	Sets the number of quadrature points used to generate integrals that arise in the one-dimensional fit (i.e. for the monopolar multipole expansion). Default: 50
<i>F2DGRIDR=value</i>	Sets the number of quadrature points used along the <i>r</i> (radial) coordinate when generating integrals for the two-dimensional fit (i.e. for the bipolar multipole expansion). Default: 50
<i>F2DGRIDP=value</i>	Sets the number of quadrature points used along the ϕ (angular) coordinate for the two-dimensional integrals. Default: 20
<i>F1DBORDER=value</i>	If greater than 0, sets the upper bound of integration (in bohr) and selects the Gauss–Legendre quadrature for the one-dimensional integrals. If 0, selects Gauss–Laguerre quadrature. Default: 0
<i>F2DBORDER=value</i>	The same for the two-dimensional integrals.
<i>F1DGAMMA=γ</i>	Sets the negative exponent of the weight function for the one-dimensional fit. Smaller values are better for more diffuse densities. Should be positive (leading to a negative exponent). Default: 1.7
<i>F2DGAMMA=γ</i>	The same for the two-dimensional fit.
<i>WEIGHT3D=option</i>	Selects what type of weight function is used for the fits. <i>option=0</i> : Flat <i>option=1</i> : Spatial Remember to change <i>F1DGAMMA</i> / <i>F2DGAMMA</i> accordingly when using a flat weight function, $\gamma = 1.0$ is then a reasonable value. Default: 1

Options for determination of batches

<i>NUMBATCH=value</i>	If 0, selects automatic determination of the number of batches. If >0, provides a manual override for this number. Default: 0
<i>BATCHDIAM=value</i>	Maximal acceptable diameter of batches. Used to automatically determine the number of batches. If zero, disables batches (i.e., only one batch will be created). Default: 35
<i>BATCHALGO=option</i>	Selects the algorithm used to determine the batches. <i>option=0</i> : Manually set batch centres. Three arrays with name <i>BATCHX</i> , <i>BATCHY</i> , <i>BATCHZ</i> have to be set in the input before the <i>MP2</i> card that contain the <i>x</i> -, <i>y</i> - and <i>z</i> -coordinates of the desired batch centres. <i>option=1</i> : Simple algorithm: Determines a path through the molecules and distributes the batch centres along this path. <i>option=2</i> : A robust simulated annealing algorithm. Will only use atom positions as batch centres, might therefore fail for strongly separated dimers and similar systems where there is no atom near the optimal batch centre position. <i>option=3,4</i> : A less robust simulated annealing algorithm that directly tries to optimize the batches instead of the batch centres. Default: 2
<i>WEIGHTPREV=value</i>	Affects the one-dimensional path laid through the molecule for <i>BATCHALGO</i> =1. Smaller values mean a more systematic, directed order of atoms from

one end to the other. Larger values mean that the distance from one atom in the path to the next will become smaller. Has to be between 0 and 1. Default: 0.5

RANSEED=*value* Negative values initialize the random number generator for the simulated annealing algorithms. Positive numbers suppress initialisation of the random number generator. Default: -1

Further numerical stability options

CUTOFF=*distance* Applies a simple cutoff to orbitals before the transformation of the multipole operators. Orbital coefficients belonging to AOs that are more than *distance* [a.u.] away from the orbital centre will be deleted. *distance*=0 means don't use a cutoff. Default: 15

MONOPOLE=*option* Specifies how to treat monopole integrals. Monopole integrals should be zero due to the orthogonality between occupied and virtual orbitals. Therefore, they are usually not included in the calculation. However, this does not hold exactly when an orbital cutoff is applied. Including monopole integrals in the calculation might therefore improve the numerical stability.

option=0: Neglect monopole integrals
option=1: Use monopole integrals in the translation, but neglect them later on
option=2: Use monopole integrals everywhere (translation and expansion), only neglect monopole- 2^{p+1} pole interactions, where p is the requested multipole level
option=3: Use monopole integrals everywhere (translation and expansion), but neglect all monopole- 2^{p+1} pole, dipole- 2^p pole, quadrupole- 2^{p-1} pole interactions and so on. This is entirely consistent, but reduces the effective multipole level by 1 as compared to the other options.

Default: 1

Multipole operators

MAXMLTPL=*option* Defines the highest level of multipole operators that are created. Has to be greater or equal than each of DSTMLT, LONGMLT and SHORTMLT, is otherwise overwritten by the default, which is max(DSTMLT, LONGMLT, SHORTMLT). Greater values are a useless waste of CPU unless you save the operators for later reuse.

MULTPAGE=*option* *option*=0: Suppress paging of multipole operators during multipole expansion. After their creation, all operators are read from disk into memory. Will crash if not enough memory is available.
option=1: Read operators from disk when needed. Small performance impact.

Default: 1

Essentially obsolete keys (for Taylor expansions)

TRUNCATE=*option* Determines if the simple or exhaustive truncation scheme will be used for multipole expansions. Exhaustive truncation means that, unlike a classical multipole expansion, all interactions of multipoles up to the highest order are taken into account, e. g. in an expansion of level two, the exhaustive scheme will include quadrupole-quadrupole terms while the simple scheme won't. Valid options are:
option=0: Always use simple truncation.

option=1: Use simple truncation for bipolar expansions, but exhaustive truncation for monopolar expansions. This significantly improves convergence when a Taylor expansion is used, but also accelerates the onset of divergent behaviour for large expansion levels and should therefore always be used in connection with a damping function for multipole operators.

option=2: Use simple truncation for monopolar expansions, but exhaustive truncation for bipolar expansions including distant pair treatment.

option=3: Always use exhaustive truncation.

The latter two options are not useful and only included for the sake of completeness. Default: 0

DAMP=*order*

Specifies the form of a damping function that is applied to orbitals before the transformation of the multipole operators. *order* is actually the order of a Taylor series that is part of the function and that mimics the behaviour of the monopolar multipole expansion. Only multiples of 4 are reasonable values. Higher values mean a harder damping function (i. e. less damping at short distances, but more at long distances). A value of 0 disables the damping. A value of 8 is recommended when using a Taylor expansion. Default: 0

SCALEDAMP=*scalefactor* The damping function can additionally be scaled by the value of *scalefactor*. A value lower than 1 means the damping function will be squeezed, while a value higher than 1 will cause it to be stretched. A value of 0 disables the scaling (as well as 1). A value of 0.9 is recommended when using a Taylor expansion. Default: 0

Stuff for debugging

PAIREN=*option*

Prints a list of uncoupled pair energies before the MP2 iterations. Can be used as a convenient diagnostic when getting totally implausible correlation energies or no convergence or 'UNREASONABLE NORM' messages.

Table 9: Summary of local (*multp*) options and their default values

Parameter	Alias	Default value	Meaning
General Parameters:			
LOCAL		4	determines which program to use
SAVE	SAVDOM	0	specifies record for saving domain info
START	RESTDOM	0	specifies record for reading domain info
PIPEK	LOCORB	0	activates or deactivates PM localization
SAVORB	SAVLOC	0	specifies record for saving local orbitals
DOMONLY		0	if set, only domains are made. if 2, only orbital domains are made.
Parameters to select weak and distant pairs:			
WEAKPAIR	WEAKP	1	criterion for selection of weak pairs
DISTPAIR	DISTP	0 (8)	criterion for selection of distant pairs
VERYDIST	VERYD	0 (15)	criterion for selection of very distant pairs
Parameters to define domains:			
DOMSEL	CHGFRAC	0.98	selection criterion for orbital domains
DELCOR	IDLCOR	2	delete projected core AOs up to certain shell
DELBAS	IBASO	0	determines how to remove redundancies
Parameters for redundancy check using DOMSEL=1			
TYPECHECK	TYPECHK	1	activates basis function type restrictions
DELSHL	IDLSHL	1	determines if whole shells are to be deleted
DELEIG	IDLEIG	1	determines how to select redundant functions
DELCMIN	CDELMIN	0.1	parameter for use with DELEIG=1
Parameter for multipole treatment of exchange operators:			
DSTMLT		0 (3)	expansion level for distant pairs
Parameters for energy partitioning:			
IEPART		0	If nonzero: do energy partitioning
EPART		3.0	cutoff parameter for determining individual monomers
Miscellaneous options:			
SKIPDIST	SKIPD	3	determines at which stage weak and distant pairs are eliminated
ASYDOM	JITERM	0	parameter for use of asymmetric domains
LOCSING	LOCSNG	0	determines virtual space used for singles
MAXANG	MAXL	99	restriction for Boughton-Pulay domain selection
CHGMIN		0.01	minimum Mulliken charge for BP domain selection
CHGMINH		0.05	minimum Mulliken charge of H-atoms for BP domain selection
CHGMAX		0.40	If charge larger than this value, atom is always included
MAXBP	MAXBP	0	determines how to rank atoms in Boughton-Pulay domain selection
MULLIKEN	LOCMUL	0	determines how to rank atoms for domains
PIPEKAO	LOCAO	0	activates AO localization criterion
NONORM		2	determines whether projected functions are normalized
LMP2ALGO	MP2ALGO	1	if nonzero, use low-order scaling method in LMP22 iterations
OLDDEF		0	allows to revert to older defaults
Thresholds:			
THRPIP		1.d-12	Threshold for Pipek-Mezey localization.
THRORB		1.d-6	Threshold for eliminating projected orbitals with small norm.
THRLOC		1.d-6	Threshold for eliminating redundant projected orbitals.
THRCOR		1.d-1	Threshold for eliminating projected core orbitals.
THRMP2		1.d-8	Threshold for neglecting small fock matrix elements in the LMP22 iterations.

Table 10: Summary of `attenuate` options and their default values

Parameter	Default value	Meaning
Most important options:		
DECAY	0.20	split parameter ω
SHORTMLT	15	level p of monopolar multipole expansion
LONGMLT	13	level p of bipolar multipole expansion
Specifying which integrals to treat by which multipole expansion type:		
RMAIN	1	when to switch from monopolar to four-block treatment
RIONIC	0	when to switch from monopolar to bipolar treatment of ionic blocks
SUPPRESS	0	when to suppress cross-excited blocks
Options for least squares fit generation of interaction coefficients:		
FITMLTP	1	use least squares fit instead of Taylor
F1DGRID	50	no. of quadrature points for 1D fit
F2DGRIDR	50	no. of quadrature points for 2D fit r
F2DGRIDP	20	no. of quadrature points for 2D fit ϕ
F1DBORDER	0	end of integration interval for 1D fit
F2DBORDER	0	end of integration interval for 2D fit r
F1DGAMMA	1.7	negative exponent of weight function for 1D fit
F2DGAMMA	1.7	negative exponent of weight function for 2D fit
WEIGHT3D	1	use spacial instead of flat weight function
Options for determination of batches:		
NUMBATCH	0	manually set number of batches
BATCHDIAM	35	maximal diameter of batches
BATCHALGO	2	algorithm to determine batches
WEIGHTPREV	0.5	parameter for algorithm BATCHALGO=1
RANSEED	-1	initialize random number generator for simulated annealing
Further numerical stability options:		
CUTOFF	15	orbital cutoff
MONOPOLE	1	if and how to treat monopole integrals
Multipole operators:		
MAXMLTPL	<i>auto</i>	manually set level of multipole operators to create
MULTPAGE	1	turn on paging of multipole operators during multipole expansion
Essentially obsolete keys (for Taylor expansion):		
TRUNCATE	0	truncation pattern of multipole expansion
DAMP	0	damping function for orbitals
SCALEDAMP	0	scaling factor for the damping function
Stuff for debugging:		
PAIREN	0	print a list of uncoupled pair energies

24 THE FULL CI PROGRAM

This module is the determinant full CI program, as described in

P.J. Knowles and N.C. Handy, Chem. Phys. Letters 111 (1984) 315,
P.J. Knowles and N.C. Handy, Comp. Phys. Commun. 54 (1989) 75.

Published work resulting from the use of this program should cite these references.

The program in normal use finds the lowest eigenvector of the complete CI hamiltonian matrix; more sophisticated use is possible, but not documented here. The program is interfaced to free standing versions such as supplied in the CPC program library by use of the DUMP option.

The program is called with the command FCI.

24.1 Defining the orbitals

ORBIT,*name.file*;

name.file specifies the record from which orbitals are read. The default is the set of orbitals from the last SCF, MCSCF or CI calculation.

24.2 Occupied orbitals

OCC,*n*₁,*n*₂, . . . ,*n*₈;

*n*_{*i*} specifies numbers of occupied orbitals (including CORE) in irreducible representation number *i*. If not given, the default is the complete basis set.

24.3 Frozen-core orbitals

CORE,*n*₁,*n*₂, . . . ,*n*₈;

*n*_{*i*} is the number of frozen-core orbitals in irrep number *i*. These orbitals are doubly occupied in all configurations, i.e., not correlated. If no CORE card is given, the program uses the same core orbitals as the last CI calculation; if there was none, then the atomic inner shells are taken as core. To avoid this behaviour and correlate all electrons, specify

CORE

24.4 Defining the state symmetry

The number of electrons and the total symmetry of the wavefunction are specified on the WF card:

WF,*elec,sym,spin*

where

<i>elec</i> :	is the number of electrons
<i>sym</i> :	is the number of the irreducible representation
<i>spin</i> :	defines the spin symmetry, <i>spin</i> = 2 <i>S</i> (singlet=0, doublet=1, triplet=2, etc.)

24.5 Printing options

PRINT,*code,value*;

Print options. Generally, the value determines how much intermediate information is printed. *value*=-1 means no print (default for all codes). if *value* is omitted, it is taken as zero, which is usually appropriate. Specification of higher values will generate more output. The following codes are allowed:

ORBITAL	Print molecular orbitals
INTEGRAL	Print integrals
TIMING	Print extra timing information
DIAGONAL	Print diagonal elements of Hamiltonian
HAMILTONIAN	Print much intermediate information

24.6 Interface to other programs

DUMP;

causes the FCI diagonalization to be bypassed, with input information and transformed integrals being written to a formatted file FCIDUMP. The format is as described in Comp. Phys. Commun. 54 (1989) 75.

25 PROPERTIES AND EXPECTATION VALUES

25.1 The property program

The property program allows the evaluation of one-electron operators and expectation values. Normally, the operators are computed automatically when using the global GEXPEC directive (see section 4.13) or the EXPECT or TRAN commands in the SCF, MCSCF, and CI programs. The explicit use of the property program is only necessary in the rare case that the user is interested in an orbital analysis of the properties.

25.1.1 Calling the property program (PROPERTY)

PROPERTY

invokes the property program.

25.1.2 Expectation values (DENSITY)

DENSITY [*record.file*] [*specifications*]

If this card is present, the density matrix will be read from record *record.file* and property expectation values will be calculated. If the specification *record.file* is omitted, the last dump record is used. Density matrices for specific states can be selected using *specifications*, as explained in section 2.16. Note that the density matrices are stored in the same record as the orbitals.

25.1.3 Orbital analysis (ORBITAL)

ORBITAL [*record.file*] [*specifications*]

If this card is present, the orbitals are read from record *record.file* and an orbital analysis of the expectation values is printed (the density matrix must also be provided!). If *record.file* is omitted, the last dump record is used. This is only meaningful for diagonal density matrices (SCF or natural orbitals). Natural orbitals for specific states can be selected using *specifications*, as explained in section 2.16.

25.1.4 Specification of one-electron operators

The required operators are specified by code words. Optionally, the geometry or the nuclear centre at which the operator is computed can be specified.

For each operator, an input card of the following form is required:

code,centre,x,y,z,,factor

code specifies the property. The available operators are given in section 4.13.

The other parameters have the following meaning:

centre row number of Z-matrix or atomic symbol defining the centre at which property shall be calculated; if *centre* ≠ 0 you need not read in coordinates.

x,y,z cartesian coordinates of the point (only if *centre*=0).

factor the operator is multiplied by this factor. The default is *factor*=1 except for REL. In this cases proper factors for relativistic corrections are used unless *factor* is given. The two commas before factor are needed to preserve compatibility with Molpro96.

25.1.5 Printing options

PRINT,*print*

This card is used to control output, mainly for debugging purposes.

print= 0 no test output (default)

print> 0 operators are printed.

25.1.6 Examples

The following example computes the dipole quadrupole moments of water and prints an orbital analysis. By default, the origin is at the centre of mass, and this is taken as origin for the quadrupole moments.

```

***,h2o properties
geometry={o;h1,o,r;h2,o,r,h1,theta} !Z-matrix geometry input
r=1 ang !bond length
theta=104 !bond angle
hf !do scf calculation
property !call property program
orbital !read scf orbitals
density !read scf density matrix
dm !compute dipole moments and print orbital contributions
qm !compute quadrupole moments and print orbital contributions
multi;state,2 !do full-valence CASSCF
natorb,state=1.1 !compute natural orbitals for state 1.1
natorb,state=2.1 !compute natural orbitals for state 2.1

property !call property program
orbital,state=1.1 !read casscf natural orbitals for state 1.1
density,state=1.1 !read casscf density matrix for state 1.1
dm !compute dipole moments and print orbital contributions
qm !compute quadrupole moments and print orbital contributions

property !call property program
orbital,state=2.1 !read casscf natural orbitals for state 2.1
density,state=2.1 !read casscf density matrix for state 2.1
dm !compute dipole moments and print orbital contributions
qm !compute quadrupole moments and print orbital contributions

```

Alternatively, the dipole and quadrupole moments can be computed directly in the SCF and MCSCF programs, but in this case no orbital contributions are printed: program

```

***,h2o properties
geometry={o;h1,o,r;h2,o,r,h1,theta} !Z-matrix geometry input
r=1 ang !bond length
theta=104 !bond angle
gexpec,dm,qm !global request of dipole and quadrupole moments
hf !do scf calculation examples/
multi;state,2 !do full-valence CASSCF h2o_gexpec1.com
natorb,state=1.1 !compute natural orbitals for state 1.1
natorb,state=2.1 !compute natural orbitals for state 2.1

```

25.2 Distributed multipole analysis

Any density matrix can be analysed using the distributed multipole analysis described by Stone, Chem. Phys. Letters (1981), 83, 233. The multipole moments arising from the overlap of each pair of primitives are calculated with respect to the overlap centre, and then shifted to the nearest of a number of *multipole sites*. By default these comprise all atoms specified in the integral input. However the list of multipole sites can be modified by deleting and/or adding sites, and also by restricting the rank of multipole which may be transferred to any given site. The atomic charges are stored in the MOLPRO variable ATCHARGE. The *i*'th element in ATCHARGE corresponds to the *i*'th row of the Z-matrix input.

Options may appear in any order, except DENSITY, which must be first if given.

The present version does not allow generally contracted AO basis sets.

25.2.1 Calling the DMA program (DMA)

```
DMA ;
```

This command initializes the DMA program.

25.2.2 Specifying the density matrix (DENSITY)

```
DENSITY,record.file [,specifications]
```

The density matrix to be analysed is that found in record *record* on file *file*. If omitted, *record.file* defaults to current orbital record. If specified, DENSITY must appear first in the input. Density matrices for specific states can be selected using *specifications*, as explained in section 2.16.

25.2.3 Linear molecules (LINEAR, GENERAL)

```
GENERAL;
```

(default) invokes the normal program, which copes with any geometry.

```
LINEAR
```

invokes a faster program which can be used when all the atoms are arranged parallel to the *z*-axis and only the $m = 0$ components of the multipoles are required.

25.2.4 Maximum rank of multipoles (LIMIT)

LIMIT,*name*,*lmax*;

lmax is the highest rank of multipole that is to be calculated by the program. Default (and maximum) is 10 for the general program and 20 for the linear one. If *name* is specified, the limit applies only to multipole site *name*.

25.2.5 Omitting nuclear contributions (NONUCLEAR)

NONUCLEAR

The nuclear contributions to properties are not to be evaluated.

25.2.6 Specification of multipole sites (ADD, DELETE)

ADD,*name*,*x*,*y*,*z*,*lmax*,*radius*;

Add a new site at (*x*, *y*, *z*) with the *name* specified. The multipole rank is limited to *lmax* if a value is specified, otherwise the value of *lmax* specified by the LIMIT directive is used. No account is taken of symmetry; every site in a symmetry-equivalent set must be specified explicitly. The *radius* of the site may also be specified (default 1.0).

DELETE,*name*

Delete all atoms with the *name* given from consideration as a multipole site. Note that original atoms from the integral program have names 1, 2, 3, ... as printed in integral output. DELETE , ALL deletes all atoms and gives the multipoles with respect to the origin only.

25.2.7 Defining the radius of multipole sites (RADIUS)

RADIUS,*name*,*r*;

Assign radius *r* to all sites with the *name* given. The program moves multipoles at an overlap centre *P* to the site *S* for which the value of $|P - S|/r(S)$ is smallest. In the absence of a RADIUS directive, all sites are given radius 1.

25.2.8 Notes and references

The multipoles produced by this analysis are given in their spherical harmonic definitions. Explicit formulae for translating between the cartesian and spherical harmonic definitions of the multipole moments are given in, *Explicit formulae for the electrostatic energy, forces and torques between a pair of molecules of arbitrary symmetry*, S. L. Price, A. J. Stone, and M. Alderton, Molec. Phys., 52, 987 (1984).

For examples of the use of DMA analysis see, Price and Stone, Chem. Phys. Lett., 98, 419 (1983); Buckingham and Fowler, J. Chem. Phys., 79, 6426 (1983).

25.2.9 Examples

The following input calculates SCF multipole moments for water.

```

***,h2o distributed multipole analysis
geometry={o;h1,o,r;h2,o,r,h1,theta} !Z-matrix geometry input
r=1 ang !bond length
theta=104 !bond angle
basis=6-311g**
hf !do scf calculation
dma;limit,,4 !results for total multipoles are

```

examples/
h2o_dma.com

25.3 Mulliken population analysis

25.3.1 Calling the population analysis program (POP)

POP;

Invokes Mulliken analysis program, which analyses any density matrix into its contributions from s,p,d,f... basis functions on each atom. The density matrix is taken from the last dump record, unless overridden with the DENSITY card. The subcommands may be abbreviated by the first four characters. The atomic charges are stored in the MOLPRO variable ATCHARGE. The *i*'th element in ATCHARGE corresponds to the *i*'th row of the Z-matrix input.

25.3.2 Defining the density matrix (DENSITY)

DENSITY,*record.file* [*,specifications*]

Take density matrix to be analysed from record *record* on file *file*. Density matrices for specific states can be selected using *specifications*, as explained in section 2.16. Note that the density matrices are stored in the same record as the orbitals.

25.3.3 Populations of basis functions (INDIVIDUAL)

INDIVIDUAL;

25.3.4 Example

```

***,h2o population analysis
geometry={o;h1,o,r;h2,o,r,h1,theta} !Z-matrix geometry input
r=1 ang !bond length
theta=104 !bond angle
basis=6-311g**
hf !do scf calculation
pop; !Mulliken population analysis using mcscf density
individual !give occupations of individual basis functions

```

examples/
h2o_pop.com

If specified, the Mulliken populations of each individual basis function are printed.

25.4 Finite field calculations

Dipole moments, quadrupole moments etc. and the corresponding polarizabilities can be obtained as energy derivatives by the finite difference approximation. This is most easily done with the DIP, QUAD, or FIELD commands. An error will result if the added perturbation is not totally symmetric (symmetry 1). Note that the orbitals must be recomputed before performing a correlation calculation.

25.4.1 Dipole fields (DIP)

DIP,*xfield,yfield,zfield*;
 DIP+,*xfield,yfield,zfield*;

Add a finite dipole field to the one electron Hamiltonian and the core energy. The field strength is given by *xfield,yfield,zfield*. DIP+ adds to any existing field, otherwise any previous field is removed.

25.4.2 Quadrupole fields (QUAD)

QUAD,*xxfield,yyfield,zzfield,xyfield,xzfield,yzfield*;
 QUAD+,*xxfield,yyfield,zzfield,xyfield,xzfield,yzfield*;

Exactly as the DIP command, but adds a quadrupole field.

25.4.3 General fields (FIELD)

FIELD,*oper1,fac1, oper2,fac2, ...*;
 FIELD+,*oper1,fac1, oper2,fac2, ...*;

Adds one-electron operators *oper1, oper2, ...* with the corresponding factors *fac1, fac2, ...* to the one-electron hamiltonian. The available operators are given in section 4.13. An error will result if the added perturbation is not totally symmetric (symmetry 1).

FIELD+ adds to any existing field, otherwise any previous field is removed.

Note that FIELD does currently not modify core polarization potentials (CPP). If CPPs are present, only DIP and QUAD should be used.

25.4.4 Examples

The first examples shows various possibilities to add perturbations to the one-electron hamiltonian.

```

***,H2O finite fields
memory,4,m
R      =    0.96488518 ANG
THETA =  101.90140469
geometry={H1
          O,H1,R;
          H2,O,R,H1,THETA}
hf;wf,10,1;          !scf without field

f=0.05

dip,,,f              !add dipole (z) field to h0
hf                   !do scf with modified h0

field,dmz,f          !add dipole (z) field to H0
                    !same result as previous example
hf                   !do scf with modified h0

quad,,,f             !add quadrupole (qmzz) field to h0
hf                   !do scf with modified h0

field,qmzz,f         !add quadrupole (qmzz) field to h0;
                    !same result as previous example
hf                   !do scf with modified h0

field,zz,f,xx,-0.5*f,yy,-0.5*f
                    !add general field; same result as quad above
hf                   !do scf with modified h0

field,zz,f           !same as before with separate field commands
field+,xx,-0.5*f
field+,yy,-0.5*f
hf                   !do scf with modified h0

field                !remove field
hf                   !scf without field

```

examples/
field.com

The second example shows how to compute dipole moments and polarizabilities using finite fields.

```

***,H2O finite field calculations

r=1.85,theta=104                !set geometry parameters
geometry={0;                    !z-matrix input
          H1,O,r;
          H2,O,r,H1,theta}
basis=avtz                      !define default basis
field=[0,0.005,-0.005]         !define finite field strengths
method=[hf,mp4,ccsd(t),casscf,mrci]

k=0
do i=1,#field                   !loop over fields
  dip,,,field(i)               !add finite field to H
  do m=1,#method               !loop over methods
    k=k+1
    $method(m)                 !calculate energy
    e(k)=energy                !save energy
  enddo
enddo

k=0
n=#method
do m=1,#method
  k=k+1
  energ(m)=e(k)
  dipmz(m)=(e(k+n)-e(k+2*n))/(field(2)-field(3)) !dipole moment as first energy derivative
  dpolz(m)=(e(k+n)+e(k+2*n)-2*e(k))/((field(2)-field(1))*(field(3)-field(1))) !polarizability
enddo

table,method,energy,dipmz,dpolz
title,results for H2O, r=$R, theta=$theta, basis=$basis
---
```

examples/
h2o_field.com

25.5 Relativistic corrections

Relativistic corrections may be calculated within the Cowan-Griffin approach by computing expectation values of the mass-velocity and 1-electron Darwin integrals; these should be generated using the property integral program with keyword REL. The expectation values can be computed within the SCF, MCSCF and CI programs in the usual way using the EXPECT command, again with the keyword REL. The mass-velocity and Darwin terms, and their sum are subsequently available through the MOLPRO variables MASSV, DARW and EREL respectively.

25.5.1 Example

```

***,ar2
geometry={ar1;ar2,ar1,r}       !geometry definition
r=2.5 ang                     !bond distance
hf;                           !non-relativistic scf calculation
expec,rel,darwin,massv        !compute relativistic correction using Cowan-Griffin operator
e_nrel=energy                 !save non-relativistic energy in variable enrel
show,massv,darwin,erel       !show individual contribution and their sum

dkroll=1                      !use douglas-kroll one-electron integrals
hf;                           !relativistic scf calculation
e_dk=energy                   !save relativistic scf energy in variable e_dk.
show,massv,darwin,erel       !show mass-velocity and darwin contributions and their sum
show,e_dk-e_nrel             !show relativistic correction using Douglas-Kroll
```

examples/
ar2_rel.com

25.6 CUBE — dump density or orbital values

CUBE,*filename*,*iflag*,*n₁*,*n₂*,*n₃*

calls a module which dumps the values of various properties on a spatial parallelepipedal grid to an external file. The purpose is to allow plotting of orbitals, densities and other quantities by external programs. The format of the file is intended to be the same as that produced by other programs.

