

Physics 476LW

Modern Physics Laboratory

Gamma-Ray Spectroscopy

Introduction

The purpose of this experiment is to acquaint the student with some of the basic techniques used for measuring g-rays. It is based on the use of a thallium-activated sodium iodide detector (NaI(Tl)). Much of this manual is adapted from the ORTEC[®] documents *Experiment 3: Gamma-Ray Spectroscopy Using NaI(Tl)* and *MAESTRO[®]-32 MCA Emulator Software User's Manual*.

There are a number of types of detectors for gamma radiation and they are used in a variety of areas of physics. For example, g-ray detectors that were used for studying g-rays produced by electron-positron annihilation near black holes detected g-rays produced by lightning on Earth. g-ray production by lightning had been unknown and the serendipitous measurements opened up a variety of questions that continue to be studied.

The premium solid-state g-ray detector is high purity germanium (HPGe). Although it produces excellent results it is very expensive and a high level of expertise is need use it. In this lab we use a NaI(Tl) detector, which is a scintillation type detector. One of the weaknesses of NaI(Tl) is that it does not have the energy resolution of HPGe, as can be seen by comparing the width of the peaks in Figure 1. Therefore, one part of our experiment will be to determine the energy resolution of our NaI(Tl) detector.

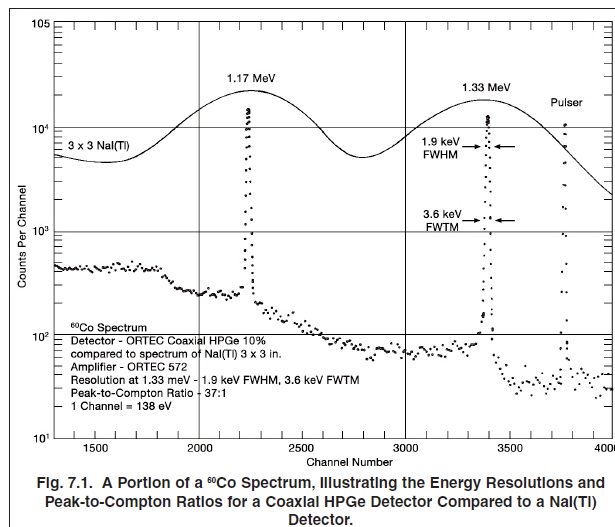


Figure 1. Comparison of peaks from HPGe and NaI(Tl) detectors.
(From ORTEC publication: *Experiment 7: High-Resolution Gamma-Ray Spectroscopy*)

2 The Experiment

2.1 The Apparatus

- Integrated NaI(Tl) Scintillator, Photomultiplier Tube, and PMT Base with Stand
- Lead shielding
- Multi Channel Analyzer (MCA)
- Windows computer with PCI card
- MAESTRO-32 software
- A set of radioactive samples