<i>filename</i>	is the unix path name of the file to be written, and its specification is mandatory.
<i>iflag</i>	If <i>iflag</i> is negative (default), a formatted file will be written, otherwise unformatted fortran i/o will be used.
<i>n₁</i> , <i>n₂</i> , <i>n₃</i>	specify the number of grid points in each of three dimensions. If not specified, sensible defaults are chosen.

By default, the last density computed is evaluated on the grid, and written to *filename*. This behaviour can be modified by one or more of the following subcommands.

25.6.1 DENSITY — source of density

DENSITY,[*density-source*]
 GRADIENT,[*density-source*]
 LAPLACIAN,[*density-source*]

Compute the density and, optionally, its gradient and laplacian. <*density-source*> may be a record number containing the required density, and may contain further qualification, such as set number, in the usual way. By default, the last computed density is taken.

25.6.2 ORBITAL — source of orbitals

ORBITAL,[*orbital-list*],[*orbital-source*]

<*orbital-list*> is a list of one or more orbital numbers of the form *number.symmetry* or keywords chosen from HOMO, LUMO, OCC, ALL. If nothing is specified, the default is HOMO. <*orbital-source*> may be a record number containing the required density, and may contain further qualification, such as set number, in the usual way. By default, the last computed orbitals are taken.

Note that the CUBE file format precludes simultaneous orbital and density dumps, but that this may be achieved in the GOPENMOL format (see 25.7).

25.6.3 AXIS — direction of grid axes

AXIS,*x*,*y*,*z*

x,*y*,*z* specify the unnormalised direction cosines of one of the three axes defining the grid. Up to three AXIS commands can be given, but none is required. Axes need not be orthogonal. By default, the first axis is the cartesian *x*, the second is orthogonal to the first and to the cartesian *z*, and the third is orthogonal to the first two.

25.6.4 BRAGG — spatial extent of grid

Based on the direction of the coordinate axes, a parallelepiped (in the usual case of orthogonal axes, a cuboid) is constructed to contain the molecule completely. The atoms are assumed to be spherical, with an extent proportional to their Bragg radii, and the constant of proportionality can be changed from the default value using

BRAGG,*scale*

After the parallelepiped has been constructed, the grid is laid out with equal spacing to cover it using the number of points specified on the CUBE command.

25.6.5 ORIGIN — centroid of grid

ORIGIN,*x,y,z*

x,y,z specify the centroid of the grid. It is usually not necessary to use this option, since the default should suffice for most purposes.

25.6.6 Format of cube file

The formatted cube file contains the following records

(A)	job title.
(A)	brief description of the file contents.
(I5,3F12.6)	number of atoms, coordinates of grid origin (bohr).
(I5,3F12.6)	number of grid points n_1 , step vector for first grid dimension.
(I5,3F12.6)	number of grid points n_2 , step vector for second grid dimension.
(I5,3F12.6)	number of grid points n_3 , step vector for third grid dimension.
(I5,4F12.6)	atomic number, charge and coordinates; one such record for each atom.
(6E13.5)	$n_1 \times n_2$ records of length n_3 containing the values of the density or orbital at each grid point. In the case of a number of orbitals m , the record length is $m \times n_3$, with the data for a single grid point grouped together. In the case of the density gradient, there is first a record of length n_3 containing the density, then one of length $3n_3$ containing the gradient, with the three cartesian components contiguous. For the laplacian, there is a further record of length n_3 .

25.7 GOPENMOL — calculate grids for visualization in gOpenMol

GOPENMOL,*filename,iflag,n1,n2,n3*

The syntax and sub-options are exactly the same as for CUBE, except that the files produced are in a format that can be used directly in the gOpenMol visualization program. The following should be noted.

- Only the base name (up to the last '.') in *filename* is used, and is appended by different suffices to create several different files:

<code>.crd</code>	A CHARMM CRD-format file containing the coordinates is always produced, and may be used in the invocation of gOpenMol: <code>rungOpenMol -ifilename.crd</code>
<code>_density.plt</code>	If DENSITY is given, then the file <code>filename_density.plt</code> is produced and contains the density grid in gOpenMol internal format.
<code>_orbital_number.symmetry.plt</code>	If ORBITAL is given, then for each orbital <code>number.symmetry</code> specified, the file <code>filename_orbital_number.symmetry.plt</code> is produced and contains the orbital grid in gOpenMol internal format.

- The default is not to produce any orbitals or densities, and so only the atomic coordinates are dumped.
- The default is to use unformatted binary files, and this should not normally be changed.
- The ORIGIN and AXIS commands should not be used.
- If

INTERACT

is given in the input, when all the grids have been calculated, an attempt is made to start gOpenMol by executing the Unix command `rungOpenMol`. If `rungOpenMol` is not in `$PATH`, then nothing happens. Otherwise, gOpenMol should start and display the molecule. Any `.plt` files produced can be added to the display by following the Plot;Contour menu item. The name of the Unix command may be changed from the default `rungOpenMol` by specifying it as the first argument to the INTERACT directive. By default, gOpenMol is not started, and this is equivalent to giving the command BATCH.

26 DIABATIC ORBITALS

In order to construct diabatic states, it is necessary to determine the mixing of the diabatic states in the adiabatic wavefunctions. In principle, this mixing can be obtained by integration of the non-adiabatic coupling matrix elements. Often, it is much easier to use an approximate method, in which the mixing is determined by inspection of the CI coefficients of the MCSCF or CI wavefunctions. This method is applicable only if the orbital mixing is negligible. For CASSCF wavefunctions this can be achieved by maximizing the overlap of the active orbitals with those of a reference geometry, at which the wavefunctions are assumed to be diabatic (e.g. for symmetry reasons). The orbital overlap is maximized using using the new DIAB command in the MCSCF program.

This procedure works as follows: first, the orbitals are determined at the reference geometry. Then, the calculations are performed at displaced geometries, and the "diabatic" active orbitals, which have maximum overlap with the active orbitals at the reference geometry, are obtained by adding a DIAB directive to the input:

Old form (Molpro96, obsolete):

```
DIAB,orbref, orbsav, orb1,orb2,pri
```

New form:

```
DIAB,orbref[,TYPE=orbtype] [,STATE=state] [,SPIN=spin] [,MS2=ms2] [,SAVE=orbsav]
[,ORB1=orb1, ORB2=orb2][,PRINT=pri]
```

Here *orbref* is the record holding the orbitals of the reference geometry, and *orbsav* is the record on which the new orbitals are stored. If *orbsav* is not given (recommended!) the new orbitals are stored in the default dump record (2140.2) or the one given on the ORBITAL directive (see section 17.5.3). In contrast to earlier versions of MOLPRO it is possible that *orbref* and *orbsav* are the same. The specifications TYPE, STATE, SPIN can be used to select specific sets of reference orbitals, as described in section 2.16. *orb1*, *orb2* is a pair of orbitals for which the overlap is to be maximized. These orbitals are specified in the form *number.sym*, e.g. 3.1 means the third orbital in symmetry 1. If *orb1*, *orb2* are not given, the overlap of all active orbitals is maximized. *pri* is a print parameter. If this is set to 1, the transformation angles for each orbital are printed for each jacobi iteration.

Using the defaults described above, the following input is sufficient in most cases:

```
DIAB,orbref
```

Using Molpro98 is is not necessary any more to give any GEOM and DISPL cards. The displacements and overlap matrices are computed automatically (the geometries are stored in the dump records, along with the orbitals).

The diabatic orbitals have the property that the sum of orbital and overlap contributions in the non-adiabatic coupling matrix elements become approximately zero, such that the adiabatic mixing occurs only through changes of the CI coefficients. This allows to determine the mixing angle directly from the CI coefficients, either in a simple way as described for instance in J. Chem. Phys. **89**, 3139 (1988), or in a more advanced manner as described by Pacher, Cederbaum, and Köppel in J. Chem. Phys. **89**, 7367 (1988).

Below we present an example for the first two excited states of H₂S, which have *B*₁ and *A*₂ symmetry in *C*_{2v}, and *A*^{''} symmetry in *C*_s. We first perform a reference calculation in *C*_{2v} symmetry, and then determine the diabatic orbitals for displaced geometries in *C*_s symmetry. Each subsequent calculation uses the previous orbitals as reference. One could also use the orbitals of the *C*_{2v} calculation as reference for all other calculations. In this case one would have to take out the second-last input card, which sets `reforb=2141.2`.

```

***,H2S diabatic A" states

basis=VDZ                                !use cc-pVDZ basis set
geometry={x;                              !use Cs symmetry
          planeyz;                        !fix orientation of the molecule
          noorient                        !dont allow automatic reorientation
          s;h1,s,r1;h2,s,r2,h1,theta}    !Z-matrix geometry input

gprint,orbitals,civector                !global print options

text,reference calculation for C2V
theta=92.12,r1=2.3,r2=2.3                !reference geometry

hf;occ,7,2;wf,18,1;                      !scf calculation for ground state

multi;occ,9,2;closed,4,1;                !define active and inactive spaces
wf,18,2;state,2;                          !two A" states (1B1 and 1A2 in C2v)
orbital,2140.2                            !save orbitals to 2140.2
reorb=2140.2                               examples/
                                           h2s_diab.com

text,calculations at displaced geometries

rd=[2.4,2.5,2.6]                          !define a range of bond distances

do i=1,#rd                                !loop over displaced geometries

r2=rd(i)                                  !set r2 to current distance

multi;occ,9,2;closed,4,1;                !same wavefunction definition as at reference geom.
wf,18,2;state,2;                          !save new orbitals to record
orbital,2141.2                            !compute diabatic orbitals using reference orbitals
diab,reorb                                !stored on record reorb
reorb=2141.2                              !set variable reorb to the new orbitals.
enddo

```


27 NON ADIABATIC COUPLING MATRIX ELEMENTS

Non-adiabatic coupling matrix elements can be computed by finite differences for MCSCF or CI wavefunctions using the DDR program. For state-averaged MCSCF wavefunctions, they can also be computed analytically (cf. section 17.9.2).

Note that present numerical procedure has been much simplified relative to `Molpro96`. No `GEOM` and `DISPL` input cards are needed any more, and the three necessary calculations can be done in any order.

27.1 The DDR procedure

In order to compute the coupling matrix elements by finite differences, one has to compute and store the wavefunctions at two (first-order algorithm) or three (second-order algorithm) slightly displaced geometries. The order of these calculations is arbitrary.

The typical strategy is as follows:

- 1.) Compute the wavefunction at the reference geometry. The wavefunctions for both states have to be stored using the `SAVE` command of the `CI` program. If the matrix elements are computed for MCSCF wavefunctions, it is necessary to recompute the wavefunction with the `CI` program, using the `NOEXC` option. The transition density matrix is stored using the `DM` directive of the `CI` program.
- 2.) Compute the wavefunctions at the (positively) displaced geometry and store the `CI` wavefunction in a second record.
- 3.) If the second-order (three-point) method is used, step (2) is repeated at a (negatively) displaced geometry.
- 4.) Compute the transition density matrices between the states at the reference geometry and the displaced geometries. This is done with the `TRANS` directive of the `CI` program.
- 5.) Finally, the `DDR` program is used to assemble the matrix element. Using the first-order two-point method, only a single input line is needed:

```
DDR, dr, orb1, orb2, trdm2
```

where *dr* is the geometry increment used as denominator in the finite difference method, *orb1* is the record holding the orbitals of the reference geometry, *orb2* is the record holding the orbitals of the displaced geometry, and *trdm2* is the record holding the transition density matrix computed from the `CI`-vectors at *R* and *R+DR*.

If central differences (three points) are used, the input is as follows:

```
DDR,2*dr
ORBITAL,orb1,orb2,orb3
DENSITY,trdm1,trdm2,trdm3
```

where *dr*, *orb1*, *orb2* are as above, and *orb3* is the record holding the orbitals at the negatively displaced geometry.

trdm1, *trdm2*, *trdm3* are the records holding the transition densities $\gamma(R|R)$, $\gamma(R|R + DR)$, and $\gamma(R|R - DR)$, respectively.

If more than two states are computed simultaneously, the transition density matrices for all pairs of states will be stored in the same record. In that case, and also when there are just two states

whose spatial symmetry is not 1, it is necessary to specify for which states the coupling is to be computed using the STATE directive:

```
STATE,state1, state2
```

where *state*_{*i*} is of the form *istate.isym* (the symmetries of both states must be the same, and it is therefore sufficient to specify the symmetry of the first state).

As an example the input for first-order and second-order calculations is given below. The calculation is repeated for a range of geometries, and at the end of the calculation the results are printed using the TABLE command.

In the calculation shown, the "diabatic" CASSCF orbitals are generated in the two CASSCF calculations at the displaced geometries by maximizing the overlap with the orbitals at the reference geometry. This is optional, and (within the numerical accuracy) does not influence the final results. However, the relative contributions of the orbital, overlap and CI contributions to the NACME are modified. If diabatic orbitals are used, which change as little as possible as function of geometry, the sum of overlap and orbital contribution is minimized, and to a very good approximation the NACME could be obtained from the CI-vectors alone.

```

***,lif non-adiabatic coupling
memory,1,m

basis,f=avdz,li=vdz           !define basis
r=[10.0,10.5,11.0,11.5,12.0]  !define bond distances
dr=0.01                        !define increment
geometry={li:f,li,rlif}       !define geometry

rlif=3                          !first calculation at R=3
hf;occ,4,1,1                   !SCF
multi;closed,3;                !CASSCF, 3 inactive orbitals
wf,12,1;state,2;              !Two 1A1 states
orbital,2140.2                 !dump orbitals to record 2140.2

do i=1,#r                       !loop over geometries
  rlif=r(i)                     !set bond distance
  multi;closed,3;               !CASSCF, 3 inactive orbitals
  wf,12,1;state,2;             !Two 1A1 states
  orbital,2140.2                !Overwrite previous orbitals by present ones

  ci;state,2;noexc;            !CI for 2 states, no excitations
  save,6000.2;                  !save wavefunction to record 6000.2
  dm,8000.2;                    !save (transition) densities to record 8000.2

  rlif=r(i)+dr                 !increment bond distance by dr

  multi;closed,3;              !same CASSCF as above
  wf,12,1;state,2;            !Two 1A1 states
  start,2140.2;                !start with orbitals from reference geometry
  orbital,2141.2;              !save orbitals to record 2141.2
  diab,2140.2                  !generate diabatic orbitals by maximizing the
                               !overlap with the orbitals at the reference geometry

  ci;state,2;noexc;save,6001.2; !CI for 2 states, wavefunction saved to record 6001.2
                               examples/
                               lif_nacme.com

  rlif=r(i)-dr                 !repeat at r-dr

  multi;closed,3;              !same CASSCF as above
  wf,12,1;state,2;            !Two 1A1 states
  start,2140.2;                !start with orbitals from reference geometry
  orbital,2142.2;              !save orbitals to record 2142.2
  diab,2140.2                  !generate diabatic orbitals by maximizing the
                               !overlap with the orbitals at the reference geometry

  ci;state,2;noexc;save,6002.2; !CI for 2 states, wavefunction saved to record 6002.2

  ci;trans,6000.2,6001.2;      !Compute overlap and transition density <R|R+DR>
  dm,8100.2;                    !Save transition density to record 8100.2

  ci;trans,6000.2,6002.2;      !Compute overlap and transition density <R|R-DR>
  dm,8200.2;                    !Save transition density to record 8200.2

  ddr,dr,2140.2,2141.2,8100.2  !compute NACME using 2-point formula (forward difference)
  nacmelp(i)=nacme             !store result in variable nacmelp
  ddr,-dr,2140.2,2142.2,8200.2 !compute NACME using 2-point formula (backward difference)
  nacmelm(i)=nacme            !store result in variable nacmelm

  ddr,2*dr                      !compute NACME using 3-point formula
  orbital,2140.2,2141.2,2142.2; !orbital records for R, R+DR, R-DR
  density,8000.2,8100.2,8200.2; !transition density records for R, R+DR, R-DR
  nacme2(i)=nacme              !store result in variable nacme2

end do                          !end of loop over differend bond distances

nacmeav=(nacmelp+nacmelm)*0.5  !average the two results forward and backward differences
table,r,nacmelp,nacmelm,nacmeav,nacme2 !print a table with results
title,Non-adiabatic couplings for LiF !title for table

```

This calculation produces the following table:

Non-adiabatic couplings for LiF

R	NACME1P	NACME1M	NACMEAV	NACME2
10.0	-0.22828936	-0.22328949	-0.22578942	-0.22578942
10.5	-0.51777034	-0.50728914	-0.51252974	-0.51252974
11.0	0.76672943	0.76125391	0.76399167	0.76399167
11.5	0.42565202	0.42750263	0.42657733	0.42657733
12.0	0.19199878	0.19246799	0.19223338	0.19223338

Note that the sign changes because of a phase change of one of the wavefunctions. In order to keep track of the sign, one has to inspect both the orbitals and the ci-vectors.

28 QUASI-DIABATIZATION

The DDR procedure can also be used to generate quasi-diabatic states and energies for MRCI wavefunctions (CASSCF case can be treated as special case using the NOEXC directive in the MRCI). The quasi-diabatic states have the property that they change as little as possible relative to a reference geometry; with other words, the overlap between the states at the current geometry with those at a reference geometry is maximized by performing a unitary transformation among the given states. Preferably, the adiabatic and diabatic states should be identical at the reference geometry, e.g., due to symmetry. For instance, in the examples given below for the 1B_1 and 1A_2 states of H_2S , C_{2v} geometries are used as reference, and at these geometries the states are unmixed due to their different symmetry. At the displaced geometries the molecular symmetry is reduced to C_s . Both states now belong to the $^1A''$ irreducible representation and are strongly mixed. For a description and application of the procedure described below, see D. Simah, B. Hartke, and H.-J. Werner, J. Chem. Phys. **111**, 4523 (1999).

This diabaticization can be done automatically and requires two steps: first, the active orbitals of a CASSCF calculation are rotated to maximize the overlap with the orbitals at the reference geometry. This is achieved using the DIAB procedure described in section 17.5.8. Secondly, the DDR procedure can be used to find the transformation among the CI vectors.

The following input is required:

DDR calls the DDR procedure.

ORBITAL,*orb1*, *orb2* *orb1* and *orb2* are the (diabatic) orbitals at the current and reference geometry, respectively.

DENSITY,*trdm1*,*trdm2* *trdm1* are the transition densities computed at the current geometry, *trdm2* are transition densities computed using the wavefunctions of the current (bra) and reference (ket) geometries.

MIXING,*state1*, *state2*, ... The given states are included in the diabaticization.

ENERGY,*e1*, *e2*, ... Adiabatic energies of the states. If this input card is present, the Hamiltonian in the basis of the diabatic states is computed and printed. Alternatively, the energies can be passed to DDR using the Molpro variable EADIA.

The results are printed and stored in the following Molpro variables, provided the ENERGY directive or the EADIA variable is found:

Results including the first-order orbital correction:

SMAT	The first $nstate \times nstate$ elements contain the state overlap matrix (bra index runs fastest).
UMAT	The first $nstate \times nstate$ elements contain the transformation matrix.
HDIA	The first $nstate \cdot (nstate + 1)/2$ elements contain the lower triangle of the diabatic hamiltonian.
MIXANG	Non-adiabatic mixing angle in degree. This is available only in the two-state case.

The corresponding results obtained from the CI-vectors only (without orbital correction) are stored in the variables [SMATCI], UMATCI, HDIACI, and MIXANGCI.

The way it works is most easily demonstrated for some examples. In the following input, the wavefunction is first computed at the C_{2v} reference geometry, and then at displaced geometries.

```

***,h2s Diabatization
memory,3,m

gprint,orbitals,civector

geometry{x;noorient      !noorient should always be used for diabatization
        s;
        h1,s,r1;
        h2,s,r2,h1,theta}

basis=avdz              !This basis is too small for real application

r1=2.5                  !Reference geometry
theta=[92]

r=[2.50,2.55,2.60]    !Displaced geometries

reorb=2140.2           !Orbital dumprecord at reference geometry
refci=6000.2           !MRCI record at reference geometry
savci=6100.2           !MRCI record at displaced geometries

text,compute wavefunction at reference geometry (C2v)
r2=r1

hf;occ,9,2;wf,18,2,4;
orbital,2100.2

multi;occ,9,2;closed,4,1;
wf,18,2;state,2;      !1B1 and 1A2 states
natorb,reorb          !Save reference orbitals on reorb
noextra               !Dont use extra symmetries

ci;occ,9,2;closed,4,1; !MRCI at reference geometry
wf,18,2,0;state,2;   !1B1 and 1A2 states
orbital,reorb         !Use orbitals from previous CASSCF
save,refci;           !Save MRCI wavefunction

Text,Displaced geometries

do i=1,#r              !Loop over different r values
data,truncate,savci+1 !truncate dumpfile after reference
r2=r(i)                !Set current r2

multi;occ,9,2;closed,4,1;
wf,18,2,0;state,2;   !Wavefunction definition
start,reorb           !Starting orbitals
orbital,3140.2;       !Dump record for orbitals
diab,reorb            !Generate diabatic orbitals relative to reference geometry
noextra               !Dont use extra symmetries

ci;occ,9,2;closed,4,1;
wf,18,2,0;state,2;   !1B1 and 1A2 states
orbital,diabatic      !Use diabatic orbitals
save,savci;           !Save MRCI for displaced geometries

e1(i)=energy(1)       !Save adiabatic energies
e2(i)=energy(2)

ci;trans,savci,savci  !Compute transition densities at R2
dm,7000.2;            !Save transition densities on this record
ci;trans,savci,refci; !Compute transition densities between R2 and R1
dm,7100.2;            !Save transition densities on this record

ddr
density,7000.2,7100.2 !Densities for <R2||R2> and <R2||R1>
orbital,3140.2,2140.2 !Orbitals for <R2||R2> and <R2||R1>
energy,e1(i),e2(i)    !Adiabatic energies
mixing,1.2,2.2        !Compute mixing angle and diabatic energies

mixci(i)=mixangci(1)  !Mixing angle obtained from ci vectors only
hllci(i)=hdiaci(1)   !Diabatic energies obtained from ci vectors only

```

This calculation produces the following results:

Diabatic energies for H2S, obtained from CI-vectors

R	E1	E2	H11CI	H22CI	H21CI	MIXCI
2.50	-398.64296319	-398.63384782	-398.64296319	-398.63384782	0.00000000	0.00
2.55	-398.64572746	-398.63666636	-398.64509901	-398.63729481	-0.00230207	15.27
2.60	-398.64911752	-398.63771802	-398.64662578	-398.64020976	-0.00471125	27.87

Diabatic energies for H2S, obtained from CI-vectors and orbital correction

R	E1	E2	H11	H22	H21	MIXTOT
2.50	-398.64296319	-398.63384782	-398.64296319	-398.63384782	0.00000000	0.00
2.55	-398.64572746	-398.63666636	-398.64509941	-398.63729441	-0.00230139	15.26
2.60	-398.64911752	-398.63771802	-398.64662526	-398.64021027	-0.00471160	27.88

The results in the first table are obtained from the CI-contribution to the state-overlap matrix only, while the ones in the second table include a first-order correction for the orbitals. In this case, both results are almost identical, since the DIAB procedure has been used to minimize the change of the active orbitals. This is the recommended procedure. If simply natural orbitals are used without orbital diabaticization, the following results are obtained from the otherwise unchanged calculation:

Diabatic energies for H2S, obtained from CI-vectors

R	E1	E2	H11CI	H22CI	H21CI	MIXCI
2.50	-398.64296319	-398.63384782	-398.64296319	-398.63384782	0.00000000	0.00
2.55	-398.64572742	-398.63666630	-398.64475612	-398.63763760	-0.00280315	19.11
2.60	-398.64911746	-398.63771803	-398.64521031	-398.64162518	-0.00541050	35.83

Diabatic energies for H2S, obtained from CI-vectors and orbital correction

R	E1	E2	H11	H22	H21	MIXTOT
2.50	-398.64296319	-398.63384782	-398.64296319	-398.63384782	0.00000000	0.00
2.55	-398.64572742	-398.63666630	-398.64509146	-398.63730226	-0.00231474	15.36
2.60	-398.64911746	-398.63771803	-398.64648358	-398.64035190	-0.00480493	28.73

It is seen that the mixing obtained from the CI vectors only is now very different and meaningless, since the orbitals change significantly as function of geometry. However, the second calculations, which accounts for this change approximately, still gives results in quite good agreement with the calculation involving diabatic orbitals.

The final examples shows a more complicated input, which also computes the non-adiabatic coupling matrix elements. In a two-state model, the NACME should equal the first derivative of the mixing angle. In the example, the NACME is computed using the 3-point DDR method (NACMECI), and also by finite difference of the mixing angle (DCHI).

```

***,h2s Diabatization and NACME calculation
memory,3,m

```

```

gprint,orbitals,civector

```

```

geometry{x;noorient      !noorient should always be used for diabatization
        s;
        h1,s,r1;
        h2,s,r2,h1,theta}

```

```

basis=avdz              !This basis is too small for real application

```

```

r1=2.5                  !Reference geometry
theta=[92]

```

```

r=[2.55,2.60]          !Displaced geometries
dr=[0,0.01,-0.01]     !Small displacements for finite difference NACME calculation

```

```

reforb1=2140.2         !Orbital dumprecord at reference geometry
refci=6000.2          !MRCI record at reference geometry
savci=6100.2          !MRCI record at displaced geometries

```

```

text,compute wavefunction at reference geometry (C2v)
r2=r1

```

```

hf;occ,9,2;wf,18,2,4;orbital,2100.2

```

```

multi;occ,9,2;closed,4,1;
wf,18,2;state,2;      !1B1 and 1A2 states
natorb,reforb1        !Save reference orbitals on reforb1
noextra               !Dont use extra symmetries

```

```

ci;occ,9,2;closed,4,1; !MRCI at reference geometry
wf,18,2,0;state,2;    !1B1 and 1A2 states
orbital,reforb1       !Use orbitals from previous CASSCF
save,refci;           !Save MRCI wavefunction

```

```

Text,Displaced geometries

```

```

do i=1,#r              !Loop over different r values
data,truncate,savci+1 !truncate dumpfile after reference
reforb=reforb1

```

```

do j=1,3              !Loop over small displacements for NACME
r2=r(i)+dr(j)        !Set current r2

```

```

multi;occ,9,2;closed,4,1;
wf,18,2,0;state,2;   !Wavefunction definition
start,reforb         !Starting orbitals
orbital,3140.2+j;    !Dumprecord for orbitals
diab,reforb         !Generate diabatic orbitals relative to reference geometry
noextra             !Dont use extra symmetries
reforb=3141.2       !Use orbitals for j=1 as reference for j=2,3

```

```

ci;occ,9,2;closed,4,1;
wf,18,2,0;state,2;   !Use diabatic orbitals
orbital,diabatic     !Use diabatic orbitals
save,savci+j;        !Save MRCI for displaced geometries

```

```

eadia=energy          !Save adiabatic energies for use in ddr
if(j.eq.1) then
e1(i)=energy(1)      !Save adiabatic energies for table printing
e2(i)=energy(2)
end if

```

```

ci;trans,savci+j,savci+j; !Compute transition densities at R2+DR(j)
dm,7000.2+j;           !Save transition densities on this record
ci;trans,savci+j,refci; !Compute transition densities between R2+DR(j) and R1
dm,7100.2+j;           !Save transition densities on this record
ci;trans,savci+j,savci+1; !Compute transition densities between R and R2+DR(j)
dm,7200.2+i;           !Save transition densities on this record

```


The calculation produces the following table

Mixing angles and non-adiabatic coupling matrix elements for H2S

R	MIXCI	MIXTOT	DCHI	NACMECI
2.55	15.2694	15.2644	-5.2226	-5.2365
2.60	27.8740	27.8772	-3.4702	-3.4794

Diabatic energies for H2S, obtained from CI-vectors

R	E1	E2	H11CI	H22CI	H21CI
2.55	-398.64572746	-398.63666636	-398.64509901	-398.63729481	-0.00230207
2.60	-398.64911752	-398.63771802	-398.64662578	-398.64020976	-0.00471125

Diabatic energies for H2S, obtained from CI-vectors and orbital correction

R	E1	E2	H11	H22	H21
2.55	-398.64572746	-398.63666636	-398.64509941	-398.63729441	-0.00230139
2.60	-398.64911752	-398.63771802	-398.64662526	-398.64021027	-0.00471160

As expected the coupling matrix elements obtained from the 3-point DDR calculation (NACMECI) and by differentiating the mixing angle (DCHI) are in close agreement.

29 THE VB PROGRAM CASVB

CASVB is a general program for valence bond calculations written by T. Thorsteinsson and D. L. Cooper (1996–2000).

This program can be used in two basic modes:

- a) variational optimization of quite general types of nonorthogonal MCSCF or modern valence bond wavefunctions
- b) representation of CASSCF wavefunctions in modern valence form, using overlap- (*relatively inexpensive*) or energy-based criteria.

Bibliography:

- T. Thorsteinsson, D. L. Cooper, J. Gerratt, P. B. Karadakov and M. Raimondi, *Theor. Chim. Acta* **93**, 343–366 (1996).
D. L. Cooper, T. Thorsteinsson and J. Gerratt, *Int. J. Quant. Chem.* **65**, 439–51 (1997).
D. L. Cooper, T. Thorsteinsson and J. Gerratt, *Adv. Quant. Chem.* **32**, 51–67 (1998).
T. Thorsteinsson and D. L. Cooper, *Prog. Theor. Chem. Phys.* (in press).

All publications resulting from use of this program must acknowledge some or all of the above. For an up-to-date bibliography see <http://rs2.ch.liv.ac.uk/dlc/CASVB.html>

29.1 Structure of the input

All CASVB sub-commands may be abbreviated by four letters. The general input structure can be summarized as follows:

- a) For generating representations of CASSCF wavefunctions, the program is invoked by the command CASVB. For variational optimization of wavefunctions it is normally invoked inside *MULTI* by the sub-command VB (see 17.10).
- b) Definition of the CASSCF wavefunction (not generally required).
- c) Definition of the valence bond wavefunction.
- d) Recovery and/or storage of orbitals and vectors.
- e) Manual input of starting guess (optional).
- g) Optimization control.
- f) Definition of molecular symmetry and possible constraints on the VB wavefunction.
- h) Wavefunction analysis.
- i) Further general options.

Items a) and b) should precede everything else in the input; apart from this, commands may come in any order.

29.2 Defining the CASSCF wavefunction

CASVB is interfaced with the determinant part of *MULTI* (i.e., *CONFIG,CSF*; must *not* be specified). When this program is run prior to *CASVB*, the CI vector must be dumped using one of the directives *SAVE*, *NATORB*, *CANONICAL*, or *LOCALI* (see section 17.5.4). The three latter are recommended.

29.2.1 The VBDUMP directive

VBDUMP[,*vbdump*];

It is advisable to restore the wavefunction definitions using VBDUMP cards here and in the CASSCF calculation (see Section 17.8.6). The default record name (*vbdump*) is 4299.2. If a VBDUMP card is not present and record 4299.2 does not exist, then *CASVB* will attempt to generate the wavefunction information automatically based on the latest MCSCF calculation (however, *STATE* and *WEIGHT* information will not be restored in such a case).

If present, the VBDUMP card must occur first in the *CASVB* input. It is not required for variational calculations.

Note that in the majority of cases (e.g., if a *CASVB* run occurs immediately after *MULTI*, or for variational calculations), explicit specification of dump records with *vbdump* is not required.

29.3 Other wavefunction directives

The definitions of the CASSCF wavefunction may also be specified manually using some or all of the directives:

OCC	Occupied orbitals.
CLOSED	Closed-shell orbitals.
CORE	Frozen-core orbitals.
WF	Wavefunction card.
STATE	Number of states for this wavefunction symmetry.
WEIGHT	Weights of states.

For the exact definition of these cards see sections 17.2 and 17.3. These commands may also be used to modify the values defined in VBDUMP. The information given on these cards should correspond to the CI vector saved in the CASSCF calculation. The cards, and their ordering, should therefore coincide with those used in *MULTI*, except for the *WEIGHT* cards which may differ. At present, the VB wavefunction must correspond to a well-defined number of electrons and total spin. Other states may be present, but an error condition will occur if non-zero weights are specified for wavefunction symmetries with varying values of *elec* or *spin*.

29.4 Defining the valence bond wavefunction

29.4.1 Specifying orbital configurations

The number of core and active orbitals (*mcore*, *mact*), active electrons (*Nact*), and the value of the total spin will be identical to that defined for the CASSCF wavefunction. The spatial VB

configurations are defined in terms of the active orbitals only, and may be specified using one or more CON cards (note that the RESTRICT and SELECT keywords are not used in CASVB):

CON, $n_1, n_2, n_3, n_4, \dots$;

The configurations can be specified by occupation numbers, as in section 13.4.3, so that n_i is the occupation of the i th valence bond orbital. Alternatively a list of N_{act} orbital numbers (in any order) may be provided – the program determines which definition applies. The two cards CON, 1, 0, 1, 2; and CON, 1, 3, 4, 4; are thus equivalent.

If no configurations are specified the single covalent configuration $\phi_1\phi_2 \cdots \phi_{N_{act}}$ is assumed.

29.4.2 Selecting the spin basis

SPINBASIS,*key*;

key may be chosen from KOTANI (default), RUMER, PROJECT or LTRUMER, specifying the basis of spin eigenfunctions used in the definition of valence bond structures. PROJECT refers to spin functions generated using a spin projection operator, LTRUMER to Rumer functions with the so-called “leading term” phase convention.

29.5 Recovering CASSCF CI vector and VB wavefunction

The appropriate MOLPRO records may be specified explicitly using the START directive (note, however, that use of the *vbdump* mechanism described in section 29.2.1 is preferable whenever possible):

START,*ci, vb, orb, trnint*;

ci: record name for the CASSCF CI vector. The CI vector must have been dumped previously using either of the SAVE, NATORB, CANONICAL, or LOCALI directives (see section 17.5.4). A default value for *ci* is determined from the most recent *vbdump* record(s).

Note that if the *ci* record is not found, only an energy-based optimization of the VB wavefunction can be carried out.

vb: record name for the valence bond orbitals and structure coefficients, as saved by a previous CASVB calculation. If the VB wavefunction was previously saved in the AO basis the orbitals will be projected onto the present active space (note that it is necessary to specify a record name for the molecular orbitals (*orb* below) for this to be possible).

orb: record name for the molecular orbitals defining the CASSCF wavefunction. This information is necessary if one wants to output the valence bond orbitals in the atomic orbital basis.

trnint: record name for the transformed CASSCF integrals. These are required for the energy-based criteria (i.e., if CRIT, ENERGY is specified), and can be saved inside MULTI by the TRNINT sub-command (see 17.8.7). The default record name, both here and in MULTI, is 1900.1.

29.6 Saving the VB wavefunction

SAVE,*vb, civb, vbao*;

vb: record name for VB wavefunction (default is first available record after 3200.2), i.e., orbitals and structure coefficients.

civb: record name for valence bond full CI vector defined in terms of the CASSCF MOs (default is 3300.2). Saving this vector is necessary for the calculation of further properties, geometry optimization, etc.

vba0: record name for valence bond wavefunction in the AO basis. Note that specifying *orb* in the START directive is a precondition for this keyword. It may be useful for plotting of orbitals, or for providing a guess to be used in the interpretation of a CASSCF solution employing a different active space.

29.7 Specifying a guess

GUESS={*key-1*,...;*key-2*,...;...}

The GUESS keyword initiates the input of a guess for the valence bond orbitals and structure coefficients. *key-i* can be either ORB, STRUC or READ. These keywords modify the guess provided by the program, or specified by the START directive. It is thus possible to modify individual orbitals in a previous solution to construct the starting guess.

29.7.1 Orbital guess

ORB,*i*, *c*₁, *c*₂,... *c*_{*mact*};

Specifies a starting guess for valence bond orbital number *i*. The guess is specified in terms of the *mact* active MOs defining the CASSCF wavefunction. (Note that the definition of these MOs will depend on how the CI vector was dumped – i.e. which of the SAVE, NATORB, CANONICAL, or LOCALI directives was used (see section 17.5.4). Use of one of the three latter keywords is recommended.)

29.7.2 Guess for structure coefficients

STRUC,*c*₁, *c*₂,... *c*_{*NVB*};

Specifies a starting guess for the *NVB* structure coefficients. If this card is not provided, and no guess specified by START, the perfect-pairing mode of spin coupling is assumed for the spatial configuration having the least number of doubly occupied orbitals. Note that the definition of structures depends on the value of SPINBASIS. Doubly occupied orbitals occur first in all configurations, and the spin eigenfunctions are based on the singly occupied orbitals being in ascending order.

29.7.3 Read orbitals or structure coefficients

The READ keyword can take one of the following forms:

READ,ORB,*iorb1*[,TO,*iorb2*] [,AS,*jorb1*[,TO,*jorb2*]] [,FROM,*record*];

READ,STRUC,*istruc1*[,TO,*istruc2*] [,AS,*jstruc1*[,TO,*jstruc2*]] [,FROM,*record*];

READ,ALL [,FROM,*record*];

In this way a subset of orbitals and/or structure coefficients may be picked out from a previous calculation. Renumbering of orbitals or structures can be done using the “AS” construct as outlined above. If the VB wavefunction was previously saved in the AO basis, the orbitals will

be projected onto the present active space (note that it is necessary to specify a record name for the molecular orbitals (*orb* in the START command) for this to be possible).

Default for *record* is the *vb* record name specified in keyword START (if applicable).

29.8 Permuting orbitals

ORBPERM,*i*₁,...,*i*_{mact};

Permutes the orbitals in the valence bond wavefunction and changes their phases according to $\phi'_j = \text{sign}(i_j)\phi_{\text{abs}(i_j)}$. The guess may be further modified using the GUESS keyword. Also the structure coefficients will be transformed according to the given permutation (note that the configuration list must be closed under the orbital permutation for this to be possible).

29.9 Optimization control

29.9.1 Optimization criterion

CRIT,*method*;

Specifies the criterion for the optimization. *method* can be OVERLAP or ENERGY (OVERLAP is default). The former maximizes the normalized overlap with the CASSCF wavefunction:

$$\max \left(\frac{\langle \Psi_{\text{CAS}} | \Psi_{\text{VB}} \rangle}{(\langle \Psi_{\text{VB}} | \Psi_{\text{VB}} \rangle)^{1/2}} \right)$$

and the latter simply minimizes the energy:

$$\min \left(\frac{\langle \Psi_{\text{VB}} | \hat{H} | \Psi_{\text{VB}} \rangle}{\langle \Psi_{\text{VB}} | \Psi_{\text{VB}} \rangle} \right).$$

29.9.2 Number of iterations

MAXITER,*N*_{iter};

Specifies the maximum number of iterations in the second order optimizations. Default is *N*_{iter}=50.

29.9.3 CASSCF-projected structure coefficients

(NO)CASPROJ;

With this keyword the structure coefficients are picked from the transformed CASSCF CI vector, leaving only the orbital variational parameters. For further details see the bibliography. This option may be useful to aid convergence.

29.9.4 Saddle-point optimization

SADDLE,*n*;

Defines optimization onto an *n*th-order saddle point. See also T. Thorsteinsson and D. L. Cooper, Int. J. Quant. Chem. **70**, 637–50 (1998).

29.9.5 Defining several optimizations

More than one optimization may be performed in the same *CASVB* deck, by the use of `OPTIM` keywords:

```
OPTIM[={... }];
```

The subcommands may be any optimization declarations defined in this section, as well as any symmetry or constraints specifications described in section 29.10. Commands given as arguments to `OPTIM` will be particular to this optimization step, whereas commands specified outside will act as default definitions for all subsequent `OPTIM` keywords.

If only one optimization step is required, the `OPTIM` keyword need not be specified.

When only a machine-generated guess is available, *CASVB* will attempt to define a sequence of optimization steps chosen such as to maximize the likelihood of successful convergence and to minimize CPU usage. To override this behaviour, simply specify one or more `OPTIM` cards.

29.9.6 Multi-step optimization

A loop over two or more optimization steps may be specified using:

```
ALTERN,Niter={... }
```

With this specification the program will repeat the enclosed optimization steps until either all optimizations have converged, or the maximum iteration count, *Niter*, has been reached.

29.10 Point group symmetry and constraints

The problems associated with symmetry-adapting valence bond wavefunctions are considered, for example, in: T. Thorsteinsson, D. L. Cooper, J. Gerratt and M. Raimondi, *Theor. Chim. Acta* **95**, 131 (1997).

29.10.1 Symmetry operations

```
SYMELM,label,sign;
```

Initiates the definition of a symmetry operation referred to by *label* (any three characters). *sign* can be + or -; it specifies whether the total wavefunction is symmetric or antisymmetric under this operation, respectively. A value for *sign* is not always necessary but, if provided, constraints will be put on the structure coefficients to ensure that the wavefunction has the correct overall symmetry (note that the configuration list must be closed under the orbital permutation induced by *label* for this to be possible).

The operator is defined in terms of its action on the active MOs as specified by one or more of the keywords `IRREPS`, `COEFFS`, or `TRANS` (any other keyword will terminate the definition of this symmetry operator). If no further keyword is supplied, the identity is assumed for *label*. The alternative format `SYMELM,label,sign={key-1,...;key-2,...;... }` may also be used.

29.10.2 The IRREPS keyword

```
IRREPS,i1, i2,...;
```

The list i_1, i_2, \dots specifies which irreducible representations (as defined in the CASSCF wavefunction) are antisymmetric with respect to the *label* operation. If an irreducible representation is not otherwise specified it is assumed to be symmetric under the symmetry operation.

29.10.3 The COEFFS keyword

COEFFS, i_1, i_2, \dots ;

The list i_1, i_2, \dots specifies which individual CASSCF MOs are antisymmetric with respect to the *label* operation. If an MO is not otherwise specified, it is assumed to be symmetric under the symmetry operation. This specification may be useful if, for example, the molecule possesses symmetry higher than that exploited in the CASSCF calculation.

29.10.4 The TRANS keyword

TRANS, $n_{dim}, i_1, \dots, i_{n_{dim}}, c_{11}, c_{12}, \dots, c_{n_{dim}n_{dim}}$;

Specifies a general $n_{dim} \times n_{dim}$ transformation involving the MOs $i_1, \dots, i_{n_{dim}}$, specified by the *c* coefficients. This may be useful for systems with a two- or three-dimensional irreducible representation, or if localized orbitals define the CASSCF wavefunction. Note that the specified transformation must always be orthogonal.

29.10.5 Symmetry relations between orbitals

In general, for a VB wavefunction to be symmetry-pure, the orbitals must form a representation (not necessarily irreducible) of the symmetry group. Relations between orbitals under the symmetry operations defined by SYMELM may be specified according to:

ORBREL, $i_1, i_2, label1, label2, \dots$;

Orbital i_1 is related to orbital i_2 by the sequence of operations defined by the *label* specifications (defined previously using SYMELM). The operators operate right to left. Note that i_1 and i_2 may coincide. Only the minimum number of relations required to define all the orbitals should be provided; an error exit will occur if redundant ORBREL specifications are found.

29.10.6 The SYMPROJ keyword

As an alternative to incorporating constraints, one may also ensure correct symmetry of the wavefunction by use of a projection operator:

(NO)SYMPROJ[,*irrep1, irrep2, \dots*];

The effect of this keyword is to set to zero coefficients in unwanted irreducible representations. For this purpose the symmetry group defined for the CASSCF wavefunction is used (always a subgroup of D_{2h}). The list of irreps in the command specifies which components of the wavefunction should be kept. If no irreducible representations are given, the current wavefunction symmetry is assumed. In a state-averaged calculation, all irreps are retained for which a non-zero weight has been specified in the wavefunction definition. The SYMPROJ keyword may also be used in combination with constraints.

29.10.7 Freezing orbitals in the optimization

FIXORB, i_1, i_2, \dots ;

This command freezes the orbitals specified in the list i_1, i_2, \dots to that of the starting guess. Alternatively the special keywords ALL or NONE may be used. These orbitals are eliminated from the optimization procedure, but will still be normalized and symmetry-adapted according to any ORBREL keywords given.

29.10.8 Freezing structure coefficients in the optimization

FIXSTRUC, i_1, i_2, \dots ;

Freezes the coefficients for structures i_1, i_2, \dots . Alternatively the special keywords ALL or NONE may be used. The structures are eliminated from the optimization procedure, but may still be affected by normalization or any symmetry keywords present.

29.10.9 Deleting structures from the optimization

DELSTRUC, i_1, i_2, \dots , [ALL], [NONE];

Deletes the specified structures from the wavefunction. The special keywords ALL or NONE may be used. A structure coefficient may already be zero by symmetry (as defined by SYMELM and ORBREL), in which case deleting it has no effect.

29.10.10 Orthogonality constraints

ORTHCON= { $key-1, \dots; key-2, \dots; \dots$ }

The ORTHCON keyword initiates the input of orthogonality constraints between pairs of valence bond orbitals. The sub-keywords $key-i$ can be one of ORTH, PAIRS, GROUP, STRONG or FULL as described below. Orthogonality constraints should be used with discretion. Note that orthogonality constraints for an orbital generated from another by symmetry operations (using the ORBREL keyword) cannot in general be satisfied.

ORTH, i_1, i_2, \dots ;

Specifies a list of orbitals to be orthogonalized. All overlaps between pairs of orbitals in the list are set to zero.

PAIRS, i_1, i_2, \dots ;

Specifies a simple list of orthogonalization pairs. Orbital i_1 is made orthogonal to i_2, i_3 to i_4 , etc.

GROUP, $label, i_1, i_2, \dots$;

Defines an orbital group to be used with the ORTH or PAIRS keyword. The group is referred to by $label$ which can be any three characters beginning with a letter a–z. Labels defining different groups can be used together or in combination with orbital numbers in ORTH or PAIRS. i_1, i_2, \dots specifies the list of orbitals in the group. Thus the combination GROUP,A,1,2; GROUP,B,3,4; ORTH,A,B; will orthogonalize the pairs of orbitals 1-3, 1-4, 2-3 and 2-4.

STRONG;

This keyword is short-hand for strong orthogonality. The only allowed non-zero overlaps are between pairs of orbitals $(2n-1, 2n)$.

FULL;

This keyword is short-hand for full orthogonality. This is mainly likely to be useful for testing purposes.

29.11 Wavefunction analysis

29.11.1 Spin correlation analysis

(NO)SCORR;

With this option, expectation values of the spin operators $(\hat{s}_\mu + \hat{s}_\nu)^2$ are evaluated for all pairs of μ and ν . Default is NOSCORR. The procedure is described by: G. Raos, J. Gerratt, D. L. Cooper and M. Raimondi, Chem. Phys. **186**, 233–250 (1994); *ibid*, 251–273 (1994); D. L. Cooper, R. Ponec, T. Thorsteinsson and G. Raos, Int. J. Quant. Chem. **57**, 501–518 (1996).

At present this analysis is only implemented for spin-coupled wavefunctions.

29.11.2 Printing weights of the valence bond structures

For further details regarding the calculation of weights in CASVB, see T. Thorsteinsson and D. L. Cooper, J. Math. Chem. **23**, 105-26 (1998).

VBWEIGHTS,*key1,key2,...*

Calculates and outputs weights of the structures in the valence bond wavefunction Ψ_{VB} . *key* specifies the definition of nonorthogonal weights to be used, and can be one of:

CHIRGWIN	Evaluates Chirgwin-Coulson weights (see: B. H. Chirgwin and C. A. Coulson, Proc. Roy. Soc. Lond. A201 , 196 (1950)).
LOWDIN	Performs a symmetric orthogonalization of the structures and outputs the corresponding weights.
INVERSE	Outputs "inverse overlap populations" as in G. A. Gallup and J. M. Norbeck, Chem. Phys. Lett. 21 , 495–500 (1973).
ALL	All of the above.
NONE	Suspends calculation of structure weights.

The commands LOWDIN and INVERSE require the overlap matrix between valence bond structures, and some computational overhead is thus involved.

29.11.3 Printing weights of the CASSCF wavefunction in the VB basis

For further details regarding the calculation of weights in CASVB, see T. Thorsteinsson and D. L. Cooper, J. Math. Chem. **23**, 105-26 (1998).

CIWEIGHTS,*key1,key2,...* [,*N_{conf}*];

Prints weights of the CASSCF wavefunction transformed to the basis of nonorthogonal VB structures. For the *key* options see VBWEIGHTS above. Note that the evaluation of inverse overlap weights involves an extensive computational overhead for large active spaces. Weights are

given for the total CASSCF wavefunction, as well as the orthogonal complement to Ψ_{VB} . The default for the number of configurations requested, N_{conf} , is 10. If $N_{\text{conf}}=-1$ all configurations are included.

29.12 Controlling the amount of output

PRINT, i_1, i_2, \dots ;

Each number specifies the level of output required at various stages of the execution, according to the following convention:

-1	No output except serious, or fatal, error messages.
0	Minimal output.
1	Standard level of output.
2	Extra output.

The areas for which output can be controlled are:

i_1	Print of input parameters, wavefunction definitions, etc.
i_2	Print of information associated with symmetry constraints.
i_3	General convergence progress.
i_4	Progress of the 2nd order optimization procedure.
i_5	Print of converged solution and analysis.
i_6	Progress of variational optimization.
i_7	Usage of record numbers on file 2.

For all, the default output level is +1. If $i_5 \geq 2$ VB orbitals will be printed in the AO basis (provided that the definition of MOs is available).

29.13 Service mode

SERVICE;

This keyword takes precedence over any others previously defined to CASVB. It provides simple facilities for retrieving orbital coefficients and VB structure coefficients. It should not be used during a run of CASVB that has been invoked from inside MULTI.

START,*record.file*;

Coefficients are taken from *record.file*. The default value is 2100.2.

WRITE,*iwrite*;

Vectors in the symmetry orbital basis are written to channel *iabs(iwrite)*. The default action is to write these vectors to the standard output. If *iwrite* is negative, then the vectors are instead written to a binary file as a single record.

SPECIAL,*idim1, idim2, idim3, idim4*;

If present, this keyword must come last. The program attempts to retrieve from *record.file* a vector of length $idim1 * idim2 + idim3$, after first skipping *idim4* elements. The vector is written according to the setting of *iwrite1*. (Default *idim* values are zero.)

29.14 Examples

```

***, ch2                                ! A1 singlet state
geometry={angstrom
c
h1,c,1.117
h2,c,1.117,h1,102.4}
int
hf
multi;occ,4,1,2;closed,1                ! 6 in 6 CASSCF
natorb,,ci,save=3500.2;vbdump
casvb                                    ! Overlap-based VB using
save,3200.2                              ! the spin-coupled wavefunction
casvb                                    ! Energy-based VB calculation
start,,3200.2;save,3220.2
crit,energy
multi;occ,4,1,2;closed,1                ! Fully variational VB calculation
vb={start,,3220.2;save,3240.2;print,,,,2}
---

***, lih                                ! Fully variational VB calculation
r=2.8,bohr                               ! and geometry optimization.
basis={
s,1,921.300000,138.700000,31.940000,9.353000,3.158000,1.157000;
k,1.6,0.001367,0.010425,0.049859,0.160701,0.344604,0.425197;
s,1,0.444600,0.076660,0.028640;
p,1,1.488000,0.266700,0.072010,0.023700;
k,1.2,0.038770,0.236257;
s,2,13.36,2.013,0.4538,.1233;
k,1.2,0.032828,0.231204;}
geometry={li;h,li,r}
int;
hf;wf,4,1;
multi
occ,4,0,0,0
closed,0,0,0,0
natorb,,ci,save=3500.2
multi
vb
optg
---
```

30 SPIN-ORBIT-COUPLING

30.1 Introduction

Spin-orbit matrix elements and eigenstates can be computed using either the Breit-Pauli (BP) operator or spin-orbit pseudopotentials (ECPs). The *state-interacting* method is employed, which means that the spin-orbit eigenstates are obtained by diagonalizing $\hat{H}_{el} + \hat{H}_{SO}$ in a basis of eigenfunctions of \hat{H}_{el} . The full Breit-Pauli SO-operator can be used only for MCSCF wavefunctions. For MRCI wavefunctions, the full BP operator is used for computing the matrix elements between *internal configurations* (no electrons in external orbitals), while for contributions of external configurations a mean-field one-electron fock operator is employed. The error caused by this approximation is usually smaller than 1 cm^{-1} .

The program allows either the computation of individual spin-orbit matrix elements for a given pair of states, or the automatic setting-up and diagonalization of the whole matrix for a given set of electronic states. In the latter case, matrix elements over one-electron operators are also computed and transformed to the spin-orbit eigenstates (by default, the dipole matrix elements are computed; other operators can be specified on the GEXPEC or EXPEC cards, see section 4.13). Since it may be often sufficient to compute the spin-orbit matrix elements in a smaller basis than the energies, it is possible to replace the energy eigenvalues by precomputed values, which are passed to the spin-orbit program by the MOLPRO variable HLSDIAG.

30.2 Calculation of SO integrals

The one- and two-electron spin-orbit integrals over the BP Hamiltonian can be precomputed and stored on disk using the command

```
LSINT [,X] [,Y] [,Z]
```

X, Y, and Z specify the components to be computed. If none of these is given, all three are evaluated. The advantage of precomputing the integrals is that they can then be used in any number of subsequent SO calculations, but this may require a large amount of disk space (note that there are 6 times as many integrals as in an energy calculation). If the LSINT card is not given, the integrals are recomputed for one component at a time whenever needed, and destroyed at the end of the SO calculation. This reduces the disk space by a factor of 3, but may be expensive in terms of CPU if several SO calculations (e.g., for MCSCF and MRCI wavefunctions) are carried out.

The input for spin-orbit ECPs is described in section 11. Of course, in ECP-LS calculations the LSINT card is not needed.

30.3 Calculation of individual SO matrix elements

Individual spin-orbit matrix elements can be computed within the MRCI program using

```
TRANLS,record1.file, record2.file, bra2ms, ket2ms, lso;
```

where

<i>record1.file</i>	Record holding the bra-wavefunction.
<i>record2.file</i>	Record holding the ket-wavefunction. Both records must have been generated using the SAVE directive of the MRCI program.

<i>bra2ms</i>	$2 \times M_S$ value of the bra-wavefunction.
<i>ket2ms</i>	$2 \times M_S$ value of the ket-wavefunction.
<i>lsop</i>	Cartesian component of the Spin-orbit Hamiltonian. This can be one of LSX, LSY, or LSZ in all electron calculations, and ECPLSX, ECPLSY, or ECPLSZ in ECP calculations.

Since the spin-orbit program is part of the MRCI program, the TRANLS card must be preceded by a [MR]CI card. For the case that the matrix elements are computed for MCSCF wavefunctions, one has to recompute and save the CI-vectors using the MRCI program (see chapter 18), using the NOEXC directive to avoid inclusion of any further excitations out of the MCSCF reference function. If in the MRCI step several states of the same symmetry are computed simultaneously using the STATE directive, the matrix elements are computed for all these states. Note that the OCC and CLOSED cards must be the same for all states used in a TRANLS calculation.

The selection rules for the M_S values are $\Delta M_S = \pm 1$ for the LSX and LSY operators, and $\Delta M_S = 0$ for the LSZ operator. Note that $2M_S$ has to be specified, and so the selection rules applying to the difference of the input values are 0 or 2.

In all-electron SO calculations the value of the calculated spin-orbit matrix element is saved (in atomic units) in the MOLPRO variables TRLSX, TRLSY and TRLSZ for the x , y , and z components respectively. For ECP-LS calculations the variables TRECPLSX, TRECPLSY, and TRECPLSZ are used. Note that for imaginary matrix elements (i.e., for the x and z components of the SO Hamiltonian) the matrix elements are imaginary and the stored real values have to be multiplied by i . If matrix elements for several states are computed, all values are stored in the respective variable-arrays with the bra-states running fastest.

30.4 Calculation and diagonalization of the entire SO-matrix

HLSMAT,*type*, *record1*, *record2*, *record3*, ...

Computes the entire SO matrix and diagonalizes it using all states which are contained in the records *record1*, *record2*, *record3*, ... All records must have been generated using the SAVE directive of the MRCI program. *type* may be either LS for Breit-Pauli calculations, or ECP for ECP-LS calculations. By default, the eigenvalues and dipole transition matrix elements between the ground and excited states are printed.