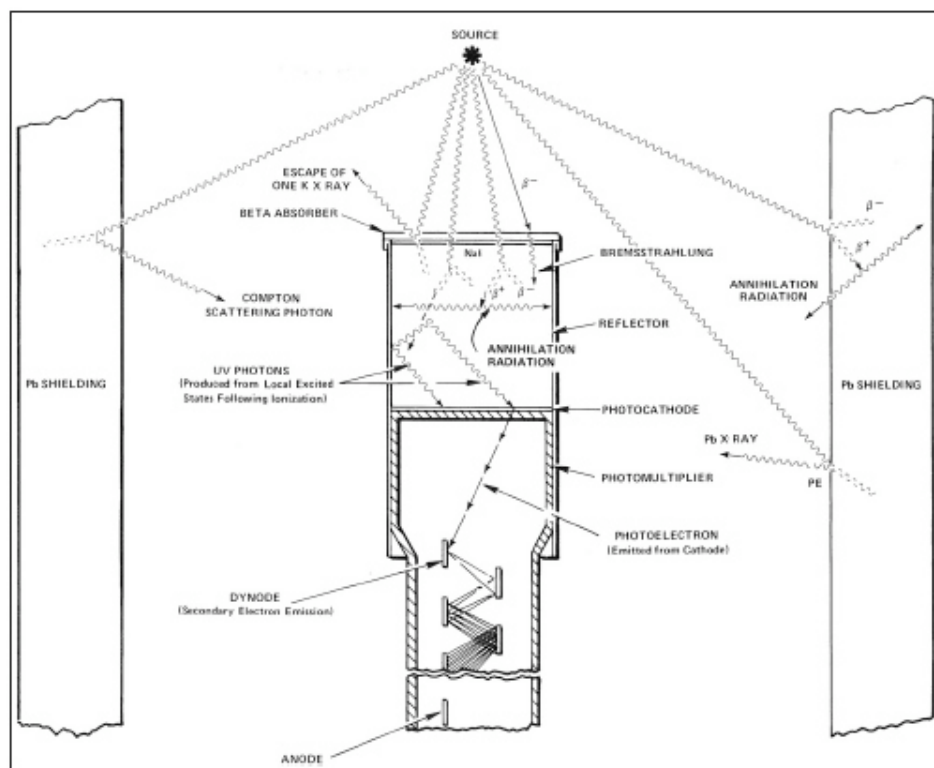


Figure 2. Diagram of NaI(Tl) g-ray scintillation detector system

2.2 Startup Procedures

The apparatus will be set up for this experiment. Complete the following steps:

- Log onto the computer using your university username and password .- Start the MAESTRO-32 software
- Under tab "Acquire" select option "MCB Properties"
- Select tab "high voltage" and select "On". The high voltage is 800 V.

2.3 Spectra Collection

Calibration

The first task that must be performed is calibration. The detector consists of 1024 channels of a multichannel analyzer (MCA) system. The system has a linear relation between the height of the pulse coming from the gamma-ray detector and the energy deposited in the detector by a g-ray. We will calibrate the system to determine the slope and y-intercept of the calibration function.

Calibration procedure

- Under "Acquire" select the "Clear" option.
- Place the ^{137}Cs sample on the blue sample holder and slide it into the top of the sample chamber rack.
- Make sure the lead absorber is in the slot of the rack just below the sample.
- Close the door on the sample chamber.

- Click the "Go" icon and collect counts until there are several hundred counts in the channels.
- Click the "Stop" icon and save the data in an SPE file on the desktop with a name that you will be able to recognize later.
- Remove the ^{137}Cs and replace it with ^{60}Co .
- Under "Acquire" select the "Clear" option.
- Close the door and collect a sample for the ^{60}Co .
- Save again.

The energy data given for ^{137}Cs and ^{60}Co and the channel number of these peaks are three data points from which you can determine the coefficients of the calibration equation:

$$E_g = m \cdot (\text{channel number}) + b \quad (1)$$

where E_g is the g-ray energy in MeV, m is the slope of the calibration line, and b is the intercept on the energy axis.

	Item	Energy (MeV)	Channel Number
1	^{137}Cs	0.662	
2	1 st ^{60}Co	1.17	
3	2 nd ^{60}Co	1.33	

Table 1

Determine m and b and plot the calibration curve.

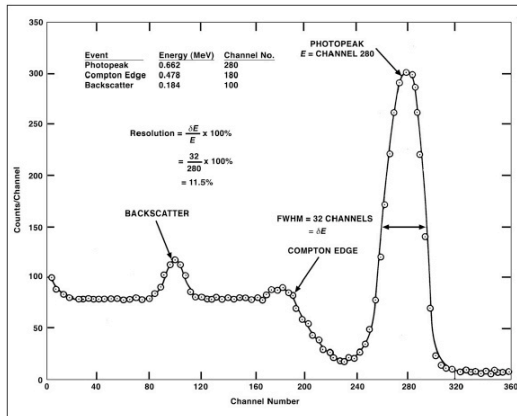


Figure 3. ^{137}Cs peak

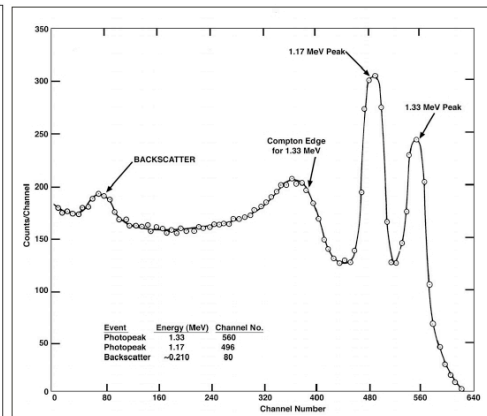


Figure 4. ^{60}Co peaks

EXPERIMENT 3.2. Energy Analysis of an Unknown Gamma Source*

3.2.1 Purpose

This experiment uses the calibrated system of Experiment 3.1 to measure the photopeak energies of an unknown g-ray emitter and to identify the unknown isotopes of the elements.