As with the TRANLS card, the HLSMAT is recognized only by the MRCI program and must be preceded by a CI card. Also, the OCC and CLOSED cards must be the same for all states used in a HLSMAT calculation.

30.5 Modifying the unperturbed energies

Often it may be sufficient to compute the spin-orbit matrix elements in a smaller basis or at a lower computational level than the energies. It is therefore possible to replace the energy eigenvalues by precomputed values, which are passed to the spin-orbit program by the MOLPRO variable HLSDIAG. The energy values in HLSDIAG must be in exactly the same order as the states in the records given on the HLSMAT card. Before any spin-orbit calculation, the variable HLSDIAG must either be undefined or cleared (then the original energies are used), or must contain exactly the number of energies as the number of states treated in the subsequent spin-orbit calculation (use CLEAR, HLSDIAG to clear any previous values in the variable). It is the user's responsibility that the order of the energies in HLSDIAG is correct!

30.5.1 Print Options for spin-orbit calculations

PRINT,*option*₁=*value*₁, *option*₂=*value*₂, ...

where option can be

HLS	HLS=-1 only the SO energies and transition matrix elements between ground and excited states are printed (default). HLS ≥ 0: The SO matrix is printed. HLS ≥ 1: The property matrices are printed. HLS ≥ 2: The individual matrix elements are printed (same as OPTION, MATEL). HLS ≥ 3: Debugging information is printed.
VLS	VLS=-1: No print of eigenvectors (default). VLS ≥ 0: The eigenvectors are printed.

30.5.2 Options for spin-orbit calculations

Some options can be set using the OPTION directive (in any order)

OPTIONS [,WIGNER=*value*] [,HLSTRANS=*value*] [,MATEL=*value*]

where

WIGNER	This option determines whether the Wigner-Eckart theorem should be used when the SO matrix is determined. WIGNER=1 (default) uses the theorem, WIGNER=0 calculates each SO matrix element individually. This option is needed for test purposes only.
HLSTRANS	This option determines whether a SO matrix calculation should be performed in the not spin-symmetry adapted basis set (HLSTRANS=0), in the spin-symmetry adapted basis set (HLSTRANS=1, default) or with both basis sets (HLSTRANS=2). At present, symmetry adaption can only be performed for triplet states, where the following notation is used to indicate the symmetry adapted spin functions: $ S, M_S\rangle_+ = \frac{1}{\sqrt{2}}(S, M_S\rangle + S, -M_S\rangle)$, $ S, M_S\rangle_- = \frac{1}{\sqrt{2}}(S, M_S\rangle - S, -M_S\rangle)$. If only singlet and triplet states are considered, the spin-orbit matrix is blocked according to double-group symmetry and the eigenvalues for each each block are printed separately. In all other cases the HLSTRANS option is ignored.
MATEL	If the entire SO matrix is calculated using HLSMAT, the individual matrix elements are normally not shown. When the option MATEL=1 is given, the individual matrix elements and the contributions of the internal and external configuration classes are printed.

30.6 Examples

30.6.1 SO calculation for the S-atom using the BP operator

```

***,SO calculation for the S-atom
geometry={s}
basis={spd,s,vtz}                                !use uncontracted basis

rhf;occ,3,2,2,,2;wf,16,4,2                        !rhf for 3P state

multi                                             !casscf
wf,16,4,2;wf,16,6,2;wf,16,7,2;wf,16,1,0;state,3; !1D and 1S states
wf,16,4,0;wf,16,6,0;wf,16,7,0;                !3P states

ci;wf,16,1,0;save,3010.1;state,3                !save casscf wavefunctions using mrci
noexc
ci;wf,16,4,0;save,3040.1
noexc
ci;wf,16,6,0;save,3060.1
noexc
ci;wf,16,7,0;save,3070.1
noexc
ci;wf,16,4,2;save,3042.1
noexc
ci;wf,16,6,2;save,3062.1
noexc
ci;wf,16,7,2;save,3072.1
noexc

ci;wf,16,1,0;save,4010.1;state,3                examples/
ed=energy(1)                                     !mrci calculations for 1D, 1S states
es=energy(3)                                     !save energy for 1D state in variable ed
ep=energy                                         !save energy for 1S state in variable es
ci;wf,16,4,2;save,4042.1                         !mrci calculations for 3P states
ep=energy                                         !save energy for 3P state in variable ep
ci;wf,16,6,2;save,4062.1                         !mrci calculations for 3P states
ci;wf,16,7,2;save,4072.1                         !mrci calculations for 3P states
text,only triplet states, casscf

lsint                                             !compute so integrals

text,3P states, casscf
ci;hlsmat,ls,3042.1,3062.1,3072.1                !Only triplet states, casscf

text,3P states, mrci
ci;hlsmat,ls,4042.1,4062.1,4072.1                !Only triplet states, mrci

text,3P, 1D, 1S states, casscf
ci;hlsmat,ls,3010.1,3040.1,3060.1,3070.1,3042.1,3062.1,3072.1 !All states, casscf

text,only triplet states, use mrci energies and casscf SO-matrix elements
hlsdiag=[ed,ed,es,ed,ed,ed,ep,ep,ep]           !set variable hlsdiag to mrci energies
ci;hlsmat,ls,3010.1,3040.1,3060.1,3070.1,3042.1,3062.1,3072.1

```

30.6.2 SO calculation for the I-atom using ECPs


```

***,I
memory,5,M;
gprint,orbitals,civector,basis;
gthresh,energy=1.d-8,civector=1.d-8;
geometry={I};

basis={
!
! Iodine-ECP (Dirac-Fock) with SO-coupling
!
ecp,I,46,4,3;
1; 2, 1.00000000, 0.00000000; ! lokal term = 0
2; 2, 3.50642001, 83.09814545; 2, 1.74736492, 5.06370919; ! s-terme
4; 2, 2.99860773, 1/3* 81.88444526; 2, 3.01690894, 2/3* 83.41280402; ! p-terms with wei
    2, 1.59415934, 1/3* 2.32392477; 2, 1.19802939, 2/3* 2.72079843;
4; 2, 1.03813792, 2/5* 6.40131754; 2, 1.01158599, 3/5* 6.21328827; ! d-terms with wei
    2, 2.04193864, 2/5* 19.11604172; 2, 1.99631017, 3/5* 19.08465909;
4; 2, 2.64971585, -3/7* 24.79106489; 2, 2.75335574, -4/7* 24.98147319; ! f-terms with wei
    2, 0.49970082, -3/7* 0.27936581; 2, 0.79638982, -4/7* 0.70184261;
4; 2, 2.99860773, -2/3* 81.88444526; 2, 3.01690894, 2/3* 83.41280402; ! ECP-SO for p-ter
    2, 1.59415934, -2/3* 2.32392477; 2, 1.19802939, 2/3* 2.72079843;
4; 2, 1.03813792, -2/5* 6.40131754; 2, 1.01158599, 2/5* 6.21328827; ! ECP-SO for d-ter
    2, 2.04193864, -2/5* 19.11604172; 2, 1.99631017, 2/5* 19.08465909;
4; 2, 2.64971585, 2/7* 24.79106489; 2, 2.75335574, -2/7* 24.98147319; ! ECP-SO for f-ter
    2, 0.49970082, 2/7* 0.27936581; 2, 0.79638982, -2/7* 0.70184261;
!
! Iodine-basis
!
s,I,0.2027624,0.4080619,0.8212297,1.6527350,3.3261500;
c,1.5,-0.4782372,-0.5811680,0.2617769,0.4444120,-0.1596560;
s,I,0.05,0.1007509;
p,I,0.2027624,0.4080619,0.8212297,1.6527350,3.3261500;
c,1.5,0.4251859,0.2995618,0.0303167,-0.2064228,0.0450858;
p,I,0.05,0.1007509,0.01; ! diffuse p-Funktion wegen evt. neg. Part.Ldg
d,I,0.2,0.4;
f,I,0.3;
}
examples/
i_ecp.com

hf;occ,1,1,1,,1;wf,7,5,1; !scf for 2Pz
multi;occ,1,1,1,,1; !casscf with minimal active space
wf,7,2,1;wf,7,3,1;wf,7,5,1 !average 2P states
ci;wf,7,2,1;noexc;save,5000.2 !save casscf vector for 2Px state
ci;wf,7,3,1;noexc;save,5100.2 !save casscf vector for 2Py state
ci;wf,7,5,1;noexc;save,5200.2 !save casscf vector for 2Pz state
ci;wf,7,2,1;save,6000.2 !mrci for 2Px state
ci;wf,7,3,1;save,6100.2 !mrci for 2Py state
ci;wf,7,5,1;save,6200.2 !mrci for 2Pz state

multi;occ,1,2,2,,2 !casscf with larger active space
wf,7,2,1;wf,7,3,1;wf,7,5,1 !average 2P states
ci;wf,7,2,1;noexc;save,5010.2
ci;wf,7,3,1;noexc;save,5110.2
ci;wf,7,5,1;noexc;save,5210.2
ci;wf,7,2,1;save,6010.2
ci;wf,7,3,1;save,6110.2
ci;wf,7,5,1;save,6210.2

text,casscf, occ,1,1,1,,1
ci;hlsmat,ecp,5000.2,5100.2,5200.2 !do spin-orbit calculations
text,casscf, occ,1,2,2,,2
ci;hlsmat,ecp,5010.2,5110.2,5210.2

text,mrci, occ,1,1,1,,1
ci;hlsmat,ecp,6000.2,6100.2,6200.2
text,mrci, occ,1,2,2,,2
ci;hlsmat,ecp,6010.2,6110.2,6210.2

```

31 ENERGY GRADIENTS

31.1 Analytical energy gradients

MOLPRO uses two different gradient programs:

The CADPAC gradient program is based on the CADPAC integral routines by R. D. Amos. Currently, this program works for closed shell SCF, high spin RHF, and MCSCF. In the MCSCF case the wavefunction must either be fully optimized, or frozen core orbitals must be taken from a closed-shell SCF calculation. The present version does not work with generally contracted basis functions.

The ALASKA gradient program is based on the SEWARD integral routines by R. Lindh. It allows the calculation of gradients of generally contracted basis functions for closed shell SCF, open shell RHF, DFT, MCSCF, MP2, LMP2, and QCISD. It does not work with state averaged SA-MCSCF wave functions.

The ALASKA gradients are default for general contracted basis sets. In all other cases the CADPAC gradients are default. On most architectures these are faster than ALASKA gradients. However, it is possible to choose the gradients by defining the variable GRADTYP before calling the gradient program.

```
GRADTYP=ALASKA  
GRADTYP=CADPAC
```

The gradient program is called using the FORCE command:

```
FORCE
```

Normally, the FORCE command is not needed, since geometry optimizations should be performed using the OPTG procedure. An exception is the optimization of counterpoise corrected energies, which requires several force calculations (cf. section 32.2.19). Note that for computing gradients for state-averaged MCSCF a CPMCSCF is required in the MCSCF calculation (see CPMCSCF).

If no further data cards are given, the default is to evaluate the gradient for the last optimized wavefunction. In this case no further input is needed for ordinary gradient cases (the program remembers the records on which the wavefunction information is stored). An exception is the unusual case that several different CPMCSCF calculations have been formed in a previous MCSCF calculation. In this case the SAMC directive must be used to select the desired record.

31.1.1 Adding gradients (ADD)

```
ADD,factor,[NOCHECK];
```

If this card is present, the current gradient and energy are added to the previous ones using the given factor. This is useful for the optimization of counterpoise corrected energies (cf. 32.2.19). By default, the program will stop with an error message unless NOORIENT has been specified in the geometry input. This behaviour can be disabled by the NOCHECK option. This option should only be given if all gradients which are added together are evaluated at exactly the same nuclear geometry; otherwise wrong results could result due to unintended rotations of the system.

31.1.2 Scaling gradients (SCALE)

```
SCALE,factor;
```

If this card is present, the current gradient and energy are scaled by the give factor. This is sometimes useful for the optimization of counterpoise corrected energies (cf. 32.2.19).

31.1.3 Defining the orbitals for SCF gradients (ORBITAL)

ORBITAL,*record.file*;

In the SCF case, *record.file* specifies the location of the orbitals, which are used for constructing density matrices, etc. This card is only needed if the SCF for which the gradient is to be computed was not the most recent energy calculation.

For MCSCF wavefunctions, the ORBITAL card is not needed, because the location of the orbitals is stored in the MCSCF dump record.

31.1.4 MCSCF gradients (MCSCF)

MCSCF,*record.file*;

Triggers code for MCSCF gradient. *record.file* specifies the location of information dumped from the MCSCF program, using a SAVE ,GRD=*recmc.filmc* card. This card is not needed if the FORCE command appears directly after the corresponding MCSCF input, since the program automatically remembers where the MCSCF information was stored. The same is true if OPTG is used.

31.1.5 State-averaged MCSCF gradients (SAMC)

Normally, no further input is required for computing gradients for state-averaged MCSCF. Note, however, that a CPMCSCF ,GRAD,*state* directive is required in the SA-MCSCF calculation (see CPMCSCF). The gradients are then computed automatically for the *state* specified on the CPMCSCF card. The same is true for difference gradients (CPMCSCF ,DGRAD,*state1*, *state2*) and non-adiabatic coupling matrix elements (CPMCSCF ,NACM,*state1*, *state2*). It is possible to do several coupled-perturbed MCSCF calculations one after each other in the same MCSCF. In this case FORCE would use the last solution by default. The information from the CPMCSCF is passed to the FORCE program in a certain records (default 5101.1, 5102.1, ...). If several CPMCSCF calculations are performed in the same MCSCF, several such records may be present, and a particular one can be accessed in the FORCE program using the SAMC directive:

SAMC,*record*.

An alias for SAMC is CPMC. For compatibility with earlier versions one can also use

NACM,*record*

for non-adiabatic couplings or

DEMC,*record*

for difference gradients.

Example:

```
multi;
....
state,3
cpmcscf,nacm,1.1,2.1,save=5101.1 !do cpmcscf for coupling of states 1.1 - 2.1
```

```

cpmcscf,nacm,1.1,3.1,save=5102.1 !do cpmcscf for coupling of states 1.1 - 3.1
cpmcscf,nacm,2.1,3.1,save=5103.1 !do cpmcscf for coupling of states 2.1 - 3.1

force:samc,5101.1; !compute NACME for states 1.1 - 2.1
force:samc,5102.1; !compute NACME for states 1.1 - 3.1
force:samc,5103.1; !compute NACME for states 2.1 - 3.1

```

See also test job `lif_nacme.test`.

31.1.6 Non-adiabatic coupling matrix elements (NACM)

see SAMC

31.1.7 Difference gradients for SA-MCSCF (DEMC)

see SAMC

31.1.8 Example

```

***, Calculate SCF-Gradients for Water
alpha=104 degree !set geometry parameters
r=1 ang
geometry={O; !define z-matrix
          H1,o,r;
          H2,o,r,H1,alpha}
basis=vdz !basis set
hf !do scf
forces !compute gradient for SCF
mp2 !mp2 calculation
forces !mp2 gradients
multi !casscf calculation
forces !casscf gradient

```

examples/
h2o_forces.com

31.2 Numerical gradients

It is possible to compute gradients by finite differences using

```
FORCE, NUMERICAL, STARTCMD=command
```

where *command* is the first command in the input needed for the current energy calculation. The *command* must be found in the input *before* the FORCE card, i.e., the sequence of input cards starting with *command* and ending with FORCE is used to compute the energies. For example, in order to compute numerical gradients for `ccsd(t)`, it is necessary to do a HF and a CCSD(T) calculation at each geometry. The input would read

```

hf
ccsd(t)
forces,numerical,startcmd=hf

```

The program will then automatically repeat HF and CCSD(T) at as many geometries as needed for evaluating the gradient. The keyword NUMERICAL implies that also STARTCMD must be given, otherwise an error results.

31.2.1 Choice of coordinates (COORD)

By default, the numerical gradients are computed relative to all variables on which the z-matrix depends. If the z-matrix depends on no variables or on $3N$ variables, the gradient is computed for all $3N$ coordinates and symmetrical displacement coordinates are used to evaluate the gradient. This yields the minimum computational effort.

These defaults can be modified using the COORD directive:

COORD,*coord_type*,[*displacement_type*]

where *coord_type* can be one of the following:

- | | |
|------------|--|
| ZMAT | Compute the numerical gradients for all variables on which the geometry depends (default). |
| 3N or CART | Compute the gradients for all $3N$ nuclear coordinates. This is the default if the z-matrix does not depend on variables or if the xyz input format is used. If this option is used and the original geometry is given in z-matrix form, the z-matrix is lost. |

The specification of *displacement_type* is optional and only affects the numerical calculation of the gradient for $3N$ coordinates. By default, SYM is used.

- | | |
|--------|---|
| SYM | Use symmetrical displacements. This yields the minimum number of displacements and always preserves the symmetry of the wavefunction. This is the only recommended option. |
| CART | Displacements are generated for all $3N$ cartesian coordinates. This is normally not recommended, since in cases in which molecular symmetry is present it generates far more displacements than needed. Also, the wavefunction symmetry is not preserved, and the calculation must be done in C1 symmetry. |
| UNIQUE | As CART, but symmetry-equivalent displacements are eliminated. Not recommended either. |

31.2.2 Numerical derivatives of a variable

Numerical derivatives of the value of a variable can be computed using

VARIABLE=*name*

The default is to compute the gradient of the current energy.

31.2.3 Stepsizes for numerical gradients

By default, the numerical step sizes are 0.01 bohr for distances or cartesian coordinates, and 1 degree for angles. These defaults can be changed using

RSTEP=*dr*

ASTEP=*da*

where *dr* is the displacement for distances (or cartesian coordinates) in bohr, and *da* is the displacement for angles in degree. The value of RSTEP is used for symmetrical displacements. The step sizes for individual variables can be modified using

STEP, *varname*=*value*,...

where the *value* must be in atomic units for distances and in degree for angles.

32 GEOMETRY OPTIMIZATION

It is possible to invoke the geometry optimization program in two ways: Automatic geometry optimization using the OPTG command or geometry optimization with an explicit input sequence using OPT. Normally, geometry optimizations should be done using the OPTG command. This calls OPT and all other necessary programs automatically.

32.1 Geometry optimization step (OPT)

The OPT program reads the geometry definitions and geometry parameters, energies, and gradients of the present and previous points from a geometry record. It is necessary that the gradients have been computed using the FORCE command (see above) before calling OPT. The program then predicts a new optimum geometry (i.e. takes one optimization step), by default optimizing all variable parameters on which the geometry depends, and writes this back to the geometry record. The optimization is performed in a space of scaled parameters, the scaling being such that the initial hessian matrix has unit diagonal elements. The variable OPTCONV is set to the length of the step taken in scaled parameter space, and can be tested after the OPT step using the IF command, to decide whether to return for another geometry. For subcommands of OPT see OPTG.

OPT can also be used for automatic geometry optimization using a sequence of input commands. In this case, one can specify the first input command needed for computing the energy and gradient using the STARTCMD option:

```
OPT , STARTCMD=command
```

Similar to the calculation of numerical gradients (see above), *command* must be found in the input *before* the OPT card, i.e., the sequence of input cards starting with *command* and ending with OPT defines one optimization step. For example, in order to optimize the geometry at the ccsd(t) level using numerical gradients the following input could be used

```
hf                !optimize orbitals
ccsd(t)           !compute ccsd(t) geometry
forces,numerical ,startcmd=hf  !compute numerical ccsd(t) gradients
opt ,startcmd=hf  !optimize geometry
```

The convergence criteria are the same as explained below for the OPTG procedure. The convergence thresholds can be modified using further options on the OPT card, exactly in the same way as explained below for OPTG. For example, the threshold for the gradient can be changed using

```
OPT , STARTCMD=command , GRAD=1 .d-4
```

Further subcommands for OPT are possible, which are the same as for OPTG described in the next section.

32.2 Automatic geometry optimization (OPTG)

The OPTG command is used to perform automatic geometry optimizations for all kinds of wavefunctions. The coordinates to be optimized can be chosen using the COORD directive (see section optgeo:coord). Various optimization methods can be selected as described in section 32.2.4. MOLPRO allows minimization (i.e. search for equilibrium geometries), transition state optimization (i.e. search for saddle points on energy surfaces), and reaction path following. The standard algorithms are based on the *rational function* approach and the *geometry DIIS* approach.

Also available is the *quadratic steepest descent following* method of Sun and Ruedenberg (see J. Sun and K. Ruedenberg, *J. Chem. Phys.* **99**, 5257 (1993)). This method is often advantageous in Transition State searches. For a detailed discussion of the various minimization algorithms see (see F. Eckert, P. Pulay and H.-J. Werner, *J. Comp. Chem* **18**, 1473 (1997)).

The OPTG must directly follow the input for the wavefunction used in the geometry optimization. It will call FORCE, OPT, INT, and, as needed, HF, RHF, MCSCF, CI, CCSD etc. For each of these programs, the input file is automatically repositioned to the last corresponding input before the OPTG card; so any input for RHF, MCSCF, CI, CCSD etc. can be used and will be correctly processed. It is essential, however, that the most recently optimized orbitals are used in the wavefunction for which the geometry is optimized. Any input needed for OPTG must directly follow the OPTG card. The gradients are computed analytically for HF, DFT, MP2, QCISD, or MCSCF wavefunctions; otherwise the gradients are computed by finite differences (see OPTG, NUMERICAL). Davidson corrected energies or excited state energies can be optimized using the VARIABLE and STATE subdirective.

Various options, in particular convergence criteria, can be specified on the OPTG command:

OPTG,*key1=value, key2=value,*

where *key* can be

MAXIT	to set the maximum number of optimization cycles. The default is 50.
GRAD	sets the required accuracy of the optimized gradient. The default is $3 \cdot 10^{-4}$.
STEP	to set the convergence threshold for the geometry optimization step; if <i>value</i> ≥ 1 , the threshold is set to 10^{-value} . The default is $3 \cdot 10^{-4}$.
ENERGY	sets the required accuracy of the optimized energy. The default is $1 \cdot 10^{-6}$.
GAUSSIAN	Use Gaussian convergence criteria.
SRMS	sets (for Gaussian convergence criterion) the required accuracy of the RMS of the optimization step. The default is 0.0012.
GRMS	sets (for Gaussian convergence criterion) the required accuracy of the RMS of the gradient. The default is $3 \cdot 10^{-4}$.
BAKER	Use Baker's convergence criteria (see J. Baker, <i>J. Comp. Chem.</i> 14 , 1085 (1993)).
NUMERICAL	Force the use of numerical gradients.

The standard MOLPRO convergence criterion requires the maximum component of the gradient to be less than $3 \cdot 10^{-4}$ [a.u.] and the maximum energy change to be less than $1 \cdot 10^{-6}$ [H] or the maximum component of the gradient to be less than $3 \cdot 10^{-4}$ [a.u.] and the maximum component of the step to be less than $3 \cdot 10^{-4}$ [a.u.].

It is also possible to use the convergence criterion of the Gaussian program package. It is somewhat weaker than the MOLPRO criterion and requires the maximum component of the gradient to be less than $4.5 \cdot 10^{-4}$ [a.u.] and the root mean square (RMS) of the gradient to be less than $3 \cdot 10^{-4}$ [a.u.] as well as the maximum component of the optimization step to be less than 0.0018 [a.u.] and the RMS of the optimization step to be less than 0.0012 [a.u.].

The defaults for the convergence parameters can also be changed by using a global GTHRESH directive, i.e.

GTHRESH, OPTSTEP=*step*, OPTGRAD=*grad*, ENERGY=*energy*;

32.2.1 Optimization coordinates (COORD)

It is possible to use various coordinate types and algorithms for the optimization. This can be controlled by additional subcommands as described in this and the following subsections.

COORD,[*opt_space*],[*opt_coord*],[*displacement_type*],[*option*]

This option chooses the optimization space and the coordinate system in which the optimization takes place.

opt_space defines the parameters to be optimized. By default, if the geometry input is given in z-matrix format, all variables on which the z-matrix depends are optimized. Subsets of the variables on which the z-matrix depends can be chosen using the ACTIVE or INACTIVE subdirectives. If the z-matrix depends on no variables or xyz input is used, all $3N$ cartesian coordinates are optimized.

opt_coord determines the coordinates in which the optimization takes place. By default, local normal coordinates are used. Optionally cartesian coordinates or natural internal coordinates can be used.

displacement_type specifies how numerical gradients and Hessians are computed. This defaults to symmetric displacement coordinates and should normally not be modified.

These defaults can be modified using the COORD directive.

opt_space can be one of the following:

ZMAT	Optimize all variables on which the z-matrix depends (default)
3N	Optimize all $3N$ cartesian coordinates. Z-Matrix input coordinates will be destroyed on this entry.

opt_coord can be one of the following:

NORMAL	Optimization in local normal coordinates. This is default if the Model Hessian is used to approximate the Hessian.
NONORM	Don't use local normal coordinates.
BMAT[= <i>filename</i>]	Use Pulay's <i>natural internal coordinates</i> , see G. Fogarasi, X. Zhou, P. W. Taylor and P. Pulay <i>J. Am. Chem. Soc.</i> 114 , 8191 (1992); P. Pulay, G. Fogarasi, F. Pang, J. E. Boggs <i>J. Am. Chem. Soc.</i> 101 , 2550 (1979)). Optionally, the created coordinates plus additional information about this optimization is written to the specified file. These coordinates resemble in part the valence coordinates used by vibrational spectroscopists, and have the advantage of decreasing coupling between different modes. This often increases the speed of convergence. The use of this option is highly recommended, especially in minimization of large organic molecules with rings. Nevertheless you should keep in mind that these coordinates are constructed automatically, and there exist exotic bond structures which might not be treated properly (e.g. weakly bonded species as in transition state optimizations). In such a case, if the BMAT optimization converges slowly or leads to symmetry-breaking errors, you should try another optimization method and/or cartesian or Z-Matrix coordinates.

displacement_type can be one of the following (affects only numerical gradients):

SYM	Use symmetric displacement coordinates (default). This is the only recommended option.
CART	Use $3N$ cartesian displacements (not recommended). This requires many more energy calculations than necessary and does not preserve the molecular symmetry.
UNIQUE	Use symmetry-unique cartesian displacements (not recommended)

If *option* is set to [NOROT], the cartesian coordinates are not transformed to minimize rotations.

32.2.2 Defining active geometry parameters (ACTIVE)

ACTIVE,*param*;

Declares variable name *param* to be active in the optimization. By default, initially all variables on which the geometry depends are active; inclusion of an ACTIVE card makes all parameters inactive unless explicitly declared active (see also INACTIVE).

32.2.3 Defining inactive geometry parameters (INACTIVE)

INACTIVE,*param*;

Declares variable name *param* to be inactive in the optimization. If any ACTIVE card appears in the input, this card is ignored! (see also ACTIVE)

32.2.4 Selecting the optimization method (METHOD)

METHOD,*key*;

key defines the optimization method.

For *minimization* the following options are valid for *key*:

RF	Rational Function method (default).
AH	Augmented Hessian method. This is similar to RF algorithm but uses a more sophisticated step restriction algorithm.
DIIS	Pulay's Geometry DIIS method. As an additional option you may add the number of geometries to be used in GDIIS interpolation (default 5) and the interpolation type (i.e. the subspace in which the GDIIS interpolation is made). METHOD, DIIS, <i>number</i> , <i>type</i> <i>type</i> may be GRAD interpolation using the gradients (default), working good for rigid molecules, STEP interpolation using Quasi-Newton steps which could be advantageous in dealing with very floppy molecules, ENER interpolation using energies, which is an intermediate between the above two.
QSD	Quadratic steepest descent method of Sun and Ruedenberg.
SRMIN	Old version of QSD.

For *transition state* searches (invoked with the ROOT option, see section 32.2.6) *key* can be

RF	Rational Function method (default).
DIIS	Pulay's Geometry DIIS method (see above).
QSD	Quadratic Steepest Descent Transition State search using the image hessian method (see J. Sun and K. Ruedenberg, <i>J. Chem. Phys.</i> 101 , 2157 (1994)) The use of this option is recommended for transition state searches – especially in complicated cases. The optimization step is checked and the hessian is recalculated when approaching a troublesome region of the PES. Thus this method is somewhat safer (and often faster) in reaching convergence than the RF or DIIS method. The hessian recalculation safeguard may be turned off using the METHOD , QSD , NOHESS input card.
SRTRANS	Old version of QSD.

For *reaction path following* the input key is

QSDPATH	Quadratic Steepest Descent reaction path following. This methods determines reaction paths (intrinsic reaction coordinates, IRCs) by following the exact steepest descent lines of subsequent quadratic approximations to the potential energy surface. The hessian matrix is calculated numerically at the first optimization step and subsequently updated by Powell or BFGS update. If a given arc length of the steepest descent lines is exceeded, the hessian is recalculated numerically (see OPTION section 32.2.17). For details see J. Sun and K. Ruedenberg, <i>J. Chem. Phys.</i> 99 , 5269 (1993) It is also possible to recalculate the hessian after each m steps using the NUMHES, m command (see section 32.2.15). If the hessian matrix is recalculated in every optimization step (NUMHES,1) a algorithm different to the one with updated Hessians is used, which is very accurate. Using the PRINT , OPT card, this algorithm prints in every optimization step a <i>reaction path point r</i> which is different from the point where the energy and the gradient is calculated but closer to the real reaction path (for further details of the algorithm see J. Sun and K. Ruedenberg, <i>J. Chem. Phys.</i> 99 , 5257 (1993)). For further input options of the QSD reaction path following see OPTION section 32.2.17.
SRSTEEP	Old Version of QSDPATH.

In *transition state searches* and *reaction path following* Z-Matrix coordinates should be used. Although it is also possible to use cartesian or BMAT coordinates, the computational effort is usually much larger, since the hessian matrix has to be calculated numerically in all $3 * N$ possible degrees of freedom.

32.2.5 Approximating hessian matrix elements (HESSIAN)

By default, the MOLPRO geometry optimization utilizes a force field approximation to the hessian (“Model Hessian”, see R. Lindh, A. Bernhardsson, G. Karlström and P. Malmqvist *Chem. Phys. Lett.* **241**, 423 (1995).) which speeds up convergence significantly. The Model Hessian is parameterized for the elements up to the third row and is used by default unless the molecule contains atoms from higher rows. Alternatively, the model Hessian of Schlegel can be used.

HESSIAN[,key], value, param1, param2;

where key can be

MODEL	Use Lindh's Model Hessian in optimization (default).
-------	--

MODEL, SCHLEGEL	Use Schlegel Model Hessian.
MODEL, VDW	Add vdW terms to Lindh's Model Hessian.
SCHLEGEL	Same as MODEL, SCHLEGEL]
VDW	Same as MODEL, VDW
NOMODEL	Don't use Model Hessian approximation to the hessian.

For minimizations, the Model Hessian provides very good approximations the Hessian matrix, improving convergence rapidly, so it was chosen as default. At present it is implemented for the first three row elements.

You may also put in individual matrix elements of the hessian: *value* sets starting value for hessian matrix element between parameters *param1*, *param2*. If *param2* is omitted it defaults to *param1* (diagonal element). If the Model Hessian is disabled, the initial hessian is diagonal, with values $1 \text{ hartree} \cdot \text{bohr}^{*-2}$ for all lengths, $1 \text{ hartree} \cdot \text{radian}^{*-2}$ for all angles. This is usually quite reasonable except for cases such as dihedral angles. A reasonable strategy for complicated cases is to perform an optimization with a small basis set at the SCF level with PRINT, HESSIAN in order to obtain an approximate starting hessian. These values are set before processing the START record (see above). This option is obsolete if the Model Hessian is used (default unless heavy elements are present).

In transition state searches the hessian matrix is evaluated numerically (see NUMHES section 32.2.15). Alternatively, the cartesian hessian matrix evaluated in a previous frequency calculation (see FREQUENCIES section 33) can be used with the HSTART command (see section 32.2.16). It is also possible to use the numerical hessian or the hessian from a frequency calculation in minimizations. Note that numerical Hessians cannot be computed when dummy atoms holding basis functions are present.

32.2.6 Transition state (saddle point) optimization (ROOT)

ROOT,*root*

Specifies the eigenvector of the hessian to be followed.

root=1 specifies a minimization (default).

root=2 specifies a transition state (saddle point) optimization.

In the present implementation a saddle point search is possible with the rational function method (METHOD, RF), the geometry DIIS method (METHOD, DIIS) and the quadratic steepest descent method of Sun and Ruedenberg (METHOD, SRTRANS).

Note that convergence is usually much more difficult to achieve than for minimizations. In particular, a good starting geometry and a good approximation to the hessian is needed. The latter is achieved by evaluating the hessian numerically (see NUMHES section 32.2.15) or using a precomputed cartesian hessian with the HSTART command (see section 32.2.16).

32.2.7 Saving optimization information (SAVE)

SAVE,*record*

Specifies a record on which the geometry definitions, parameters, energies, and gradients for the present optimization are stored. By default, this is record 700, and it is overwritten in each new optimization. If it is intended to use a START directive in a subsequent optimization, a different number should be given (e.g. 710, 720). The geometry record is saved on all permanent files.

32.2.8 Restarting a geometry optimization (START)

START,*record,first,last*

Specifies a record from which the geometry definitions, parameters, energies, and gradients of a previous optimization are read. This record must have been written with a SAVE card in a previous optimization. It should not be the geometry default record 700, except if the optimization was aborted and no other calculations have been performed since then.

The information from the previous calculation is to be used to construct an approximate hessian for a starting guess. *first* and *last* specify the first and last geometry blocks, respectively, to be used from *record*. The default is to use all blocks.

It is not necessary that the method used for the previous optimization is identical to the present one. This can be useful if several optimizations for different methods follow each other. For example, the hessian from an SCF optimization can be used as starting guess for a subsequent MCSCF optimization. However, it is required in such a case that the SAVE and START records are different.

32.2.9 Setting a maximum step size (STEP)

STEP,*steplength,drmax,damax,drmax1,damax1*

steplength is the initial step length in the scaled parameter space (default 0.3). In the AH-method this is dynamically adjusted, and can have a maximum value *ahmax* (see TRUST).

drmax is the initial max change of distances (in bohr, default 0.3). In the AH-method this is dynamically adjusted up to a maximum value of *drmax1* (default 0.5 bohr).

damax is the initial max change of angles (in degree, default 2). In the AH-method this is dynamically adjusted up to a maximum value of *damax1* (default 10 degrees).

32.2.10 Number of point used in hessian update (UPDATE)

UPDATE,*type=nstep*

This option chooses the update type to be used and limits the number of points used for the hessian update to *nstep*. In minimizations *type* may be

BFGS	Use BFGS update of hessian (default).
IBFGS	Use BFGS update of the inverse hessian.
CGRD	Use Conjugate Gradient update (see also CUT,TRUST).
NONE	Don't do any update.

The default number of steps used in hessian update procedures is 5. If there are symmetry constraint in the coordinates of the optimization, the default number may be lower than five. On input, *nstep* steps will be used for update, ignoring any symmetry constraints.

In transition state optimizations *type* may be

PMS	Combined Powell/Murtagh-Sargent update of hessian (default).
POWELL	Use Powell's update of the hessian.
MS	Use update procedure of Murtagh and Sargent.
NONE	Don't do any update.

32.2.11 Redefining the trust ratio (TRUST)

TRUST,*ratio,ahmax*

ratio determines the radius around the current minimum in which points are used to update the Hessian with the conjugate gradient method (default 0.5; see also UPDATE).

ahmax is the maximum step size allowed in the Augmented Hessian procedure. This refers to the scaled parameter space (default 0.5). The initial step size is *stepmx* (see STEP card).

32.2.12 Setting a cut parameter (CUT)

CUT,*threshold*

Specifies a threshold for orthonormalization used in conjugate gradient update of hessian (default 1.d-3; see also UPDATE).

32.2.13 Line searching (LINESEARCH)

LINESEARCH,*iflag,thrlmin,thrlmax*

Interpolate the geometry of the stationary point (minimum or saddle point) by a quartic polynomial between the current and the previous geometry. If *iflag=0* or no *iflag* is set, the next optimization step will be taken from the interpolated geometry using the interpolated energy and gradient. If *iflag=1* the energy and gradient will be recalculated at the interpolated geometry before taking the new optimization step. Note though, that the additional effort of recalculating the energy and gradient is usually not met by the increase of the convergence rate of the optimization. *thrlmin* and *thrlmax* are min and max thresholds for the recalculation of the energy and the gradient in case *iflag=1*. I.e. the recalculation just takes place if the interpolated geometry isn't too close to the actual geometry *thrlmin* and isn't too remote from the actual geometry *thrlmax*. Default values are *thrlmin=0.001* and *thrlmax=0.05* in the scaled parameter space of the optimization.

32.2.14 Numerical gradients (NUMERICAL)

NUMERICAL,*active₁=step₁, active₂=step₂ . . . ;*

With this option the gradients are computed by finite differences. *step_i* is the increment for the active geometry parameter *active_i*. For active parameters which are not specified, the default values are used. By default, the increment is 0.01 bohr for bond distances and 0.5 or 1 degree for angles less than or greater than 90 degrees, respectively. These defaults can be modified by specifying *RSTEP=value*, *ASTEP=value* on the numerical card. Note that step sizes which are too large may lead to optimization failure. For each active variable, two additional energy calculations are necessary for each geometry optimization step – so this may be expensive! Symmetrical displacement coordinates are normally used to minimize the number of energy calculations. (see COORD keyword section 32.2.1).

For optimization of special energies see VARIABLE section 32.2.18.

32.2.15 Numerical Hessian (NUMHES)

NUMHES,*mstep,type,icalc,thresh;*

This option allows you to calculate numerical second derivatives of the energy by finite differences. If you use analytical gradients, these are differentiated once whereby it is possible to use forward differences (needs one additional gradient calculation for each coordinate) or central differences (more accurate, needs two additional gradient calculations for each coordinate). For transition state optimizations it is usually sufficient to use forward differences. If you use numerical gradients the energy is differentiated twice. In this case only central differences are possible. *mstep* defines the number of optimization steps after which the numerical hessian is recalculated:

<i>mstep</i> =-1	Don't calculate numerical hessian (default for minimization).
<i>mstep</i> =0	Calculate numerical hessian only once at the start of the optimization (default for transition state searches).
<i>mstep</i> = <i>n</i>	Calculate numerical hessian after each <i>n</i> optimization steps. This is useful for difficult transition state optimizations (e.g. if the eigenvalue structure of the hessian changes during the optimization).

type defines the finite differences to be used:

<i>type</i> =0	Use forward differences (default).
<i>type</i> =1	Use the more accurate central differences.

icalc defines in which way the Hessian matrix shall be recalculated:

<i>icalc</i> =0	Static regeneration: Recalculate complete Hessian matrix numerically after each <i>mstep</i> optimization steps (default).
<i>icalc</i> =1	Partial regeneration: Recalculate selected Hessian matrix elements if the relative deviation of this element before and after update (see UPDATE, section 32.2.10) is larger than <i>thresh</i> . If <i>thresh</i> is not specified, a default value of <i>thresh</i> = 0.05 (i.e. a maximum deviation of 5%) is used.
<i>icalc</i> =2	Dynamic regeneration: Recalculate complete Hessian matrix if the RMS deviation of the Hessian matrix before and after update is larger than <i>thresh</i> . If <i>thresh</i> is not specified a default value of <i>thresh</i> = 0.5 a.u. is used.

Note that the static regeneration of the complete Hessian matrix after *mstep* iterations is not disabled if the partial (*icalc*=1) or dynamical (*icalc*=2) regeneration is used; i.e. if you want to use only the partial or dynamical regeneration, you should set *mstep* to zero.

If cartesian coordinates are used, the molecular symmetry can be used to shorten the number of gradient and/or energy calculations (see COORD keyword section 32.2.1).

32.2.16 Hessian starting guess from a frequency calculation (HSTART)

HSTART,*irec.ifil*;

This option allows you to use a cartesian hessian matrix, computed in a FREQUENCIES calculation (or using the *cpmcscf*, *hessian* option in multi) as a starting guess to the Hessian in a geometry-optimization. The Hessian is transformed automatically into the coordinate system the optimization is performed in. This option might be very useful in transition state optimization, i.e. by using a cheap analytical SCF Hessian as a starting guess to a higher level

optimization (See example `hcn_mp2_ts.com` below). Note that dummy atoms are ignored, and therefore the coordinates of dummy atoms should not depend on variables.

The starting hessian is read from record *irec* on file *ifil* (compare the SAVE subcommand of FREQUENCIES, section 33ff.). If no *irec.ifil* is given, the last FREQ record found on file *ifil=2* will be used.

Note that the HSTART option has higher priority than the NUMHES card. If a hessian is found using the HSTART command, the NUMHES command will be ignored.

Further note that if the hessian is read from file 1, it is assumed that it has been calculated at the current geometry and the hessian update is disabled. On the other hand, if the hessian is found on file 2 or 3, it is assumed that it has only been computed in the first iteration, and the hessian update is performed in subsequent iterations.

```

***, HCN <--> NHC Isomerization - Transition State Optimization and Frequencies

memory,1,m

l1=1.18268242 ang
l2=1.40745082 ang
a1=55.05153416 degree

basis=3-21G

geometry={nosymm;
          C
          N,1,l1
          H,2,l2,1,a1}

int
hf                ! HF-SCF

frequencies,analytical ! Vibrational frequencies for HF-SCF (analytical Hessian)
save,5300.2        ! Save Hessian in record 5300.2
print,low          ! Print low vibrational frequencies
print,imag        ! Print imaginary vibrational frequency

mp2                ! MP2

optg               ! Optimize Geometry
root,2             ! Transition State Search
method,rf          ! Rational Function Optimizer
step,1.0,0.3,10   ! Allow large stepsize
hstart,5300.2     ! Use SCF Hessian from FREQ record 5300.2 as starting guess

frequencies       ! Vibrational frequencies for MP2 (numerical Hessian)
print,low         ! Print low vibrational frequencies
print,imag       ! Print imaginary vibrational frequency
---
```

examples/
hcn_mp2_ts.com

32.2.17 Reaction Path Following options

OPTION,*key=param*;

where *key* can be

IDIR If starting at a transition state (or near a transition state) determine where to take the first step. If IDIR=0 is chosen, the first step will be towards

the transition state. This is the default. If `IDIR=1` is given in the input the first optimization step will be along the "transition vector" i.e. the hessian eigenvector to the smallest eigenvalue which points down towards the minimum. If using a larger `IDIR` parameter, the first step will be larger; if using a negative value, the first step will be in the opposite direction.

<code>STPTOL</code>	If using an updated hessian matrix, this parameter determines what update to take. If the step size between two subsequent points on which the steepest decent lines are puzzled together is smaller than <i>stptol</i> (i.e. if we are close to a minimum) the BFGS update is used, otherwise it is Powell update. The default value of <i>stptol</i> is $1.d - 6$.
<code>SLMAX</code>	This option is only valid with the old version of the reaction path following algorithm (i.e. <code>METHOD, SRSTEEP</code>). In this algorithm <code>slmax</code> determines the frequency of the recalculation of the numerical hessian. If the total step size of the last steps exceeds <i>slmax</i> the hessian will be recalculated, otherwise it will be updated. By default <i>slmax</i> is two times the maximum step size of the optimization step <i>steplength</i> (see <code>STEP</code> section 32.2.9). If you are using <code>METHOD, QSD</code> , the <code>SLMAX</code> option is obsolete and the <code>NUMHES</code> command (see above) should be used instead.

32.2.18 Optimizing energy variables (`VARIABLE`)

`VARIABLE, name;`

Defines a variable *name* which holds the energy value to be optimized in using finite differences. By default, this is `ENERGY(1)` as set by the most recent program. Other variables which can be used are

<code>ENERGY(i)</code>	holds last energy for state <i>i</i> .
<code>ENERGR(i)</code>	holds last reference energy for state <i>i</i> .
<code>ENERGD(i)</code>	holds last Davidson corrected energy for state <i>i</i> .
<code>ENERGP(i)</code>	holds last Pople corrected energy for state <i>i</i> .
<code>ENERGC</code>	holds CCSD (QCI, BCCD) energy in CCSD(T) [QCI(T), BCCD(T)] calculations (single state optimization).
<code>ENERGT(1)</code>	holds CCSD(T) energy in CCSD(T) calculations (single state)
<code>ENERGT(2)</code>	holds CCSD[T] energy in CCSD(T) calculations (single state).
<code>ENERGT(3)</code>	holds CCSD-T energy in CCSD(T) calculations (single state).

These variables are set automatically by the CI and/or CCSD programs. It is the user's responsibility to use the correct variable name; an error exit occurs if the specified variable has not been defined by the last program or the user.

Note: The use of the `VARIABLE` option triggers `NUMERICAL`, so optimization can be very inefficient!

32.2.19 Optimizing counterpoise corrected energies

Geometry optimization of counterpoise corrected energies is possible by performing for the total system as well as for each individual fragment separate `FORCE` calculations. The gradients and

energies are added using the `ADD` directive. This requires that `NOORIENT` has been specified in the geometry input, in order to avoid errors due to unintended rotation of the system. This default can be disabled using the `NOCHECK` option, see `ADD` above.

The way a counterpoise corrected geometry optimization works is shown in the following example. Note that the total counterpoise corrected energy must be optimized, not just the interaction energy, since the interaction energy depends on the monomer geometries and has a different minimum than the total energy. The interaction energy could be optimized, however, if the monomer geometries were frozen. In any case, the last calculation before calling `OPT` must be the calculation of the total system at the current geometry (in the example below the dimer calculation), since otherwise the optimizer gets confused.

```

***,HF dimer mp2/CP optimization

basis=avdz

maxit=20                !max number of iterations

text,OPTIMIZED VALUES OF GEOMETRY VARIABLES

RFF=      5.31431160
R1=      1.75768738
R2 =     1.75298524
THETA1 =  7.03780227
THETA2 = 111.25930975

geometry={x;noorient    !noorient must be specified since gradients are added
          f1
          f2 f1 rff
          h1 f1 r1  f2 theta1
          h2 f2 r2  f1 theta2 h1 180.}

do iter=1,maxit        !optimization loop

text, CALCULATION AT LARGE SEPARATION

rff_save=rff          !save current rff distance
rff=1000              !calculation at large separation

text, HF1 MONOMER

dummy,f2,h2;          !second hf is now dummy
hf;                   !scf for first monomer
mp2;                  !mp2 for first monomer
ehflinf=energy        !save mp2 energy in variable
forces;               !compute mp2 gradient for first monomer

text, HF2 MONOMER

dummy,f1,h1;          !first hf is now dummy
hf;                   !scf for second monomer
mp2;                  !mp2 for second monomer
ehf2inf=energy        !save mp2 energy in variable
forces;               !compute mp2 gradient for second monomer
add,1                 !add from previous gradient

rff=rff_save          !reset HF - HF distance to current value

text, HF1 CP MONOMER

dummy,f2,h2;          !second hf is now dummy
hf;                   !scf for first monomer
mp2;                  !mp2 for first monomer
ehf1=energy           !save mp2 energy in variable
forces;               !compute mp2 gradient for first monomer
add,-1                !subtract from previous gradient

text, HF2 CP MONOMER

dummy,f1,h1;          !first hf is now dummy
hf;                   !scf for second monomer
mp2;                  !mp2 for second monomer
ehf2=energy           !save mp2 energy in variable
forces;               !compute mp2 gradient for second monomer
add,-1                !subtract from previous gradient

dummy                !reset dummies

text, DIMER CALCULATION
hf;                   !scf for dimer
mp2;                  !mp2 for dimer
edimer=energy         !save mp2 energy in variable
forces;               !compute mp2 gradient for dimer

```

In the next example the monomer structures are kept fixed, and the interaction energy is optimized.

```
***,HF dimer mp2/CP optimization with fixed monomers
```

```
basis=avdz
```

```
maxit=20                !max number of iterations
```

```
text,OPTIMIZED VALUES OF GEOMETRY VARIABLES
```

```
RFF=      5.31431160
R1=      1.75768738
R2 =     1.75298524
THETA1 =  7.03780227
THETA2 = 111.25930975
```

```
geometry={x;noorient    !noorient must be specified since gradients are added
          f1
          f2 f1 rff
          h1 f1 1.75768738 f2 theta1
          h2 f2 1.75298524 f1 theta2 h1 180.}
```

```
text, CALCULATION AT LARGE SEPARATION
```

```
rff_save=rff           !save current rff distance
rff=1000               !dimer calculation at large separation
```

```
text, HF1 MONOMER
```

```
dummy,f2,h2;          !second hf is now dummy
hf;                   !scf for first monomer
mp2;                  !mp2 for first monomer
ehflinf=energy        !save mp2 energy in variable
```

```
text, HF2 MONOMER
```

```
dummy,f1,h1;          !first hf is now dummy
hf;                   !scf for second monomer
mp2;                  !mp2 for second monomer
ehf2inf=energy        !save mp2 energy in variable
```

```
rff=rff_save          !reset HF - HF distance to initial value
```

```
do iter=1,maxit      !optimization loop
```

```
text, HF1 CP MONOMER
```

```
dummy,f2,h2;          !second hf is now dummy
hf;                   !scf for first monomer
mp2;                  !mp2 for first monomer
ehf1=energy           !save mp2 energy in variable
forces;               !compute mp2 gradient for first monomer
scale,-1              !multiply gradient by -1
```

examples/
hfdimer_cpcopt2.com

```
text, HF2 CP MONOMER
```

```
dummy,f1,h1;          !first hf is now dummy
hf;                   !scf for second monomer
mp2;                  !mp2 for second monomer
ehf2=energy           !save mp2 energy in variable
forces;               !compute mp2 gradient for second monomer
add,-1                !subtract from previous gradient
```

```
dummy                !reset dummies
```

```
text, DIMER CALCULATION
```

```
hf;                   !scf for dimer
mp2;                  !mp2 for dimer
edimer=energy         !save mp2 energy in variable
forces;               !compute mp2 gradient for dimer
add,1                 !add to previous gradient
```

32.2.20 Printing options (PRINT)

PRINT,*code=level,...*;

Enables printing options. Usually level should be omitted or 0; values of *level* > 0 produce output useful only for debugging. *code* can be

HESSIAN	prints the updated hessian matrix. Note that its diagonal elements are printed anyway.
HISTORY	prints the complete set of previous geometries, gradients and energies.
GRADIENT	prints extended gradient information
OPT	prints detailed information about the optimization process (mainly for debugging).

Several print options can be specified with one PRINT command.

32.2.21 Conical Intersection optimization (CONICAL)

To optimize a Conical Intersection (CI) between two electronic states having the same spin, three vectors must be evaluated at SA-CPMCSCF level:

- 1) Non-Adiabatic Derivative Coupling (DC).
- 2) Gradient of the lower state (LSG).
- 3) Gradient of the upper state (USG).

NOTE: Previous versions required an explicit Gradient Difference calculation which is now replaced (point 2) by a second gradient calculation.

This can be done by adding three different CPMCSCF cards in the MULTI input

CPMCSCF , NACM, S_i, S_j , ACCU=1.0d-7, record=*record1.file*

CPMCSCF , GRAD, S_i , SPIN=*Spin of state S_i* , ACCU=1.0d-7, record=*record2.file*

CPMCSCF , GRAD, S_j , SPIN=*Spin of state S_j* , ACCU=1.0d-7, record=*record3.file*

where S_i, S_j are the electronic states in the usual format *istate.istsym* and *record[n].file* specifies the name and the file number where CPMCSCF infos should be stored. Parameter SPIN is half of the value in the WF card used to define the electronic state.

One must remember to:

- i) specify always three different *record.file* in the CPMCSCF cards.
- ii) evaluate the CPMCSCF for USG always as last.
- iii) skip the DC evaluation if the CI involves states with different spin (eg a Singlet/Triplet crossing) because the vector would be identically zero.

Three sets of FORCE cards (only two for Singlet/Triplet CI) follow the MULTI input. They will be like:

FORCE

SAMC,*record[n].file*

CONICAL,*record4.file*[,NODC]

where *record.file* is one of the records containing CPMCSCF info (again, FORCE card that evaluates the USG must be the last one) and *record4.file* points to a free record used for internal storage by the CONICAL code.

record4.file must be the same in all the CONICAL cards.

WARNING: The present implementation works properly only if *file=1* in the CONICAL cards.

The optional keyword NODC is used in case of different spins (eg S/T crossing) when DC is not evaluated and therefore not used.