3.2.2. Procedure

1. Erase the ^{60}Co spectrum from the multichannel analyzer (MCA), but do not change any of the gain calibration settings of the system.

2. Obtain an unknown gamma source from the instructor. Accumulate a spectrum for the unknown sources for a period of time long enough to clearly identify its photopeak(s). From the calibration curve, determine the energy for each photopeak.

EXERCISE

Use the library services of MAESTRO-32 to identify the unknown isotope.

EXPERIMENT 3.3. Spectrum Analysis of ^{60}Co and $^{137}\text{Cs}^*$

3.3.1. Purpose

The purpose of this experiment is to explain some of the features, other than the photopeaks, usually present in a pulse-height spectrum. These are the Compton edge, the Compton continuum and the backscatter peak.

3.3.2. Background Information

The photopeak is created when a gamma-ray photon interacts in the scintillator via the photoelectric effect. The photon encounters an electron that is tightly bound to a nucleus. The entire energy of the photon is transferred to the electron, causing the electron to escape from the atom and the gamma-ray disappears in the process. As the photoelectron travels through the scintillator, it loses its energy by causing additional ionization. At the end of the process, the number of ionized atoms is proportional to the original energy of the photon. As the electrons refill the vacancies in the ionized atoms, visible light photons are generated. This is the source of the scintillation, wherein the number of visible photons is proportional to the original energy of the gamma-ray. Consequently, the event populates the photopeak in the spectrum. This peak is often called the full-energy peak, because a two-step interaction, a Compton scattering followed by a photoelectric interaction, also contributes a small number of events to the full-energy peak.

The Compton interaction is a pure, kinematic collision between a gamma-ray photon and what might be termed a free electron in the NaI(Tl) crystal. By this process, the incident gamma-ray photon gives up only part of its energy to the electron as it bounces off the free electron. The recoiling electron loses energy by causing ionization as it travels through the crystal. Thus the number of visible photons in the resulting scintillation is proportional to the recoil energy of the Compton electron. The amount of energy transferred from the gamma-ray photon to the recoiling electron depends on whether the collision is head-on or glancing. For a head-on collision, the gamma-ray transfers the maximum allowable energy for the Compton interaction. Although it involves a photon and an electron, the interaction is similar to a billiard-ball collision. The reduced energy of the scattered gamma-ray can be determined by solving the energy and momentum conservation equations for the collision. The solution for these equations in terms of the scattered gamma-ray energy can be written as:

$$E_{\gamma'} = \frac{E_{\gamma}}{1 + \frac{E_{\gamma}}{m_0 c^2} (1 - \cos \theta)} \quad (2)$$

where

$E_{\gamma'}$ is the reduced energy of the scattered g-ray, γ' , in MeV,

θ is the scattering angle for the scattered g-ray, relative to the incident g-ray,

E_{γ} is the energy of the incident g-ray, γ , in MeV,

$m_0 c^2 = 0.511$ MeV is the equivalent energy of the rest mass, m_0 , of the electron, and

c is the speed of light.

For a head-on collision, the gamma-ray is scattered backwards along its initial trajectory, and $\theta = 180^\circ$. For this condition, the backscattered gamma-ray energy becomes

$$E_{\gamma'} = \frac{E_\gamma}{1 + 4E_\gamma} \quad (3)$$

Where the convenient approximation, $(m_0c^2)^{-1} \approx 2$, has been used. (Note, for the units to work, the E_g in the denominator is written in terms of MeV, but units are not included with the numerical value.)

If this backscatter event happens in the detector, the maximum energy transferred to the recoiling electron will be

$$E = E_\gamma - E_{\gamma'} \quad (4)$$

Thus the maximum energy that can be recorded in the spectrum for a gamma-ray that interacts in the detector by Compton scattering is given by equation (4). This defines the energy of the Compton edge in Figures 3 and 4. For an initial gamma-ray energy of 1 MeV, equations (5) and (6) predict that the Compton edge will occur at 0.80 MeV, and the energy of the backscattered gamma ray will be 0.20 MeV. Because the gamma-ray photon can be scattered through any angle from 0 to 180° , and the scattered photon can escape the detector, the energy deposited in the detector can vary from the maximum at the Compton edge through all values down to zero. This is the genesis of the Compton continuum in Figures 3 and 4. Find a backscatter peak in this spectrum. It's energy plus the Compton edge energy equals the energy of the main photopeak.