The actual optimization is performed by repeatedly calling OPT inside a DO-LOOP cycle until the variable OPTCONV is below a predefined threshold. The example below optimizes the CI D0/D1 in LiH_2 (ground and excited states are both Doublets).

```

***, H2o CI
memory,3,M
basis=sto-3g
geometry={nosym
Li;h1,Li,r;h2,Li,r,h1,theta}
r=3.31510281
theta=30.57744006

maxstep=20
do i=1,maxstep
If(I.eq.1) then
int,cart
pri,2
hf
wf,4,1,0
else
int,cart
end if
multi
occ,7
core,0,,
closed,0
wf,5,1,1
state,2
CPMCSCF,NACM,1.1,2.1,accu=1.0d-7,record=5100.1
CPMCSCF,GRAD,1.1,spin=0.5,accu=1.0d-7,record=5101.1
CPMCSCF,GRAD,2.1,spin=0.5,accu=1.0d-7,record=5102.1
orbprint,5
NATORB,6666,ci,,,0

Force
SAMC,5100.1
CONICAL,6100.1

Force
SAMC,5101.1
CONICAL,6100.1

Force
SAMC,5102.1
CONICAL,6100.1

opt
STEP,0.15,0.2,2
COORD,bmat
UPDATE,bfgs=5

if(optconv.lt.1.d-4) exit

enddo

```

examples/
lih2_D0D1.com

This second example optimizes the CI S0/T0 in $LiH_2(+)$ (ground state is Singlet, excited state is Triplet).


```
***, H2o CI
memory,3,M
basis=sto-3g
geometry={nosym
Li;h1,Li,r;h2,Li,r,h1,theta}
r=3.31510281
theta=30.57744006

maxstep=40
do i=1,maxstep
If(I.eq.1) then
int,cart
pri,2
hf
wf,4,1,0
else
int,cart
end if
multi
occ,7
core,0,,
closed,0
wf,4,1,0
state,1
wf,4,1,2
state,1
CPMCSF,GRAD,1.1,spin=0,accu=1.0d-7,record=5101.1
CPMCSF,GRAD,1.1,spin=1,accu=1.0d-7,record=5100.1
orbprint,5
NATORB,6666,ci,,,0

Force
SAMC,5101.1
CONICAL,6100.1,NODC

Force
SAMC,5100.1
CONICAL,6100.1,NODC

opt,gradient=1.d-5,step=1.d-4
!STEP,0.15,0.2,2
COORD,bmat
UPDATE,bfgs=5

if(optconv.lt.1.d-4) exit

enddo
```

examples/
lih2+_SOT0.com

32.3 Examples

32.3.1 Allene Z-matrix

```

***, Allene geometry optimization using Z-Matrix
memory,1,m
basis=sto-3g

rcc=1.32 ang
rch=1.08 ang
acc=120 degree
Geometry={C1                                !Z-matrix input
          C2,c1,rcc
          Q1,c1,rcc,c2,45
          C3,c2,rcc,c1,180,q1,0
          h1,c1,rch,c2,acc,q1,0
          h2,c1,rch,c2,acc,h1,180
          h3,c3,rch,c2,acc,h1,90
          h4,c3,rch,c2,acc,h2,90}

hf
optg                                !default optimization using model hessian

```

examples/
allene_optscf.com

Results:

ITER.	ENERGY(OLD)	ENERGY(NEW)	DE	GRADMAX
1	-114.41781970	-114.42128294	-0.00346324	0.10090139
2	-114.42128294	-114.42170099	-0.00041806	0.02723890
3	-114.42170099	-114.42171728	-0.00001629	0.00670578
4	-114.42171728	-114.42171910	-0.00000182	0.00220625
5	-114.42171910	-114.42171910	0.00000000	0.00000895

32.3.2 Allene in natural internal coordinates

```

***, Allene geometry optimization using natural internal coordinates
memory,1,m
basis=sto-3g

rcc=1.32 ang
rch=1.08 ang
acc=120 degree
Geometry={nosym;
          C1;                                !Z-matrix input
          C2,c1,rcc
          Q1,c1,rcc,c2,45
          C3,c2,rcc,c1,180,q1,0
          h1,c1,rch,c2,acc,q1,0
          h2,c1,rch,c2,acc,h1,180
          h3,c3,rch,c2,acc,h1,90
          h4,c3,rch,c2,acc,h2,90}

hf;
optg                                !default optimization using model hessian
coord,bmat                          !use natural internal coordinates

```

examples/
allene_opt_bmat.com

Results:

ITER.	ENERGY(OLD)	ENERGY(NEW)	DE	GRADMAX
1	-114.41781970	-114.42134430	-0.00352460	0.06107869
2	-114.42134430	-114.42168324	-0.00033894	0.00464120
3	-114.42168324	-114.42171906	-0.00003582	0.00181807
4	-114.42171906	-114.42171910	-0.00000004	0.00020759

32.3.3 Allene MP2 optimization

```

***, Allene geometry optimization using Z-Matrix
memory,2,m
basis=vdz

rcc=1.32 ang
rch=1.08 ang
acc=120 degree
Geometry={C1                                !Z-matrix input
          C2,c1,rcc
          Q1,c1,rcc,c2,45
          C3,c2,rcc,c1,180,q1,0
          h1,c1,rch,c2,acc,q1,0
          h2,c1,rch,c2,acc,h1,180
          h3,c3,rch,c2,acc,h1,90
          h4,c3,rch,c2,acc,h2,90}

optmp2      !use default procedure optmp2

```

examples/
allene_optmp2.com

32.3.4 Caffeine XYZ

```

***, CAFFEINE cartesian coordinates (XYZ format)
memory,1,m
basis=sto-3g
geomtyp=xyz
geometry={nosym
24
          CAFFEINE CARTESIAN COORDINATES
C          0.8423320060          -0.3654865620          0.0000000000
C          -0.2841017540          -1.1961236000          0.0000000000
N          2.0294818880          -1.1042264700          0.0000000000
N          0.0774743850          -2.5357317920          0.0000000000
N          -1.6472646000          -0.6177952290          0.0000000000
C          1.4531962870          -2.3678913120          0.0000000000
C          0.6373131870          1.1735112670          0.0000000000
C          -1.7812691930          0.7688916330          0.0000000000
N          -0.6771444680          1.6306355000          0.0000000000
O          1.6106752160          1.9349693060          0.0000000000
O          -2.9202890400          1.2510058880          0.0000000000
C          -0.9202462430          3.1094501020          0.0000000000
C          -2.8623938560          -1.4824503660          0.0000000000
C          3.4552156930          -0.6811094280          0.0000000000
H          2.0878150460          -3.2451913360          0.0000000000
H          -1.4989252090          3.4222116470          -0.8897886280
H          -1.4989252090          3.4222116470          0.8897886280
H          0.0071905670          3.7148499490          0.0000000000
H          -3.4903070930          -1.2888938190          -0.8907763360
H          -3.4903070930          -1.2888938190          0.8907763360
H          -2.6289534570          -2.5638654230          0.0000000000
H          4.1360211370          -1.5529079440          0.0000000000
H          3.6817059520          -0.0685850980          0.8931597470
H          3.6817059520          -0.0685850980          -0.8931597470
}

hf
optg
coord,bmat      !Optimization in natural internal coordinates
method,diis     !Optimization method: Geometry DIIS

```

examples/
caffeine_opt_diis.com

Results:

ITER.	ENERGY(OLD)	ENERGY(NEW)	DE	GRADMAX
1	-667.68596573	-667.72289939	-0.03693366	0.08872281
2	-667.72289939	-667.73501326	-0.01211387	0.04342539
3	-667.73501326	-667.73561597	-0.00060271	0.01213581
4	-667.73561597	-667.73564985	-0.00003388	0.00342477
5	-667.73564985	-667.73565369	-0.00000384	0.00098604
6	-667.73565369	-667.73565399	-0.00000029	0.00029103

32.3.5 Transition State of Bicyclo[1.1.0]butane ring opening

```

***, Bicyclo-[1.1.0]-butane Transition State
memory,1,m
basis=3-21G

```

```

L1=1.495  ang          ! Define Active Variables
L2=1.418  ang
L3=1.463  ang
L4=1.093  ang
L5=1.111  ang
L6=1.098  ang
L7=1.097  ang
L8=1.110  ang
L9=1.106  ang
A1=92.1   degree
A2=62.1   degree
A3=136.0  degree
A4=123.5  degree
A5=122.4  degree
A6=124.7  degree
A7=126.7  degree
A8=117.9  degree
D1=-120.4 degree
D2=4.4    degree
D3=108.8  degree
D4=-107.5 degree
D5=84.2   degree
D6=109.3  degree
D7=-106.1 degree

```

examples/
butane_opt_transition.com

```

geometry={C1          ! Geometry Specification Z-Matrix
          C2  1  L1
          C3  2  L2  1  A1
          C4  1  L3  2  A2  3  D1
          H5  1  L4  2  A3  3  D2
          H6  2  L5  1  A4  4  D3
          H7  3  L6  2  A5  1  D4
          H8  3  L7  2  A6  1  D5
          H9  4  L8  1  A7  2  D6
          H10 4  L9  1  A8  2  D7}

int
rhf
optg
root,2          ! Transition State search
method,qsd     ! Use Quadratic Steepest Descent Method

```

Results:

ITER.	ENERGY(OLD)	ENERGY(NEW)	DE	GRADMAX
-------	-------------	-------------	----	---------

```

1 -153.88760305 -153.90344614 -0.01584309 0.05684518
2 -153.90344614 -153.90480163 -0.00135549 0.01484734
3 -153.90480163 -153.90492606 -0.00012443 0.00659821
4 -153.90492606 -153.90492110 0.00000497 0.00428601
5 -153.90492110 -153.90493986 -0.00001876 0.00329871
6 -153.90493986 -153.90494020 -0.00000034 0.00030504
7 -153.90494020 -153.90494021 -0.00000001 0.00015487

```

32.3.6 Reaction path of the HCN – HNC isomerization

```

***, HCN <---> NHC Isomerization Reaction Path
memory,1,m
basis=3-21G

l1=1.18282 ang      ! Starting geometry is transition state
l2=1.40745 ang
a1=55.05 degree

geometry={x;          ! Cs Symmetry
         C
         N,1,l1
         H,2,l2,1,a1}

int
rhf
optg          ! Enter geometry optimization
method,qsdpath ! Reaction path following
option,idir=1 ! First step is along the transition vector
print,history ! Print optimization history

```

examples/
hcn_isomerization.com

Results: The minimum reached is the HCN molecule

OPTIMIZATION HISTORY:

```

ENERGIES -92.246043 -92.246064 -92.247536 -92.249729 -92.254961 -92.260937 -92.269447
          -92.276316 -92.284467 -92.293352 -92.302671 -92.314152 -92.322784 -92.337223
          -92.343620 -92.349792 -92.353062 -92.354060 -92.354083 -92.354084

```

Positions

	1	2	3	4	5
L1	1.1828200	1.1818752	1.1760568	1.1711937	1.1636706
L2	1.4074500	1.4155787	1.4753427	1.5158853	1.5757858
A1	55.0500000	54.4419262	50.1407672	47.8796949	43.9463232
	6	7	8	9	10
L1	1.1560908	1.1498119	1.1391517	1.1385201	1.1324402
L2	1.6297641	1.6853237	1.7427947	1.7751395	1.8259215
A1	41.5474485	36.9015764	36.6753140	32.4681651	30.1961654
	11	12	13	14	15
L1	1.1293530	1.1232020	1.1236534	1.1164836	1.1212118
L2	1.8727459	1.9341224	1.9767135	2.0569670	2.0879631
A1	26.8338853	23.3625506	19.9713780	14.2222411	10.5176033
	16	17	18	19	20
L1	1.1227807	1.1279904	1.1355902	1.1374274	1.1371834
L2	2.1282973	2.1591890	2.1884619	2.1868721	2.1873023
A1	6.3306094	2.8192878	-0.4100429	0.0400977	-0.0160280

Reaction path following using option,idir=-1 (First step in the opposite direction)

Results: The minimum reached is the HNC molecule

OPTIMIZATION HISTORY:

ENERGIES	-92.246043	-92.246064	-92.247544	-92.249550	-92.252805	-92.257800	-92.264640
	-92.273233	-92.282800	-92.293268	-92.303710	-92.311460	-92.319520	-92.327790
	-92.331223	-92.334149	-92.336380	-92.337934	-92.339027	-92.339543	-92.339713
	-92.339713						

Positions

	1	2	3	4	5
L1	1.1828200	1.1834846	1.1893895	1.1920069	1.1947440
L2	1.4074500	1.3993195	1.3377010	1.3025012	1.2644292
A1	55.0500000	55.6613279	60.2156762	63.4020361	67.3580592
	6	7	8	9	10
L1	1.1976978	1.2003447	1.2020825	1.2019915	1.1983793
L2	1.2228083	1.1809654	1.1408173	1.1047430	1.0727752
A1	72.3210587	78.2505762	85.1291242	92.6038298	101.1239958
	11	12	13	14	15
L1	1.1936923	1.1870887	1.1795205	1.1706286	1.1673774
L2	1.0422291	1.0249059	1.0084811	0.9953467	0.9917554
A1	110.3660334	118.2669557	127.8629833	140.1807944	146.5637757
	16	17	18	19	20
L1	1.1645479	1.1623854	1.1611436	1.1599052	1.1595435
L2	0.9885425	0.9862022	0.9847864	0.9834921	0.9830566
A1	153.0819778	159.2729194	164.9184482	170.6662651	175.3627943
	21	22			
L1	1.1594176	1.1596834			
L2	0.9828926	0.9831460			
A1	179.8343282	179.9969372			

33 VIBRATIONAL FREQUENCIES (FREQUENCIES)

FREQUENCIES,*method*,SYMM=*flag*,START=*rec.ifil*,DUMP=*dumprec.ifil*;

Calculate harmonic vibrational frequencies and normal modes. To get reasonable results it is necessary to do a geometry optimization before using the frequency calculation. This option uses a hessian matrix calculated numerically from $3N$ cartesian coordinates. Z-Matrix coordinates will be destroyed on this entry. The hessian is calculated analytically or numerically by finite differences from the input coordinates. In numerical differentiation, if analytic gradients are available, these are differentiated once to build the hessian, otherwise the energy is differentiated twice. Using numerical differentiation the dipole derivatives and the IR intensities are also calculated. Note that numerical Hessians cannot be computed when dummy atoms holding basis functions are present.

The accuracy of the hessian is determined by *method*, which can be one of the following :

ANALYTICAL	use analytical second derivatives of the energy. At present, analytical second derivatives are only possible for closed shell Hartree-Fock (HF) and MCSCF wavefunctions without symmetry. It is not yet possible to calculate IR-intensities analytically. Note that, due to technical reasons, the analytical MCSCF second derivatives have to be computed in the MCSCF-program using e.g. <code>multi; cpmcscf, hess</code> (see MULTI) before they can be used in FREQUENCIES. If analytical MCSCF second derivatives are available, FREQUENCIES will use them by default.
CENTRAL	use central differences/high quality force constants (default).
NUMERICAL	differentiate the energy twice, using central differences.
FORWARD	use forward differences/low quality force constants.

During the numerical calculation of the hessian, the symmetry of the molecule may be lowered. Giving SYMM=AUTO the program uses the maximum possible symmetry of the molecular wavefunction in each energy/gradient calculation, and this option therefore minimizes the computational effort. With SYMM=NO no symmetry is used during the frequency calculation (default). For single reference calculations like HF, MP2, CCSD, RCCSD the AUTO option can be safely used and is recommended. However, it should be noted that SYMM=AUTO cannot be used for MRCI calculations, since the MRCI energy is slightly different with and without symmetry (this is due to first-order interacting space restrictions and can be avoided using REF cards, see section 18.3). Furthermore, certain input, which depends on orbital occupations or symmetry labels, cannot be used in frequency calculations with symmetry: for instance, the use of RESTRICT, SELECT, REF, PROJECT, LOCAL, state-averaged MCSCF will lead on an error unless the calculation is performed in C_1 symmetry (NOSYM option in the geometry input).

If the energy second derivatives of a given wavefunction have been calculated numerically or analytically in a previous FREQUENCIES run, the frequency calculation can be restarted from a given frequencies-record *irec* on file *ifil* using the command FREQUENCIES,START=*irec.ifil*; If no *irec.ifil* is given, information is recovered from the latest FREQUENCIES calculation. By default frequency information is saved in record 5300 on file 2. After completion of the frequency calculation, the normal modes and frequencies are dumped to record 5400 on file 2. This default record can be changed using the DUMP option. The normal modes stored in this record can be visualized using MOLDEN (see PUT command, section 9.4). By default, imaginary and low frequency modes are not stored. By specifying DUMPALL rather than DUMP all modes are written out.

By default, all computed frequencies (including low and imaginary ones) are printed. The following options can be used to modify the print level

PRINT, HESSIAN print the force constant matrix (hessian) i.e. the second derivative matrix of the energy and the mass weighted hessian matrix.

PRINT, LOW print low vibrational frequencies (i.e. the 5 or 6 frequencies belonging to rotations and translations) and their normal modes (default; PRINT, LOW=-1 suppresses the print).

PRINT, IMAG print imaginary vibrational frequencies and their normal modes (default; PRINT, IMAG=-1 suppresses the print). Imaginary frequencies appear at transition states. The normal mode of an imaginary frequency represents the transition vector of that state.

The threshold for low vibrations (default 150 cm^{-1}) can be changed using

THRESH, LOW=*value*

where *value* is the threshold in cm^{-1} .

Other subcommands of FREQUENCIES are:

STEP,*rstep* determines the step size of the numerical differentiation of the energy. Default step size *rstep*=0.001 [bohr].

NOPROJECT don't project translations and rotations out of the hessian.

SAVE,*irec.ifil* Save information of numerical frequency calculation to record *irec*. By default frequencies are saved on record 5300.2.

START,*irec.ifil* Restart numerical frequency calculation from record *irec* on file *ifil* (usually the *.wfu*-file 2).

VARIABLE,*variable* Name of a variable for which the hessian is computed

COORD=UNIQUE Use symmetry-unique displacements in the numerical calculation of the hessian (default).

COORD=3N Don't use symmetry-unique displacements (not recommended). using finite differences.

33.1 Numerical hessian using energy variables (VARIABLE)

VARIABLE,*name*;

Defines a variable *name* which holds the energy value to be used for computing the hessian using finite differences. By default, this is ENERGY(1) as set by the most recent program. For other other variables which can be used see section 32.2.18. Note that numerical Hessians cannot be computed when dummy atoms holding basis functions are present.

33.2 Thermodynamical properties (THERMO)

It is also possible to calculate the thermodynamical properties of the molecule. Since MOLPRO can only handle Abelian point groups it is necessary to give the point group of the molecule in the input file:

THERMO, SYM=*pointgroup*

pointgroup has to be the Schoenflies Symbol (e.g. C_{3v} for ammonia; linear molecules have to be C*_v or D*h respectively). If no point group card is given, rotational degeneracy will be set to 1, eventually causing deviations in the rotational entropy. If no other input card is given the

zero-point vibrational energy and the enthalpy $H(t) - H(0)$ [kJ/mol], heat capacity C_v [J/mol K] and entropy S [J/mol K] are calculated for standard Temperature and Pressure ($T = 298.150$ [K], $p = 1$ [atm]).

Subcommands of THERMO are

PRINT, THERMO additional information (such as atomic masses, partition functions and thermodynamical function in calories) is printed to the output.

SCALE, *factor* in calculating the thermodynamical properties use vibrational frequencies scaled with *factor*, in order to take account of systematic errors of the wavefunction (e.g. using SCF wavefunctions *factor*=0.89 is reasonable).

TEMP, *tmin, tmax, tstep* calculate the thermodynamical properties at different temperatures, starting with *tmin* [K] up to *tmax* [K] in steps of *tstep* [K].

PRESSURE, *p* calculate the thermodynamical properties at a given pressure of *p* [atm].

The FREQUENCIES program sets the variable *zpe* containing the zero-point-energy of the harmonic vibrations in atomic units. If the THERMO option is used, the variables *htotal* and *gtotal*, containing the enthalpy and the free enthalpy of the system in atomic units, are also set.

33.3 Examples

```
***,formaldehyde frequency calculation
memory, 8, m
```

```
basis=vdz
gthresh, energy=1.d-8
```

```
geomtyp=xyz
geometry={nosym;
  4
```

```
FORMALDEHYDE
C      0.0000000000      0.0000000000     -0.5265526741
O      0.0000000000      0.0000000000      0.6555124750
H      0.0000000000     -0.9325664988     -1.1133424527
H      0.0000000000      0.9325664988     -1.1133424527
}
```

examples/
form_freq.com

```
hf; accu, 14
optg; coord, 3n;
```

```
frequencies, analytic
thermo, sym=c2v
print, thermo
```

```
mp2
optg; coord, 3n
frequencies
thermo, sym=c2v
print, thermo
```

***, Phosphorous-pentafluoride Vibrational Frequencies

memory,1,m

basis=3-21G

geomtyp=xyz ! use cartesian coordinates xmol style

geometry={nosym; ! geometry input; don't use symmetry

6

PF5

P	0.00000	0.00000	0.00000
F	0.00000	1.11100	-1.12400
F	0.00000	-1.52800	-0.40100
F	0.00000	0.41700	1.52500
F	-1.60400	0.00000	0.00000
F	1.60400	0.00000	0.00000}

examples/
pf5_freq.com

rhf

optg ! optimize geometry

frequencies ! calculate vibrational frequencies

print,low ! print frequencies+modes of zero frequencies

thermo,sym=d3h ! calculate thermodynamical properties

temp,200,400,50 ! temperature range 200 - 400 [K]

34 ORBITAL MERGING

Orbitals can be manipulated using the MERGE facility. For instance, this allows the construction of molecular orbitals from atomic orbitals, to merge and orthogonalize different orbital sets, or to perform 2×2 rotations between individual orbitals. Other orbital manipulations can be performed using the LOCALI program (see section 16) or the MATROP program (section 35).

The merge program is called using

```
MERGE [,namout.file]
```

All subcommands described in the following sections may be abbreviated by three characters. *namout.file* specifies the output data set (see also SAVE command). If *namout.file* is omitted and no SAVE card is present, the new orbitals are not saved. All output orbitals must be supplied via ORBITAL and ADD, MOVE, EXTRA, or PROJECT directives before they can be saved.

34.1 Defining the input orbitals (ORBITAL)

```
ORBITAL,namin.file,specifications
```

Reads an input orbital set from a dump record. *specifications* can be used to select specific orbital sets, as described in section 2.16. Subsets of these orbitals can be added to the output set by the ADD, MOVE, or EXTRA commands.

34.2 Moving orbitals to the output set (MOVE)

```
MOVE,orb1.sym1,orb2.sym2,orb3.sym3,ioff,fac,istart,iend
```

Moves orbitals *orb1.sym1* to *orb2.sym2* from the input set to the first vector of symmetry *sym3* in the output set which is undefined so far. The first *orb3-1* vectors in the output set are skipped regardless of whether they have been defined before or not. If *sym2 > sym1*, *sym3* will run from *sym1* to *sym2* and the input for *sym3* has no effect. If *orb1.sym1* is negative, *abs(orb1)* is the maximum number of orbitals to be moved, starting with orbital *1.sym1*, up to *orb2.sym2*. If *orb2.sym2* is negative, *abs(orb2)* is the maximum number of vectors to be moved, starting at *orb1.isym1* up to the last orbital in symmetry *sym2*.

Orbitals from the input set which have already been moved or added to the output set are generally skipped. If *orb1* and *orb2* are zero, the whole input set is moved to the output set. In this case the input and output dimensions must be identical. If *orb1* is nonzero but *orb2* is zero, *orb2* is set to the last orbital in symmetry *sym2*. If *sym2=0*, *sym2* is set to *sym1*. *ioff* is an offset in the output vector, relative to the global offset set by OFFSET directive. *fac* has no effect for move. The elements *istart* to *iend* of the input vector are moved. If *istart=0* and *iend=0*, the whole input vector is moved.

The usage of the MOVE directive is most easily understood by looking at the examples given below. See also ADD and EXTRA commands.

34.3 Adding orbitals to the output set (ADD)

```
ADD,orb1.sym1,orb2.sym2,orb3.sym3,ioff,fac,istart,iend
```

This adds orbitals *orb1.sym1* to *orb2.sym2* to the output vectors, starting at *orb3.sym3*. The input vectors are scaled by the factor *fac*. If *fac=0*, *fac* is set to 1.0. For other details see

MOVE command. Note, however, that the output vectors which have already been defined are not skipped as for MOVE.

See also MOVE and EXTRA commands.

34.4 Defining extra symmetries (EXTRA)

EXTRA,*exsym,orb1.sym1,orb2.sym2,orb3.sym3,ioff,fac,istart,iend*

Works exactly as MOVE, but only input vectors with extra symmetry *exsym* are considered. If *orb1.sym1* and *orb2.sym2* are zero, all input vectors are moved to the output set ordered according to increasing extra symmetries.

Examples:

EXTRA , 1 , -4 . 1 will move the next 4 orbitals in symmetry 1 which have extra symmetry 1. Orbitals which have been moved before are skipped.

EXTRA , 2 , 1 . 1 will move all orbitals of symmetry 1 which have extra symmetry 2. Orbitals which have been moved before are skipped.

EXTRA will move all orbitals (all symmetries) and order them according to extra symmetries.

EXTRA , 3 , 1 . 1 , 0 . 8 Will move all orbitals which have extra symmetry 3 in all symmetries. Orbitals which have been moved before are skipped.

See also ADD and MOVE commands.

34.5 Defining offsets in the output set (OFFSET)

OFFSET,*iof₁,iof₂,...,iof₈*;

Sets offsets in the output vector for symmetries 1 to 8. In subsequent MOVE or ADD commands, the input vectors are moved to the locations iof_i+1 in the output vectors. The offset for individual ADD or MOVE commands can be modified by the parameter *ioff* on these cards. This card should immediately follow the orbital directive to which it applies. Generally, this card is only needed if the dimensions of input and output vectors are not identical.

If the dimensions of the input orbital sets are smaller than the current basis dimension, the offsets are determined automatically in the following way: each time an orbital set is read in, the previous input orbital dimensions are added to the offsets. Hence, this works correctly if the orbital sets are given in the correct order and if the individual dimensions add up to the current total dimension. If this is not the case, the offsets should be specified on an OFFSET card which must follow the orbital directive.

34.6 Projecting orbitals (PROJECT)

PROJECT,*namin.file*

This command will read vectors from record *namin.file*. These vectors must have the same dimension as those of the current calculation. All orbitals defined so far by the ORBITAL, MOVE, and ADD directives are projected out of the input set. The projected orbitals are then orthonormalized and moved to the undefined output vectors. This should always yield a complete set of vectors.

34.7 Symmetric orthonormalization (ORTH)

ORTH, n_1, n_2, \dots, n_8

Symmetrically orthonormalizes the first n_i vectors in each symmetry i . These vectors must be supplied before by ORBITAL and MOVE or ADD directives.

34.8 Schmidt orthonormalization (SCHMIDT)

SCHMIDT, n_1, n_2, \dots, n_8

Schmidt orthonormalizes the first n_i vectors in each symmetry i . These vectors must be supplied before by ORBITAL and MOVE or ADD directives.

34.9 Rotating orbitals (ROTATE)

ROTATE,*iorb1.sym,iorb2,angle*

Will perform 2×2 rotation of orbitals *iorb1* and *iorb2* in symmetry *sym* by the specified *angle* (in degree). *angle=0* means to swap the orbitals (equivalent to *angle=90*) These vectors must be supplied before by ORBITAL and MOVE or ADD directives.

34.10 Initialization of a new output set (INIT)

INIT,*namout.file*

Will initialize a new output set. All previous vectors in the output set are lost unless they have been saved by a SAVE directive!

34.11 Saving the merged orbitals

SAVE,*namout.file*

Saves the current output set to record *namout.file*. The current output set must be complete and will be Schmidt orthonormalized before it is saved. If the SAVE directive is not supplied, the output vectors will be saved after all valid commands have been processed to the record specified on the MERGE card.

34.12 Printing options (PRINT)

PRINT,*iprint,ideb*

Specifies print options.

$iprint = 0$	no print
$iprint \geq 1$:	orthonormalized orbitals specified on ORTH card are printed.
$iprint \geq 2$:	orbitals are also printed before this orthonormalization.
$iprint \geq 3$:	all final vectors are printed.
$ideb \neq 0$:	the overlap matrices are printed at various stages.

34.13 Examples

34.13.1 H₂F

This example merges the orbitals of H₂ and F

```

***,example for merge
print,orbitals,basis
rh2=1.4
rhf=300.
basis=vdz
geometry={x,y;F}                                !use C2v symmetry

text,F
rhf;wf,9,1,1;occ,3,1,1;orbital,2130.2          !rhf for f-atom

text,H2
geometry={x,y;                                !use C2v symmetry
          H1,
          H2,H1,rh2}

hf;orbital,2100.2;                              !scf for h2
multi;occ,2;orbital,2101.2;                     !mcsf for h2

text,FH2
geometry{F;                                     !linear geometry for F+H2
        H1,F,rhf
        H2,H1,rh2,F,180}

merge
orbital,2130.2                                  !rhf orbitals for F-atom
move,1.1,2.1,1.1                                !move orbitals 1.1, 2.1
move,3.1,0.4,4.1;                               !move all remaining, starting at 4.1
orbital,2100.2                                  !hf orbitals for H2
move,1.1,0.4                                    !move these to free positions
save,2131.2                                      !save merged orbitals

rhf;occ,4,1,1;start,2131.2                      !rhf for F+H2
orbital,2132.2

merge
orbital,2130.2                                  !rhf orbitals for F-atom
move,1.1,2.1,1.1                                !move orbitals 1.1, 2.1
move,3.1,3.1,4.1;                               !move orbital 3.1 to 4.1
move,4.1,0.4,6.1                                !move all remaining, starting at 6.1
orbital,2101.2                                  !mcsf orbitals for H2
move,1.1,0.4                                    !move these to free positions
save,2141.2                                      !save merged orbitals

multi;occ,5,1,1;start,2141.2;                  !casscf for F+H2 using valence space

```

examples/
h2f_merge.com

34.13.2 NO

This example merges the SCF orbitals of N and O to get a full valence space for NO. In the simplest case the atomic calculations are performed in the individual separate basis sets, but using the same symmetry (C_{2v}) as the molecular calculation.

```

***,NO merge
r=2.1

geometry={x,y;n}      !N-atom, c2v symmetry

rhf;occ,3,1,1;        !rhf nitrogen
wf,7,4,3;             !4S state
orbital,2110.2        !save orbitals to record 2110 on file 2

geometry={x,y;o}

rhf;occ,3,1,1;        !rhf for oxygen
wf,8,4,2;             !3P state
orbital,2120.2        !save orbitals to record 2120 on file 2

geometry={n;o,n,r}    ! NO molecule, c2v symmetry

MERGE
ORBITAL,2110.2        ! read orbitals of N atom
MOVE,1.1,1.1          ! move 1s orbital to output vector 1.1
MOVE,2.1,2.1,3.1      ! move 2s orbital to output vector 3.1
MOVE,3.1,3.1,5.1      ! move 2pz orbital to output vector 5.1
MOVE,1.2,1.2          ! move 2px orbital to output vector 1.2
MOVE,1.3,1.3          ! move 2py orbital to output vector 1.3
MOVE,4.1,,7.1         ! move virtual orbitals of symmetry 1
MOVE,2.2,,3.2         ! move virtual orbitals of symmetry 2
MOVE,2.3,,3.3         ! move virtual orbitals of symmetry 2
MOVE,1.4              ! move virtual orbitals of symmetry 2
ORBITAL,2120.2        ! read orbitals of O atom
MOVE,1.1,0.4          ! move all oxygen orbitals into place
ROT,3.1,4.1,45;       ! rotate 2s orbitals to make bonding and antibonding
                        ! linear combinations
ROT,5.1,6.1,-45;     ! rotate 2pz orbitals to make bonding and antibonding
                        ! linear combinations
PRINT,1              ! set print option
ORTH,6,2,2           ! symmetrically orthonormalize the valence orbitals
                        ! the resulting orbitals are printed
save,2150.2          ! save merged orbitals to record 2150.2

multi;occ,6,2,2       ! perform full valence casscf for NO
wf,15,2,1            ! 2Pix state
wf,15,3,1            ! 2Piy state
start,2150.2         ! start with merged orbitals

```

examples/
no_merge1.com

One can also do the atomic calculations in the total basis set, using dummy cards. In this case the procedure is more complicated, since the union of the two orbital spaces is over-complete. The calculation can be done as follows:

- a) SCF for the total molecule, orbitals saved to 2100.2
- b) SCF for the N atom with dummy basis on the O atom, orbitals saved on 2110.2
- c) SCF for the O atom with dummy basis on the N atom, orbitals saved on 2120.2
- d) Merge the atomic SCF orbitals. Finally, obtain the virtual orbitals by projecting the merge orbitals out of the SCF orbitals for NO.

```

***,NO merge
geometry={n;o,n,r}
r=2.1

rhf;occ,5,2,1      !rhf for NO
wf,15,2,1         !2Pi state
orbital,2100.2    !save orbitals to record 2100 on file 2

dummy,o           !oxygen is dummy
rhf;occ,3,1,1;    !rhf nitrogen
wf,7,4,3;        !4S state
orbital,2110.2    !save orbitals to record 2110 on file 2

dummy,n           !nitrogen is dummy
rhf;occ,3,1,1;    !rhf for oxygen
wf,8,4,2         !3P state
orbital,2120.2    !save orbitals to record 2120 on file 2

MERGE             !call merge program

ORBITAL,2110.2    ! read orbitals of N atom
MOVE,1.1,1.1      ! move input vector 1.1 to output vector 1.1
MOVE,2.1,3.1,3.1 ! move input vectors 2.1,3.1 to output vectors
                 ! 3.1 and 4.1
MOVE,1.2,1.2      ! move input vector 1.2 to output vector 1.2
MOVE,1.3,1.3      ! move input vector 1.3 to output vector 1.3
ORBITAL,2120.2    ! read orbitals of O atom
MOVE,1.1,3.1      ! move input vectors 1.1 to 3.1 to output vectors
                 ! 2.1, 5.1, 6.1
MOVE,1.2,1.2      ! move input vector 1.2 to output vector 2.2
MOVE,1.3,1.3      ! move input vector 1.3 to output vector 2.3
ROT,3.1,5.1,45;   ! rotate 2s orbitals to make bonding and antibonding
                 ! linear combinations
ROT,4.1,6.1,-45; ! rotate 2pz orbitals to make bonding and antibonding
                 ! linear combinations
PRINT,1          ! set print option
ORTH,6,2,2       ! symmetrically orthonormalize the valence orbitals
                 ! the resulting orbitals are printed
PROJ,2100.2      ! Project valence orbitals out of scf orbitals of the
                 ! molecule and add virtual orbital set.
SAVE,2150.2      ! save merged orbitals to record 2150 on file 2

dummy           ! remove dummies
multi;occ,6,2,2 ! perform full valence casscf for NO
wf,15,2,1       ! 2Pi state
wf,15,3,1       ! 2Pi state
start,2150.2    ! start with merged orbitals

```

examples/
no_merge2.com

35 MATRIX OPERATIONS

MATROP;

MATROP performs simple matrix manipulations for matrices whose dimensions are those of the one particle basis set. To do so, first required matrices are loaded into memory using the LOAD command. To each matrix an internal *name* (an arbitrary user defined string) is assigned, by which it is referenced in further commands. After performing operations, the resulting matrices can be saved to a dump record using the SAVE directive. Numbers, e.g. traces or individual matrix elements, can be saved in variables.

code may be one of the following:

LOAD	Loads a matrix from a file
SAVE	Saves a matrix to a file
ADD	Adds matrices
TRACE	Forms the trace of a matrix or of the product of two matrices
MULT	Multiplies two matrices
TRAN	Transforms a matrix
DMO	Transforms density into MO basis
NATORB	Computes natural orbitals
DIAG	Diagonalizes a matrix
OPRD	Forms an outer product of two vectors
DENS	Forms a closed-shell density matrix
FOCK	Computes a closed-shell fock matrix
COUL	Computes a coulomb operator
EXCH	Computes an exchange operator
PRINT	Prints a matrix
PRID	Prints diagonal elements of a matrix
PRIO	Prints orbitals
ELEM	Assigns a matrix element to a variable
READ	Reads a square matrix from input
WRITE	Writes a square matrix from input
SET	Assigns a value to a variable

See the following subsections for explanations.

35.1 Calling the matrix facility (MATROP)

The program is called by the input card MATROP without further specifications.

MATROP

It can be followed by the following commands in any order, with the restriction that a maximum of 50 matrices can be handled. The first entry in each command line is a command keyword, followed by the name of the result matrix. If the specified result matrix *result* already exists, it

is overwritten, otherwise a new matrix is created. All matrices needed in the operations must have been loaded or defined before, unless otherwise stated.

If a backquote (‘) is appended to a name, the matrix is transposed.

35.2 Loading matrices (LOAD)

All matrices which are needed in any of the subsequent commands must first be loaded into memory using the LOAD command. Depending on the matrix type, the LOAD command has slightly different options. In all forms of LOAD *name* is an arbitrary string (up to 16 characters long) by which the loaded matrix is denoted in subsequent commands.

35.2.1 Loading orbitals

LOAD,*name*,ORB [,*record*] [,*specifications*]

loads an orbital coefficient matrix from the given dump record. If the record is not specified, the last dump record is used. Specific orbitals sets can be selected using the optional *specifications*, as explained in section 2.16. The keyword ORB needs not to be given if *name*=ORB.

35.2.2 Loading density matrices

LOAD,*name*,DEN [,*record*] [,*specifications*]

loads a density matrix from the given dump record. If the record is not given, the last dump record is used. Specific orbitals sets can be selected using the optional *specifications*, as explained in section 2.16. The keyword DEN needs not to be given if *name*=DEN.

35.2.3 Loading the AO overlap matrix S

LOAD,*name*,S

loads the overlap matrix in the AO basis. The keyword S needs not to be given if *name*=S.

35.2.4 Loading $S^{-1/2}$

LOAD,*name*,SMH

loads $S^{-1/2}$, where **S** is the overlap matrix in the AO basis. The keyword SMH needs not to be given if *name*=SMH.

35.2.5 Loading the one-electron hamiltonian

LOAD,*name*,H0

LOAD,*name*,H01

loads the one-electron hamiltonian in the AO basis. H01 differs from H0 by the addition of perturbations, if present (see sections 25.4.1, 25.4.2). The keyword H0 (H01) needs not to be given if *name*=H0 (H01). The nuclear energy associated to H0 or H01 is internally stored.

35.2.6 Loading the kinetic or potential energy operators

LOAD,*name*,EKIN

LOAD,*name*,EPOT

loads the individual parts of the one-electron hamiltonian in the AO basis. EPOT is summed for all atoms. The nuclear energy is associated to EPOT and internally stored. The keyword EKIN (EPOT) needs not to be given if *name*=EKIN (EPOT).

35.2.7 Loading one-electron property operators

LOAD,*name*,OPER,*opname*,[*isym*],*x,y,z*

loads one-electron operator *opname*, where *opname* is a keyword specifying the operator (a component must be given). See section 4.13 for valid keys. *isym* is the total symmetry of the operator (default 1), and *x,y,z* is the origin of the operator. If the operator is not available yet in the operator record, it is automatically computed. The nuclear value is associated internally to *name* and also stored in variable OPNUC (this variable is overwritten for each operator which is loaded, but can be copied to another variable using the SET command. Note that the electronic part of dipole and quadrupole operators are multiplied by -1.

35.2.8 Loading matrices from plain records

LOAD,*name*,TRIANG,*record*,[*isym*]

LOAD,*name*,SQUARE,*record*,[*isym*]

Loads a triangular or square matrix from a plain record (not a dump record or operator record). If *isym* is not given, 1 is assumed.

35.3 Saving matrices (SAVE)

SAVE,*name*,*record* [,*type*]

At present, *type* can be DENSITY, ORBITALS, FOCK, H0, ORBEN, OPER, TRIANG, SQUARE, or VECTOR. If *type* is not given but known from LOAD or another command, this is assumed. Orbitals, density matrices, fock matrices, and orbital energies are saved to a dump record (the same one should normally be used for all these quantities). If *type* is H0, the one-electron hamiltonian is overwritten by the current matrix and the nuclear energy is modified according to the value associated to *name*. The nuclear energy is also stored in the variable ENUC. All other matrices can be saved in triangular or square form to plain records using the TRIANG and SQUARE options, respectively (for triangular storage, the matrix is symmetrized before being stored). Eigenvectors can be saved in plain records using the VECTOR option. Only one matrix or vector can be stored in each plain record.

One-electron operators can be stored in the operator record using

SAVE,*name*,OPER [PARITY=*np*], [NUC=*opnuc*], CENTRE=*icen*],[COORD=[*x,y,z*]]

The user-defined operator *name* can then be used on subsequent EXPEC or GEXPEC cards. *np* = 1, 0, -1 for symmetric, square, antisymmetric operators, respectively (default 1). If CENTRE is specified, the operator is assumed to have its origin at the given centre, where *icen* refers to the row number of the z-matrix input. The coordinates can also be specified explicitly using COORD. By default, the coordinates of the last read operator are assumed, or otherwise zero.

If NATURAL orbitals are generated and saved in a dump record, the occupation numbers are automatically stored as well. This is convenient for later use, e.g., in MOLDEN.

35.4 Adding matrices (ADD)

ADD,*result* [,*fac1*],*mat1* [,*fac2*],*mat2*,...

calculates $result = fac1 \cdot mat1 + fac2 \cdot mat2 + \dots$

The strings *result*, *mat1*, *mat2* are internal names specifying the matrices. *mat1*, *mat2* must exist, otherwise an error occurs. If *result* does not exist, it is created.

The factors *fac1*, *fac2* are optional (may be variables). If not given, one is assumed.

The nuclear values associated to the individual matrices are added accordingly and the result is associated to *result*.

35.5 Trace of a matrix or the product of two matrices (TRACE)

TRACE,*variable*, *mat1*, [,*factor*]

Computes $variable = factor * tr(mat1)$.

TRACE,*variable*, *mat1*, *mat2*, [,*factor*]

Computes $variable = factor * trace(mat1 \cdot mat2)$.

The result of the trace operation is stored in the MOLPRO variable *variable*, which can be used in subsequent operations.

If *factor* is not given, one is assumed.

35.6 Setting variables (SET)

SET,*variable*,*value*

Assigns *value* to MOLPRO variable *variable*, where *value* can be an expression involving any number of variables or numbers. Indexing of *variable* is not possible, however.

35.7 Multiplying matrices (MULT)

MULT,*result*, *mat1*, *mat2*, [,*fac1*], [,*fac2*]

calculates $result = fac2 * result + fac1 * mat1 \cdot mat2$

The strings *result*, *mat1*, *mat2* are the internal names of the matrices. If *fac1* is not given, *fac1*=1 is assumed. If *fac2* is not given, *fac2*=0 is assumed. If a backquote (‘) is appended to *mat1* or *mat2* the corresponding matrix is transposed before the operation. If a backquote is appended to *result*, the resulting matrix is transposed.

35.8 Transforming operators (TRAN)

TRAN,*result*, *Op*, *C*

calculates $result = C(T)*Op*C$. The strings *result*, *C*, and *Op* are the internal names of the matrices. If a backquote (‘) is appended to *C* or *Op* the corresponding matrix is transposed before the operation. Thus,

TRAN,*result*, *Op*, *C‘*

computes $result = C*Op*C(T)$.

35.9 Transforming density matrices into the MO basis (DMO)

DMO,*result*, *D*, *C*

calculates $result = C(T)*S*D*S*C$. The strings *result*, *C*, and *D* are internal names.

35.10 Diagonalizing a matrix (DIAG)

DIAG,*eigvec*,*eigval*,*matrix* [,*iprint*]

Diagonalizes *matrix*. The eigenvectors and eigenvalues are stored internally with associated names *eigvec* and *eigval*, respectively (arbitrary strings of up to 16 characters). The if *iprint.gt.0*, the eigenvalues are printed. If *iprint.gt.1*, also the eigenvectors are printed.

35.11 Generating natural orbitals (NATORB)

NATORB,*name*,*dens*,*thresh*

computes natural orbitals for density matrix *dens*. Orbitals with occupation numbers greater or equal to *thresh* (default 1.d-4) are printed.

35.12 Forming an outer product of two vectors (OPRD)

OPRD,*result*,*matrix*,*orb1*,*orb2*,*factor*

Takes the column vectors *v1* and *v2* from *matrix* and adds their outer product to *result*. *v1* and *v2* must be given in the form *icol.isym*, e.g., 3.2 means the third vector in symmetry 2. The result is

$$result(a,b) = result(a,b) + factor * v1(a) * v2(b)$$

If *result* has not been used before, it is zeroed before performing the operation.

35.13 Forming a closed-shell density matrix (DENS)

DENS,*density*,*orbitals*,*iocc1*, *iocc2* ...

Forms a closed-shell density matrix *density* from the given *orbitals*. The number of occupied orbitals in each symmetry *i* must be provided in *iocc_i*.

35.14 Computing a fock matrix (FOCK)FOCK,*f,d*

computes a closed shell fock matrix using density *d*. The result is stored in *f*.

35.15 Computing a coulomb operator (COUL)COUL,*J,d*

computes a coulomb operator J(*d*) using density *d*.

35.16 Computing an exchange operator (EXCH)COUL,*K,d*

computes an exchange operator K(*d*) using density *d*.

35.17 Printing matrices (PRINT)PRINT,*name*, [*ncol*(1), *ncol*(2), ...]

prints matrix *name*. *ncol*(*isym*) is the number of columns to be printed for row symmetry *isym* (if not given, all columns are printed). For printing orbitals one can also use ORB.

35.18 Printing diagonal elements of a matrix (PRID)PRID,*name* prints the diagonal elements of matrix *name*.**35.19 Printing orbitals (PRIO)**PRIO,*name*, *n*₁, *n*₂, *n*₃, ..., *n*₈

prints orbitals *name*. The first *n*_{*i*} orbitals are printed in symmetry *i*. If *n*_{*i*} = 0, all orbitals of that symmetry are printed.

35.20 Assigning matrix elements to a variable (ELEM)ELEM,*name*,*matrix*, *col*,*row*

assigns elements (*col*,*row*) of *matrix* to variable *name*. *col* and *row* must be given in the form *number.isym*, where *number* is the row or column number in symmetry *isym*. The product of the row and column symmetries must agree with the matrix symmetry.

35.21 Reading a matrix from the input file (READ)

READ,*name*,[*type*,[*subtype*]],[*symmetry*]
values

Reads a square matrix (symmetry 1) from input. The *values* can be in free format, but their total number must be correct. Comment lines starting with '#', '*', or '!' are skipped. The matrix can be read from another ASCII file by including this into the input using the INCLUDE command (see section 2.2). *type* is a string which can be used to assign a matrix type. If appropriate, this should be any of the ones used in the LOAD command. In addition, SUBTYPE can be specified if necessary. This describes, e.g., the type of orbitals or density matrices (e.g., for natural orbitals TYPE=ORB and SUBTYPE=NATURAL). The matrix symmetry needs to be given only if it is not equal to 1.

35.22 Writing a matrix to an ASCII file (WRITE)

WRITE,*name*,[*filename* [*status*]]

Writes a matrix to an ASCII file. If *filename* is not given the matrix is written to the output file, otherwise to the specified file (*filename* is converted to lower case). If *filename*=PUNCH it is written to the current punch file.

If *status*=NEW, ERASE or em REWIND, a new file is written, otherwise as existing file is appended.

35.23 Examples

The following example shows various uses of the MATROP commands.

```

***,h2o matrop examples
geometry={o;h1,o,r;h2,o,r,h1,theta} !Z-matrix geometry input
r=1 ang !bond length
theta=104 !bond angle
hf !do scf calculation
multi
natorb
canonical
matrop
load,D_ao,DEN,2140.2 !load mcscf density matrix
load,Cnat,ORB,2140.2,natural !load mcscf natural orbitals
load,Ccan,ORB,2140.2,canonical !load mcscf canonical orbitals
load,Dscf,DEN,2100.2 !load scf density matrix
load,S !load overlap matrix

prio,Cnat,4,1,2 !prints occupied casscf orbitals

elem,d11,Dscf,1.1,1.1 !print element D(1,1)
elem,d21,Dscf,2.1,1.1 !print element D(2,1)
elem,d12,Dscf,1.1,2.1 !print element D(1,2)

tran,S_mo,s,Cnat !transform s into MO basis (same as above)
print,S_mo !print result - should be unit matrix

trace,Nao,S_mo !trace of S_MO = number of basis functions
trace,Nel,D_ao,S !form trace(DS) = number of electrons examples/  
matrop.com

mult,SC,S,Cnat !form SC=S*Cnat
tran,D_nat,D_ao,SC !transform density to natural MO (could also be done usi

prid,D_nat !print diagonal elements (occupation numbers)

dmo,D_can,D_ao,Ccan !transform D_ao to canonical MO basis. Same as above sim
add,D_neg,-1,D_can !multiply d_can by -1
diag,U,EIG,D_neg !diagonalizes density D_can
mult,Cnat1,Ccan,U !transforms canonical orbitals to natural orbitals
prio,Cnat1,4,1,2 !prints new natural orbitals

natorb,Cnat2,D_ao !make natural orbitals using MCSCF density D_ao directly
prio,Cnat2,4,1,2 !prints new natural orbitals (should be the same as above)

add,diffden,D_ao,-1,Dscf !form mcscf-scf difference density
natorb,C_diff,diffden !make natural orbitals for difference density

write,diffden,denfile !write difference density to ASCII file denfile
save,C_diff,2500.2 !store natural orbitals for difference density in dump r

```

This second example adds a quadrupole field to H0. The result is exactly the same as using the QUAD command. H0 is overwritten by the modified one-electron matrix, and the nuclear energy is automatically changed appropriately. The subsequent SCF calculations use the modified one-electron operator.

Note that it is usually recommended to add fields with the \hat{DIP} , QUAD, or FIELD commands.


```

memory,2,m
R = 0.96488518 ANG
THETA= 101.90140469
geometry={H1
          O,H1,R;
          H2,O,R,H1,THETA}
hf;wf,10,1;

field=0.05          !define field strength

matrop
load,h0,h0          !load one-electron hamiltonian
load,xx,oper,xx    !load second moments
load,yy,oper,yy
load,zz,oper,zz
add,h01,h0,field,zz,-0.5*field,xx,-0.5*field,yy !add second moments to h0 and store in h01
save,h01,1210.1,h0 !save h0
hf                 !do scf with modified h0

matrop
load,h0,h0          !load h0
load,qmzz,oper,qmzz !load quadrupole moment qmzz
add,h01,h0,field,qmzz !add quadrupole moment to h0 (same result as above with second moments)
save,h01,1210.1,h0 !save h0
hf                 !do scf with modified h0

quad,,,field       !add quadrupole field to h0
hf                 !do scf with modified h0 (same result as above with matrop)

field,zz,field,xx,-0.5*field,yy,-0.5*field ! (add general field; same result as above)
hf                 !do scf with modified h0 (same result as above with matrop)

field,zz,field     !same as before with separate field commands
field+,xx,-0.5*field
field+,yy,-0.5*field
hf                 !do scf with modified h0 (same result as above with matrop)

```

examples/
matropfield.com

35.24 Exercise: SCF program

Write a closed-shell SCF program for H₂O using MATROP!

Hints:

First generate a starting orbital guess by finding the eigenvectors of $h0$. Store the orbitals in a record. Basis and geometry are defined in the usual way before the first call to MATROP.

Then use a MOLPRO DO loop and call MATROP for each iteration. Save the current energy in a variable (note that the nuclear energy is stored in variable ENUC). Also, compute the dipole moment in each iteration. At the end of the iteration perform a convergence test on the energy change using the IF command. This must be done outside MATROP just before the ENDDO. At this stage, you can also store the iteration numbers, energies, and dipole moments in arrays, and print these after reaching convergence using TABLE. For the following geometry and basis set

```

geometry={o;h1,o,r;h2,o,r,h1,theta} !Z-matrix geometry input
r=1 ang !bond length
theta=104 !bond angle
basis=vdz !basis set
thresh=1.d-8 !convergence threshold

```

the result could look as follows:

SCF has converged in 24 iterations

ITER	E	DIP
1.0	-68.92227207	2.17407361
2.0	-71.31376891	-5.06209922
3.0	-73.73536433	2.10199751
4.0	-74.64753557	-1.79658706
5.0	-75.41652680	1.43669203
6.0	-75.77903293	0.17616098
7.0	-75.93094231	1.05644998
8.0	-75.98812258	0.63401784
9.0	-76.00939154	0.91637513
10.0	-76.01708679	0.76319435
11.0	-76.01988143	0.86107911
12.0	-76.02088864	0.80513445
13.0	-76.02125263	0.83990621
14.0	-76.02138387	0.81956198
15.0	-76.02143124	0.83202128
16.0	-76.02144833	0.82464809
17.0	-76.02145450	0.82912805
18.0	-76.02145672	0.82646089
19.0	-76.02145752	0.82807428
20.0	-76.02145781	0.82711046
21.0	-76.02145792	0.82769196
22.0	-76.02145796	0.82734386
23.0	-76.02145797	0.82755355
24.0	-76.02145797	0.82742787

It does not converge terribly fast, but it works!

A Installation of MOLPRO

A.1 Obtaining the distribution materials

MOLPRO is distributed to licensees on a self-service basis using the world-wide web. Those entitled to the code should obtain it from <http://www.molpro.net/distrib/>, supplying the username and password given to them. The web pages contain both source code and binaries, although not everyone is entitled to source code, and binaries are not available for every platform.

Execution of MOLPRO, whether a supplied binary or built from source, requires a valid licence key. Note that the key consists of two components, namely a list of comma-separated key=value pairs, and a password string, and these are separated by '&'. The whole key needs to be inserted in its entirety as directed below. The web pages provide a facility to have this key delivered by email to a registered licensee.

A.2 Installation of pre-built binaries

Binaries are given as RPM (see <http://www.rpm.org>) packages which are installed in the standard way. There are two different RPMs: one contains the program, and the other documentation, and either may be installed independently of the other. At present these RPMs are not relocatable, and will install under `/usr/local`.

Once the program has been installed using, for example,

```
rpm -Uhv molpro-2002.6.i386.rpm
```

some post-installation configuration is necessary, and this is accomplished by changing or adding options in the script file that runs the program, usually `/usr/local/bin/molpro`. Most importantly, a valid licence token must be given in a `-k` option. Other configuration options as described in section A.3.6 may also be specified.

A.3 Installation from source files

A.3.1 Overview

There are usually four distinct stages in installing MOLPRO from source files:

- | | |
|---------------|---|
| Configuration | A shell script that allows specification of configuration options is run, and creates a configuration file that drives subsequent installation steps. |
| Compilation | The program is compiled and linked, and other miscellaneous utilities and files, including the default options file, are built. The essential resulting components are <ol style="list-style-type: none">1. The <code>molpro</code> executable, which is a small front-end that parses options, performs housekeeping functions, and starts the one or more processes that do computation.2. The <code>molpro.exe</code> executable, which is the main back-end. For parallel computation, multiple copies of <code>molpro.exe</code> are started by a single instance of <code>molpro</code> using the appropriate system utility (e.g. <code>mpirun</code>, <code>parallel</code>, <code>poe</code>, etc.).3. The <code>molpro.rc</code> file which contains default options for <code>molpro</code> (cf. section A.3.6). |

4. The `molproi.rc` file which contains MOLPRO-script procedures.
 5. Machine-ready basis-set, and other utility, libraries.
- Validation A suite of self-checking test jobs is run to provide assurance that the code as built will run correctly.
- Installation The program can be run directly from the source tree in which it is built, but it is usually recommended to run the procedure that installs the essential components in standard system directories.

A.3.2 Prerequisites

The following are required or strongly recommended for installation from source code.

1. A Fortran 90 compiler. Fortran77-only compilers will not suffice. On most systems, the latest vendor-supplied compiler should be used. On Intel Linux, the Portland `pgf90` compiler, version 3.2 or higher, is recommended. For using this compiler it is required that the environment variable `PGI` is set and points to the root directory where the compiler is installed, e.g., `/usr/local/pgi`. MOLPRO has also been tested with the Intel Compiler `ifc`, Version 5.0.1. In order to use this compiler the environment variable `IA32ROOT` must be set, and point to the appropriate directory, normally `/opt/intel/compiler50/ia32`. On Alpha Linux, the Compaq compiler is recommended. The directory in which the compiler is located must be in your `PATH`.
2. GNU `make`, freely available from <http://www.fsf.org> and mirrors. GNU `make` must be used; most system-standard makes do not work. In order to avoid the use of a wrong `make`, and to suppress extensive output of GNU `make`, it may be useful to set an alias, e.g., `alias make='gmake -s'`.
3. The GNU `wget` utility for batch-mode http transfers, although not needed for installation, is essential for any subsequent application of patches that implement bug fixes.
4. About 200Mb disk space (strongly system-dependent; more with large-blocksize file systems, and where binary files are large) during compilation. Typically 30Mb is needed for the finally installed program.
5. One or more large scratch file systems, each containing a directory that users may write on. There are parts of the program in which demanding I/O is performed simultaneously on two different files, and it is therefore helpful to provide at least two filesystems on different physical disks if other solutions, such as striping, are not available. The directory names should be stored in the environment variables `$TMPDIR`, `$TMPDIR2`, `$TMPDIR3`,... These variables should be set before the program is installed (preferably in `.profile` or `.cshrc`), since at some stages the installation procedures will check for them (cf. section A.3.6).
6. If the program is to be built for parallel execution, the Global Arrays toolkit, version 3.1 or later is needed. This is available from <http://www.emsl.pnl.gov:2080/docs/global/ga.html>, and should be installed prior to compiling MOLPRO. In some installations, GA uses the `tcgmsg` parallel harness; on others, it sits on an existing MPI subsystem, and on others, it makes use of the native parallel subsystem (e.g., LAPI). MOLPRO can be built to use any of these, although it is not normally recommended to use MPI where other possibilities exist. For more information, see section A.3.3.

7. The source distribution of MOLPRO, which consists of a base compressed tar archive with a file name of the form `molpro.2002.6.tar.gz`, together, possibly, with one or more *module* archives with file names of the form `molpro.module.2002.6.tar.gz`. The modules contain code which is not generally distributed, or features which are not always required to install the code. An example of the former is the program developers' kit (*module=develop*); an example of the latter is the documentation (*module=doc*). The archives can be unpacked using `gunzip` and `tar`. All archives must be unpacked in the same directory. It is essential that the base archive is unpacked first, and advisable that any modules are unpacked before further installation.

Under some circumstances, MOLPRO is delivered as a single tar file with a name of the form `molpro.all.2002.6.tar`. This archive contains all necessary base and module compressed tar archives, together with a shell script `unpack` which performs the unpacking described above.

A.3.3 Configuration

Once the distribution has been unpacked, identify the root directory that was created (normally `molpro2002.6`). In the following description, all directories are given relative to this root. Having changed to the root directory, you should check that the directory containing the Fortran compiler you want to use is in your `PATH`. Then run the command

```
./configure
```

which creates the file `CONFIG`. This file contains machine-dependent parameters, such as compiler options. Normally `CONFIG` will not need changing, but you should at the least examine it, and change any configuration parameters which you deem necessary. For further information, see any comments in the `CONFIG` file.