Note that there is a small, but non-zero, probability that the Compton scattered photon will be subsequently absorbed in the crystal by the photoelectric process. This two-step interaction will generate a pulse that falls in the full-energy peak. Compton scattering from an entirely different location causes the backscatter peak in Figures 3 and 4. Consider a gamma ray emitted by the radioactive source in a direction heading away from the detector. This gamma ray can encounter material in the neighborhood of the radioactive source and undergo Compton scattering. If the scattering angle is 180° the scattered gamma ray travels back towards the detector with an energy defined by equation (5). If this lower energy gamma ray interacts in the scintillator by the photoelectric effect, it will contribute to a photopeak at the lower energy. Typically this backscatter peak will be of low intensity, if there is minimal material behind the radioactive source. Usually, the backscatter peak is rather broad, because of the range of directions that can contribute to the peak. For initial gamma-ray energy of 1 MeV, equation (5) predicts that the backscatter peak will occur at 0.20 MeV.

Figure 2 illustrates some of the types of interactions that can take place in the NaI(Tl) detector and the surrounding shielding material.

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- Calculate the energy of the Compton edge for the 0.662-MeV gamma ray from ^{137}Cs . Enter this value in Table 1. From your plot and calibration curve, does this calculation agree with your measured value?
 - Calculate the Compton edge energy for the 1.33-MeV gamma ray from ^{60}Co , and enter that result in Table 1. Is that value in agreement with your ^{60}Co spectrum?
 - From Eq. (3) calculate the backscatter peak energies for the gamma rays from ^{137}Cs and ^{60}Co . Fill in the rest of Table 1. How do your measured energies compare with the theoretical energies from Eq. (3)?
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Energy Resolution

Purpose

The purpose of this part of the experiment is to measure the resolution of the NaI(Tl) detector.

Relevant Equations

The resolution of a spectrometer is a measure of its ability to resolve (i.e., separate) two peaks that are fairly close together in energy. Fig. 3 shows the gamma spectrum that was plotted for the ^{137}Cs source. The resolution of the photopeak is calculated from the following equation:

$$R = \frac{\delta E}{E} \times 100\% \quad (5)$$

Where δE is the Full Width of the peak at Half of the Maximum count level (FWHM) measured in number of channels, and E is the channel number at the centroid, i.e., the mid-point of the photopeak.

In Fig. 3 the photopeak is in channel 280 and its FWHM = 32 channels. From Eq. (5), the resolution is calculated to be 11.5%. It is standard practice to specify the quality of the NaI(Tl) detector by stating its measured percent resolution on the 662-keV photopeak from a ^{137}Cs radioactive source.

The resolution of a scintillation detector depends upon several factors. Among them are:

1. number of photons per scintillation event
2. number of photons that strike the photocathode
3. number of photoelectrons released per photon hitting the photocathode
4. number of photoelectrons that strike the first dynode
5. multiplication factor of the photomultiplier tube

In general, the greater each one of the above quantities is, the better (i.e., smaller) the resolution. Factors 1 and 2 refer to properties of the scintillator while factors 3, 4, and 5 refer to the phototube. Mathematically we say that

$$R = \alpha + \beta/\sqrt{E_\gamma} \quad (6)$$

where R is the resolution, α and β are constants referring to the scintillator and phototube respectively, and E_γ is the g-ray energy. Note that R varies inversely as $\sqrt{E_\gamma}$ and thus the higher E_γ is, the smaller R is.

Procedure

1. Using the same amplifier and analyzer gain settings as before, measure the g-ray spectra of ^{57}Co and ^{54}Mn . Tabulate these results along with those obtained for ^{60}Co and ^{137}Cs .
2. Using the equation (5), which defines resolution, calculate the resolution for each of the g-ray energies and put them in your table.
3. Plot the resolution R vs. $1/\sqrt{E_\gamma}$ and on a second graph, R vs. E_γ . Determine the parameters α and β from the plot of R vs. $1/\sqrt{E_\gamma}$ and equation (6). Discuss the plots that you obtained.

References and Acknowledgements

Sections marked with an * are from the ORTEC® *Experiments in Nuclear Science, Experiment 3: Gamma-Ray Spectroscopy Using NaI(Tl) Detectors*.