The `configure` procedure may be given command line options, and, normally, additionally prompts for a number of parameters:

1. On certain machines it is possible to compile the program to use either 32 or 64 bit integers, and in this case `configure` may be given a command-line option `-i4` or `-i8` respectively to override the default behaviour. Generally, the 64-bit choice allows larger calculations (files larger than 2Gb, more than 16 active orbitals), but can be slower if the underlying hardware does not support 64-bit integers (e.g., some IBM RS6000 hardware). Note that if `-i4` is used large files (>2Gb) are supported on most systems, but even then the sizes of MOLPRO records are restricted to 16 Gb since the internal addressing in MOLPRO uses 32-bit integers. If `-i8` is used, the record and file sizes are effectively unlimited.
2. In the case of building for parallel execution, the option `-mpp` must be given on the command line. At present, Molpro supports several different cases: the GA library can be either built on top of `tcgmsg`, `mpi`, or `myrinet`; on the IBM SP platform, it can also be built with a GA library made with the `LAPI` target. `configure` prompts for the type (default `tcgmsg`), and then for the directory holding the associated libraries. Normally, `tcgmsg` is recommended, which is most efficient on most systems and also most easily installed. If a `myrinet` network is available, `myrinet` should be chosen. This requires in addition to the usual MPI libraries the `gm` library and `mpirun_gm` rather than `mpirun`. At present, the `myrinet` option has been tested only on Linux systems. The name of the MOLPRO executable is generated from the program version number, the library type and

the machine architecture. It is then possible to install different versions simultaneously in the same MOLPRO tree; see section A.3.4.

3. If any system libraries are in unusual places, it may be necessary to specify them explicitly as the arguments to a `-L` command-line option.
4. `configure` prompts for the licence key, obtainable as described above. The key may also be given using the `-k` option on the command line, or given through the environment variable `$MOLPRO_KEY`.
5. `configure` asks whether you wish to use system BLAS subroutine libraries. MOLPRO has its own optimized Fortran version of these libraries, and this can safely be used. On most machines, however, it will be advantageous to use a system-tuned version instead. In the case of BLAS, you should enter a number between 1, 2 and 3; if, for example, you specify 2, the system libraries will be used for level 2 and level 1 BLAS, but MOLPRO's internal routines will be used for level 3 (i.e., matrix-matrix multiplication). Normally, however, one would choose either 0 or 3. If a system BLAS is chosen, you will be prompted to enter appropriate linker options (e.g. `-L/usr/lib -lblas`) to access the libraries.

A special situation arises if 64-bit integers are in use (`-i8`), since on many platforms the system BLAS libraries only supports 32-bit integer arguments. In such cases (e.g., IBM, SGI, SUN) either 0 or 4 can be given for the BLAS level. `BLAS=0` should always work and means that the MOLPRO Fortran BLAS routines are used. On some platforms (IBM, SGI, SUN) `BLAS=4` will give better performance; in this case some 32-bit BLAS routines are used from the system library (these are then called from wrapper routines, which convert 64 to 32-bit integer arguments. Note that this might cause problems if more than 2 GB of memory is used).

For good performance it is important to use appropriate BLAS libraries; in particular, a fast implementation of the matrix multiplication `dgemm` is very important for MOLPRO. Therefore you should use a system tuned BLAS library whenever available. For Linux PCs we recommend the following BLAS libraries:

Intel PIII: `lsblaspi1.2f_03.00.a` in the ASCI library, which can be obtained from <http://www.cs.utk.edu/~ghenry/distrib/>

To use this library, link it to a file name that the linker can understand, for example, `libblas.a`

```
ln -s lsblaspi1.2f_03.00.a libblas.a
```

and when `configure` prompts you for the library, type

```
-L blasdir -lblas
```

where `blasdir` is the absolute path of the directory holding the BLAS library.

AMD Athlon: Atlas library, obtainable from <http://www.netlib.org/atlas/>. The easiest and safest is to use a pre-built library, and we found that `atlas3.2.1_Linux_ATHLON` works very well on current hardware. The appropriate linker options to provide are

```
-L blasdir -lcblas -lf77blas -latlas
```

An even faster BLAS library `libathlonblas.a` is available on the MOLPRO web page, see below.

Intel PIV: Atlas library `atlas3.3.0_Linux_P4SSE2.tgz`, otherwise as for Athlon above. Alternatively, the Intel `mk1` libraries can be used in conjunction with the Intel fortran compiler.

For the cases where copyright rules permit, these libraries, as well as BLAS libraries for other systems such as HP (PA-RISC 32- and 64-bit, and IA64, Intel mkl), can be obtained from <http://www.molpro.net/blaslibs>. Specification of these libraries can be simplified by placing any relevant downloaded libraries in the directory `blaslibs`; `configure` searches this directory (and then, with lower priority, some potential system directories) for libraries relevant to the hardware, including that specified by a `-p3`, `-p4`, `-athlon` command line option (see below). Any directory structure in the web copy of these libraries should be preserved in the local copy. The simplest way to ensure all this is to fetch complete set of libraries using

```
wget --cut-dirs=1 -nH -np -r http://www.molpro.net/blaslibs
```

6. `configure` prompts for the destination directory (`INSTBIN`) for final installation of the MOLPRO executable. This directory should be one normally in the `PATH` of all users who will access MOLPRO, and its specification will depend on whether the installation is private or public.
7. `configure` prompts for the destination directory (`INSTLIB`) for installation of ancillary files which are required for program execution.
8. `configure` prompts for the destination directory for documentation. This should normally be a directory that is mounted on a worldwide web server.
9. `configure` prompts for the destination directory for the CGI scripts that control the delivery of documentation. This might be the same directory as (h), but some web servers require a particular special directory to be used.

The latter two parameters are relevant only if the documentation is also going to be installed from this directory (see below).

The following command-line options are recognized by `configure`.

- `-batch` disables the prompting described above.
- `-k key` specifies the licence key.
- `-i8 | -i4` forces the use of 8- or 4-byte integers respectively.
- `-L lib` specifies any additional directories containing system libraries to be scanned at link time.
- `-blas 0|1|2|3|4` specifies system BLAS level, as described above.
- `-mpp | -nompp` controls whether compilation is to be for MPP parallelism (see above).
- `-ifc | -pgf` controls whether the Intel (`ifc`) or Portland (`pgf`) compiler is be used on Linux IA32 systems. Note that appropriate environment variables must be set, see section A.3.2.
- `-f ftcflag` adds a token to the specifiers for the Fortran preprocessor `ftc`.
- `-largefiles | -nolargefiles` controls whether large file (> 2Gb) support is wanted. This option is not relevant or used on all architectures. For Linux PC, it should be specified only if the kernel and system libraries also support large files.
- `-p3 | -p4 | -athlon` specifically identifies a particular hardware in order to force appropriate run-time libraries where possible. These options are supported only on Linux systems. If any of these options is given, the MOLPRO executable will

be named `molpro_p3.exe`,
`molpro_p4.exe`, or `molpro_athlon.exe` (in the mpp case, e.g., `molpro_p3_tcgmsg`).
It is possible to install different platform variants simultaneously in the same
MOLPRO tree; see section A.3.4.

A.3.4 Configuration of multiple executables in the same MOLPRO tree

On Linux systems, it may be desirable to have optimized versions for different hardware architectures, like `p3`, `p4`, or `athlon` (see section A.3.3). Provided the compiler options are the same (i.e. neither `p4`, nor `athlon` specific), the different versions differ only by the use of specific BLAS libraries. It is then possible to install different executables for each case in the same MOLPRO tree, without the need to recompile the program. To do so, one first needs to run `configure` for each case, and specify the appropriate libraries when `configure` prompts for them. These library paths are all stored in the file `CONFIG`, generated by `configure`. Subsequently,

```
make ARCH=procname
```

will link the desired version, where *procname* can be `p3`, `p4`, or `athlon`. This will generate the executable `molpro_procname.exe`. If the `ARCH` option is not given, the last one configured will be generated.

In addition, a file `molpro_procname.rc` will be generated for each case, which defines the running environment and may also contain system dependent tuning parameters (see section A.3.7). A specific executable can then be requested using

```
molpro -rcfile molpro_procname.rc input
```

More conveniently, one can set the Unix environment variable `MOLPRO_RCFILE` to `molpro_procname.rc` and then simply use `molpro` without an option. The recommended mechanism is to set the environment variable `MOLPRO_RCFILE` in the default environment (`.cshrc`, `.profile`) as appropriate on a given machine.

Similarly, different MPP version can also be installed in one MOLPRO tree (but the tree for parallel and serial versions must be distinct!). In this case, one can run `configure` for `tcgmsg`, `mpi`, and/or `myrinet` (and in addition with `-p3`, `-p4`, and/or `-athlon`), and then link using

```
make MPPLIB=libname
```

where *libname* can be `tcgmsg`, `mpi`, or `myrinet`. The `ARCH` and `MPPLIB` options can be combined, e.g.,

```
make MPPLIB=libname ARCH=procname
```

and this will generate the executable `molpro_procname_libname.exe` and the default file `molpro_procname_libname.rc`.

As described above, the different executables can then be chosen on a specific machine by setting the environment variable `MOLPRO_RCFILE` to `molpro_procname_libname.rc`.

Note that if `MOLPRO_RCFILE` is not set, `molpro.rc` will be used by default, which will correspond to the last `molpro_procname_libname.rc` generated.

A.3.5 Compilation and linking

After configuration, the remainder of the installation is accomplished using the GNU *make* command. Remember that the default *make* on many systems will not work, and that it is essential to

use GNU *make* (cf. section A.3.2). Everything needed to make a functioning program together with all ancillary files is carried out by default simply by issuing the command

```
make
```

in the MOLPRO base directory. Most of the standard options for GNU *make* can be used safely; in particular, `-j` can be used to speed up compilation on a parallel machine. The program can then be accessed by making sure the `bin/` directory is included in the `PATH` and issuing the command `molpro`.

A.3.6 Adjusting the default environment for MOLPRO

The default running options for MOLPRO are stored in the file `bin/molpro.rc`. After program installation, either using RPMs or from source files, this file should be reviewed and adjusted, if necessary. Particular attention should be paid to some or all of the following (see User's manual for full discussion of options).

`-d dir1:dir2:...` where `dir1:dir2:...` is a list of directories which may be used for creating scratch files. Each of the directories should be writable by those who will use the program, and the directory specification may contain embedded environment variables in shell form, for example `$TMPDIR` or `/tmp/$USER`; these will be expanded at run time. If multiple scratch file systems are available, it is advantageous to present a list of directories of which there is one in each file system. Some parts of MOLPRO present extreme I/O demands, and it is therefore important to be careful in optimizing the provision and specification of scratch directories.

Note that in the building of `bin/molpro.rc`, the environment variables `$TMPDIR`, `$TMPDIR2`, `$TMPDIR3`,... are used to construct the list of scratch directories for the `-d` option. Thus, these environment variables should at make time be filled with the names of directories on each available scratch file system (cf. section A.3.3).

`-I directory` This determines the destination of permanent integral files. At run time this file is located in the first directory specified after `-d`, (i.e., `dir1`, see above), but after completion of the job the file will be copied to the directory given after `-I`. Since the integral file can be very large, it is normally recommended that `directory` is identical to `dir1` (this is the default). Then no copying will take place. On some main frames, the scratch directory is erased automatically after a job has terminated, and in such cases a different `-I` directory, e.g., `$HOME/int`, can be specified (environment variables will be expanded at run time). In view of the large integral file sizes, this should be used with care, however. Note that in parallel runs with more than 1 processor the integral file will never be copied, and cannot be restarted.

`-W directory` This determines the destination of permanent wavefunction (dump) files used for storing information like orbitals or CI-vectors etc. These files are essential for restarting a job. As explained for the integral files above, permanent wavefunction files will be copied to `directory` after completion of the job. The default for `directory` is `$HOME/wfu`.

`-k key` where `key` is the licence key, obtainable as described in section A.1.

`-m, -G` The default local memory and GA memory should be checked to be appropriate for the hardware environment.

`-n, -N` The number of processors or their identity can be specified explicitly in the configuration file, but very often it is neither desirable nor necessary to do so. Where possible, the `molpro` program extracts a reasonable default for the node specification from the controlling batch system (e.g. LoadLeveler, PBS). Usually the user will want to either specify `-n` explicitly on the command line, or rely on `molpro`'s attempts to get it from the batch system.

A.3.7 Tuning

MOLPRO can be tuned for a particular system by running in the root directory the command

```
molpro tuning.com
```

This job automatically determines a number of tuning parameters and appends these to the file `bin/molpro.rc`. Using these parameters, MOLPRO will select the best BLAS routines depending on the problem size. This job should run on an empty system. It may typically take 10 minutes, depending on the processor speed, and you should wait for completion of this run before doing the next steps.

A.3.8 Testing

At this stage, it is essential to check that the program has compiled correctly. The makefile target *test* (i.e., command `make test`) will do this using the full suite of test jobs, and although this takes a significantly long time, it should always be done when porting for the first time. A much faster test, which checks the main routes through the program, can be done using `make quicktest`. For parallel installation, it is highly desirable to perform this validation with more than one running process. This can be done conveniently through the make command line as, for example,

```
make MOLPRO_OPTIONS=-n2 test
```

If any test jobs fail, the cause must be investigated. It may be helpful in such circumstances to compare the target platform with the lists of platforms on which MOLPRO is thought to function at <http://www.molpro.net/machines.html>. If, after due efforts to fix problems of a local origin, the problem cannot be resolved, the developers of MOLPRO would appreciate receiving a report. There is a web-based mechanism at <http://www.molpro.net/bug> at which as many details as possible should be filled in. `make test` produces a file of the form `testjobs/report.*.tar.gz` that contains some details of the MOLPRO installation, and the output files of the failing test jobs. You should normally attach this file to the bug report. Please note that the purpose of such bug reports is to help the developers improve the code, and not for providing advice on installation or running.

A.3.9 Installing the program for production

Although the program can be used in situ, it is usually convenient to copy only those files needed at run time into appropriate installation directories as specified at configuration time (see section A.3.3) and stored in the file `CONFIG`. To install the program in this way, do

```
make install
```

The complete source tree can then be archived and deleted. If multiple Linux executables have been generated (see section A.3.4), they can be installed using

```
make MPPLIB=libname ARCH=procname install
```

into the same INSTBIN and INSTLIB directories (but note that the INSTLIB directories must be distinct for i4 and i8 versions). The overall effect of this is to create in the INSTBIN directory an executable command file of the form *name_arch_mplib*, where *name* is one of *molpros*, *molprop*, corresponding to serial or parallel execution. If the file INSTBIN/*name* does not already exist, or if the variable DEFAULT is set during `make install` (i.e., `make DEFAULT=1 install`), then a symbolic link is made to INSTBIN/*name*. Furthermore, If the file INSTBIN/*molpro* does not already exist, or if the variable DEFAULT is set to *molpro* during `make install` then a symbolic link is made from INSTBIN/*name* to INSTBIN/*molpro*. The overall effect of this cascade of links is to provide, in the normal case, the commands *molpro* and one or both of *molpros* (serial) and *molprop* (parallel) for normal use, with the long names remaining available for explicit selection of particular variants. As with the uninstalled program, the environment variable MOLPRO_RCFILE can be used to override the choice of configuration file.

For normal single-variant installations, none of the above has to be worried about, and the *molpro* command will be available from directory INSTLIB.

When the program has been verified and/or installed, the command `make clean` can be used to remove compilation logs. `make veryclean` will remove all binary and object files, retaining only those files included in the original distribution; it is usually recommended that this is not done, as it implies that to apply future updates and bug fixes, the whole program will have to be recompiled.

A.3.10 Getting and applying patches

Normally, the distribution when downloaded is fully up to date, and initial patching is not necessary. However, bug fixes and updates may be desired subsequently. The mechanism for updating MOLPRO source code with bug fixes and new features is through the provision of self-contained patch files, which, when applied, replace or add files, and store the replaced code in order to allow later reversion to the original. Those patches that are available can be seen at <http://www.molpro.net/patch/2002.6>, whilst a list of those already installed is printed when running the program. Patch files automatically outdate any targets that need rebuilding as a result of the patch; for example, relevant object files are removed. Thus, after all patches have been applied, it is usually necessary to rebuild the program using `make`.

The order in which patches are applied and removed is important. Some patches are prerequisites of others, and some patches are ‘parents’ of one or more ‘children’: the parent and child patches have one or more files in common, but the parent is older than the child. Individual patch scripts will themselves refuse to apply or revert if rules based on these considerations would be violated. In order to deal with this issue smoothly, a program *patcher* is provided to manage the application and removal of one or more patches. *patcher* attempts to sort the order in which patches are applied or reverted so as to avoid such conflicts; it will also, if necessary, revert and reapply patches.

Note that if you need to run *patcher* before compiling MOLPRO, you can build it from the top-level directory with

```
make -C utilities patcher
```

or if you need to run *patcher* before running `configure`, with

```
make -C utilities patcher.bootstrap
```

To use the `patcher` program, in the top-level directory issue the command

```
./patcher [--apply | --revert | --list]
          [--cache-directory] [--user] [--password]
          [--url] [--local]
          [--verbose] [--no-action] patch1 patch2 ....
```

It can operate in one of three possible modes according to the options

```
--apply, -a      (default) Apply (i.e. install) patches
--revert, -r     Revert (i.e. remove) patches
--list, -l       List available and installed patches
```

The list of patches to remove or install can be given on the command line after all options as an explicit list of either patch names or, in the case of application, patch files. Alternatively and usually, for the case of application, one can through options request either all patches that are in a local cache, or all patches that are available.

The MOLPRO patches from the central web server (default `http://www.molpro.net`), are cached by this program in a local directory (default `$HOME/.molpro/cache`). Access to the web server typically has to be authenticated; the first time you run this program, you can specify your username and password through command-line options, or else the program will prompt for them. They are then remembered in the file `CONFIG` in the cache directory.

In case of problems, first consult the file `patcher.log`, which contains the output from individual patch applications and reversions.

The following options can be given.

```
--cache-directory, -c d location of cache directory.
--verbose, -v          Increase amount of information printed. Multiple --verbose options can be
                      used.
--noverbose           Decrease amount of information printed.
--url                 URL of web server.
--user, -u u          Username for web server.
--password, -p p      Password for web server.
--noaction, -n        No applications or reversions are actually done. Useful for seeing what
                      would happen without doing it.
--local              Don't attempt to access the web server, but use only local files.
```

Examples:

```
patcher
```

Applies all patches that are available, but not yet installed. This is the normal use of the utility in bringing the copy of the source tree up to date with all available updates.

```
patcher -l
```

Lists installed and available patches.

```
patcher -r xx yy
```

Reverts patches xx and yy.

```
patcher -n
```

Loads all uninstalled patches into the cache for later use.

```
patcher --local
```

Applies all patches in the cache; no network connection needed.

A.3.11 Installation of documentation

The documentation is usually supplied as a compressed *tar* file with a name of the form `molpro.doc.2002.6.tar.gz`. This file unpacks to a top-level directory `molpro2002.6`; this may either coincide or not with the top-level directory containing source code or not, according to taste. The postscript and PDF user's manual is found in the directory `molpro2002.6/doc`, with the HTML version in the directory `molpro2002.6/doc/manual` (top level file is `manual.html`). It is generally recommended that the documentation is unpacked in the installation source tree, so that the documentation can be copied to its final destination as specified in the `CONFIG` file generated by the `configure` command. To install the documentation and interactive basis set tool, issue `make install` in the `doc` directory. Numerous example input files are included in the manual, and can alternatively be seen in the directory `molpro2002.6/examples`.

B Recent Changes

B.1 New features of MOLPRO2002.6

Relative to version 2002.1, there are the following changes and additions:

1. Support for IA-64 linux systems (HP and NEC) and HP-UX 11.22 for IA-64 (Itanium2).
2. Support for NEC-SX systems.
3. Support for IBM-power4 systems.
4. Modified handling of Molpro system variables. The SET command has changed (see sections 6 and 6.3).
5. The total charge of the molecule can be specified in a variable CHARGE or on the WF card, see section 2.14.
6. Improved numerical geometry optimization using symmetrical displacement coordinates (see sections 31.2 and 32).
7. Improved numerical frequency calculations using the symmetry (AUTO option, see section 33).

B.2 New features of MOLPRO2002

Relative to version 2000.1, there are the following principal changes and additions:

1. Modules direct and local are now included in the base version. This means that integral-direct procedures as described in
M. Schütz, R. Lindh, and H.-J. Werner, *Mol. Phys.* 96, 719 (1999),
linear-scaling local MP2, as described in
G. Hetzer, P. Pulay, and H.-J. Werner, *Chem. Phys. Lett.* 290, 143 (1998),
M. Schütz, G. Hetzer, and H.-J. Werner, *J. Chem. Phys.* **111**, 5691 (1999),
G. Hetzer, M. Schütz, H. Stoll, and H.-J. Werner, *J. Chem. Phys.* **113**, 9443 (2000),
as well as LMP2 gradients as described in
A. El Azhary, G. Rauhut, P. Pulay, and H.-J. Werner, *J. Chem. Phys.* **108**, 5185 (1998)
are now available without special licence. The linear scaling LCCSD(T) methods as described in
M. Schütz and H.-J. Werner, *J. Chem. Phys.* **114**, 661 (2001),
M. Schütz and H.-J. Werner, *Chem. Phys. Lett.* **318**, 370 (2000),
M. Schütz, *J. Chem. Phys.* **113**, 9986 (2000)
will be made available at a later stage.
2. QCISD gradients as described in *Phys. Chem. Chem. Phys.* **3**, 4853 (2001) are now available.
3. Additional and more flexible options for computing numerical gradients and performing geometry optimizations.

4. A large number of additional density functionals have been added, together with support for the automated functional implementer described in *Comp. Phys. Commun.* **136** 310–318 (2001).
5. Multipole moments of arbitrary order can be computed.
6. Further modules have been parallelized, in particular the CCSD(T) and direct LMP2 codes. The parallel running procedures have been improved. The parallel version is available as an optional module.
7. The basis set library has been extended.
8. Some subtle changes in the basis set input: it is not possible any more that several one-line basis input cards with definitions for individual atoms follow each other. Each new basis card supercedes previous ones. Either all specifications must be given on *one* BASIS card, or a basis input block must be used. `BASIS,NAME` is now entirely equivalent to `BASIS=NAME`, i.e. a global default basis set is defined and the variable `BASIS` is set in both cases.
9. Pseudopotential energy calculations can now be performed with up to *i*-functions, gradients with up to *h*-functions.
10. Many internal changes have been made to make MOLPRO more modular and stable. Support has been added for recent operating systems on Compaq, HP, SGI, SUN, and Linux. The patching system has been improved.

B.3 Features that were new in MOLPRO2000

Relative to version 98.1, there are the following principal changes and additions:

1. There was a fundamental error in the derivation of the spin-restricted open-shell coupled-cluster equations in *J. Chem. Phys.* **99**, 5129 (1993) that is also reflected in the RCCSD code in MOLPRO version 98.1 and earlier. This error has now been corrected, and an erratum has been published in *J. Chem. Phys.* **112**, 3106 (2000). Fortunately, the numerical implications of the error were small, and it is not anticipated that any computed properties will have been significantly in error.
2. There was a programming error in the transformation of gradients from cartesian to internal coordinates, which in some cases resulted in slow convergence of geometry optimizations. The error is now fixed.
3. Vibrational frequencies formerly by default used average atomic masses, rather than those of the most common isotopes, which is now the default behaviour.
4. MCSCF second derivatives (author Riccardo Tarroni) added (preliminary version, only without symmetry). Frequency and geometry optimization programs are modified so that they can use the analytic hessian.
5. New internally contracted multireference second-order perturbation theory code (author Paolo Celani) through command `RS2C`, as described in P. Celani and H.-J. Werner, *J. Chem. Phys.* **112**, 5546 (2000).
6. EOM-CCSD for excited states (author Tatiana Korona).
7. QCISD dipole moments as true analytical energy derivatives (author Guntram Rauhut).

8. Linear scaling (CPU and memory) LMP2 as described by G. Hetzer, P. Pulay, and H.-J. Werner, Chem. Phys. Lett. 290, 143 (1998).
M. Schütz, G. Hetzer, and H.-J. Werner, J. Chem. Phys. **111**, 5691 (1999).
9. Improved handling of basis and geometry records. 98.1 and 99.1 dump files can be restarted, but in case of problems with restarting old files, add `RESTART,NOGEOM` immediately after the `file` card. Also, if there are unjustified messages coming up in very large cases about "ORBITALS CORRESPOND TO DIFFERENT GEOMETRY" try `ORBITAL,record,NOCHECK`. (This can happen for cases with more than 100 atoms, since the old version was limited to 100).
10. Reorganization and generalization of basis input. Increased basis library.
11. Counterpoise geometry optimizations.
12. Improved running procedures for MPP machines. Parallel direct scf and scf gradients are working. These features are only available with the MPP module, which is not yet being distributed.
13. Important bugfixes for DFT grids, CCSD with paging, finite field calculations without core orbitals, spin-orbit coupling.
14. Many other internal changes.

As an additional service to the MOLPRO community, an electronic mailing list has been set up to provide a forum for open discussion on all aspects of installing and using MOLPRO. The mailing list is intended as the primary means of disseminating hints and tips on how to use Molpro effectively. It is not a means of raising queries directly with the authors of the program. For clearly demonstrable program errors, reports should continue to be sent to `molpro-support@molpro.net`; however, 'how-to' questions sent there will merely be redirected to this mailing list.

In order to subscribe to the list, send mail to `molpro-user-request@molpro.net` containing the text `subscribe`; for help, send mail containing the text `help`.

Messages can be sent to the list (`molpro-user@molpro.net`), but this can be done only by subscribers. Previous postings can be viewed in the archive at <http://www.molpro.net/molpro-user/archive> irrespective of whether or not you subscribe to the list. Experienced Molpro users are encouraged to post responses to queries raised. Please do contribute to make this resource mutually useful.

B.4 Facilities that were new in MOLPRO98

MOLPRO98 has the full functionality of MOLPRO96, but in order to make the code more modular and easier to use and maintain, a number of structural changes have been made. In particular, the number of different records has been significantly reduced. The information for a given wavefunction type, like orbitals, density matrices, fock matrices, occupation numbers and other information, is now stored in a single dump record. Even different orbital types, e.g., canonical, natural, or localized orbitals, are stored in the same record, and the user can subsequently access individual sets by keywords on the `ORBITAL` directive. New facilities allow the use of starting orbitals computed with different basis sets and/or different symmetries for SCF or MCSCF calculations. The default starting guess for SCF calculations has been much improved, which is most useful in calculations for large molecules. The use of special procedures for computing non-adiabatic couplings or diabatization of orbitals has been significantly simplified. We hope that these changes make the program easier to use and reduce the probability of input errors.

However, in order to use the new facilities efficiently, even experienced MOLPRO users should read the sections *RECORDS* and *SELECTING ORBITALS AND DENSITY MATRICES* in the manual. It is likely that standard MOLPRO96 inputs still work, but changes may be required in more special cases involving particular records for orbitals, density matrices, or operators.

All one-electron operators needed to compute expectation values and transition quantities are now stored in a single record. Operators for which expectation values are requested can be selected globally for all programs of a given run using the global GEXPEC directive, or for a specific program using the EXPEC directive. All operators are computed automatically when needed, and the user does not have to give input for this any more. See section *ONE-ELECTRON OPERATORS AND EXPECTATION VALUES* of the manual for details.

Due to the changed structure of dump and operator records, the utility program MATROP has a new input syntax. MOLPRO96 inputs for MATROP do not work any more.

In addition to these organizational changes, a number of new programs have been added. Analytic energy gradients can now be evaluated for MP2 and DFT wavefunctions, and harmonic vibrational frequencies, intensities, and thermodynamic quantities can be computed automatically using finite differences of analytical gradients. Geometry optimization has been further improved, and new facilities for reaction path following have been added.

An interface to the graphics program MOLDEN has been added, which allows to visualize molecular structures, orbitals, electron densities, or vibrations.

Integral-direct calculations, in which the two-electron integrals in the AO basis are never stored on disk but always recomputed when needed, are now available for all kinds of wavefunctions, with the exception of perturbative triple excitations in MP4 and CCSD(T) calculations. This allows the use of significantly larger basis sets than was possible before. The direct option can be selected globally using the GDIRECT command, or for a specific program using the DIRECT directive. See section *INTEGRAL DIRECT METHODS* in the manual for details. Note that the DIRECT module is optional and not part of the basic MOLPRO distribution.

Local electron correlation methods have been further improved. In combination with the integral-direct modules, which implement efficient prescreening techniques, the scaling of the computational cost with molecular size is dramatically reduced, approaching now quadratic or even linear scaling for MP2 and higher correlation methods. This makes possible to perform correlated calculations for much larger molecules than were previously feasible. However, since these methods are subject of active current research and still under intense development, we decided not to include them in the current MOLPRO release. They will be optionally available in one of the next releases.

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